



HUNTING FOR SUSY AT ATLAS



Antonella De Santo

Royal Holloway, University of London

Edinburgh, October 3, 2008

Outline

Introduction

The LHC

The ATLAS experiment

Supersymmetry

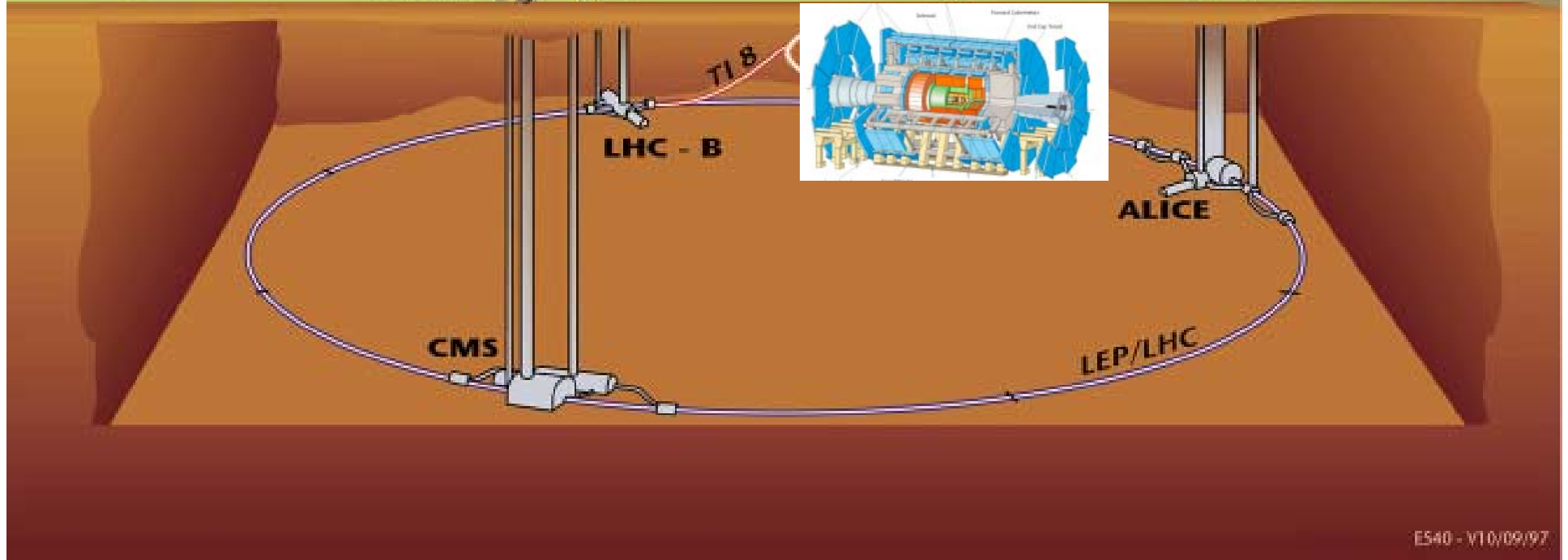
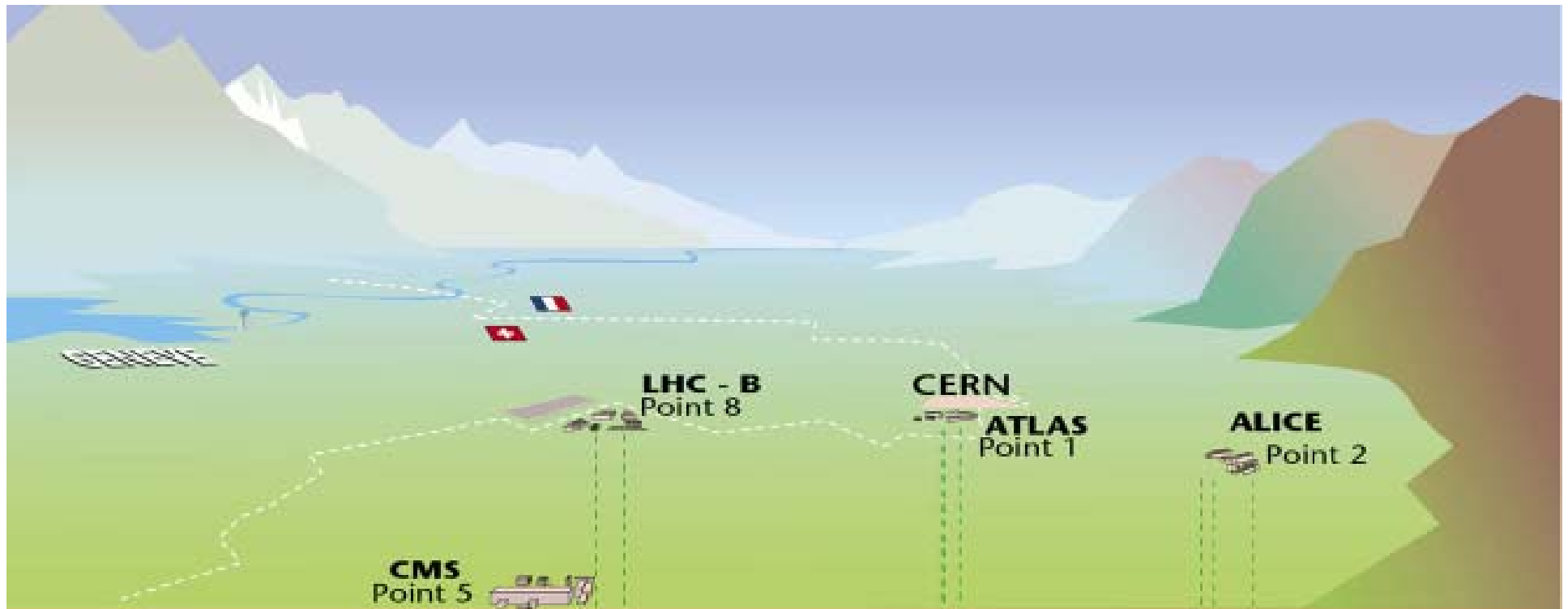
Motivations

Benchmarks and strategy

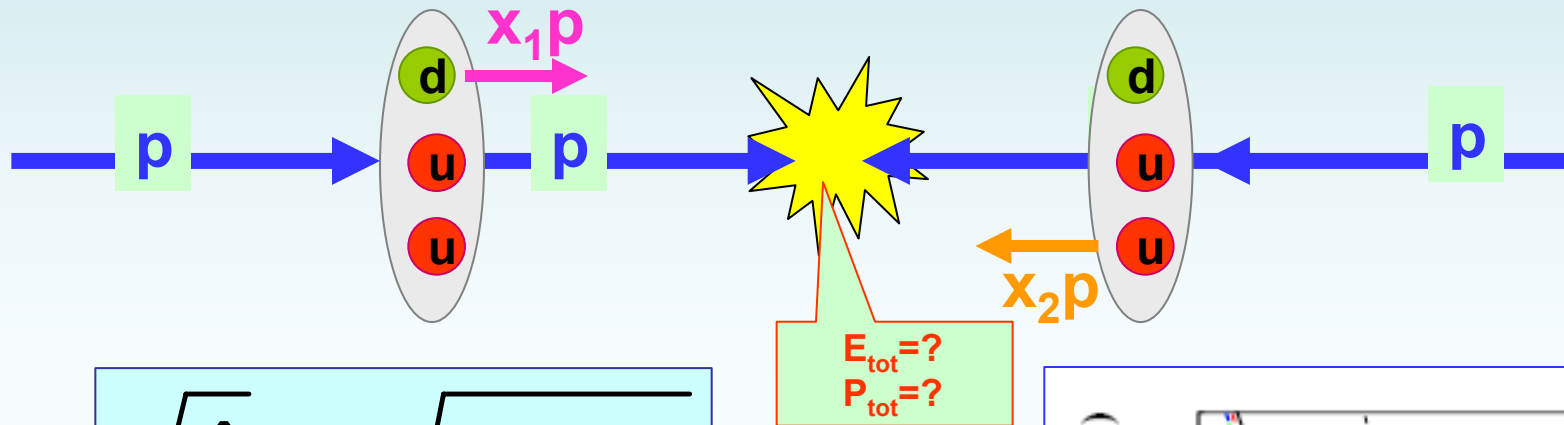
Early physics reach

Conclusions

THE LHC AND ATLAS



Collisions at a hadron collider

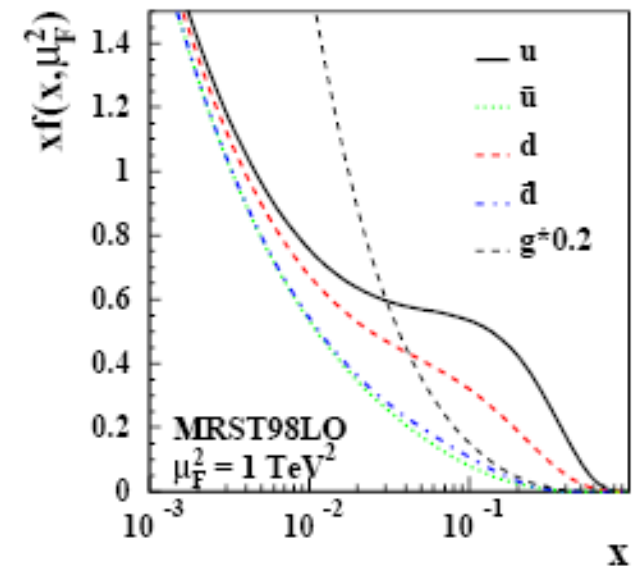


$$\sqrt{\hat{s}} = \sqrt{x_1 x_2 S}$$

No constraints on total initial energy

Broad range of $\sqrt{s} \rightarrow$ good for discovery!

All possible processes are “on” simultaneously



PDFs

LHC's Main Goals

Elucidate mechanism for EW symmetry breaking

Search for Higgs boson in $O(100 \text{ GeV})$ - $O(1 \text{ TeV})$ range

If no light Higgs is found, study WW scattering at high mass

Look for evidence of new physics at TeV-scale

Deviations from Standard Model predictions in data

Supersymmetry

(and Exotics...)

LHC's Main Goals

Elucidate mechanism for EW symmetry breaking

Search for Higgs boson in $\mathcal{O}(100 \text{ GeV})$ - $\mathcal{O}(1 \text{ TeV})$ range

Not Today!

If no light Higgs is found, study WW scattering at high mass

Look for evidence of new physics at TeV-scale

Deviations from Standard Model predictions in data

Supersymmetry

(and Exotics...)

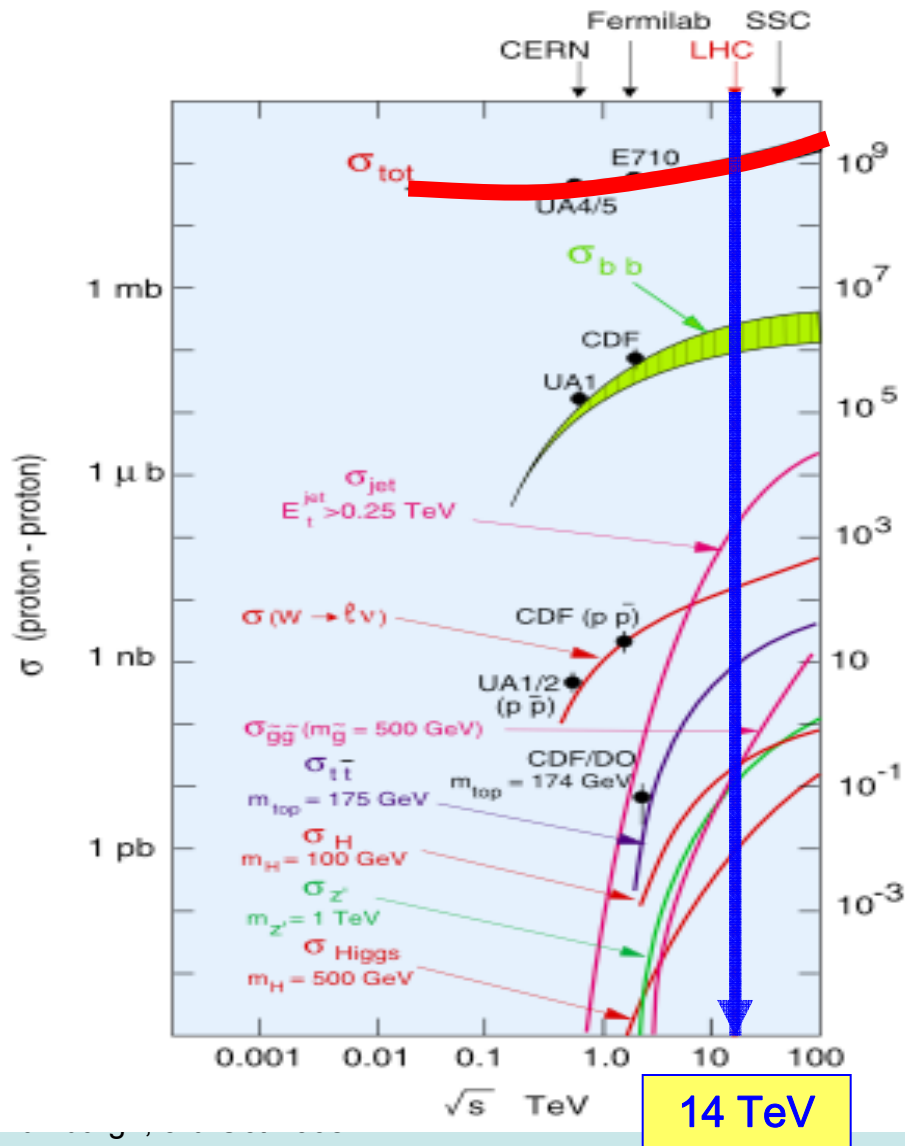
Design luminosity:

“Low” = $(2 \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1})$ ($O(10 \text{ fb}^{-1}) / \text{yr}$)

“High” = $10^{34} \text{ cm}^{-2} \text{ s}^{-1}$ ($O(100 \text{ fb}^{-1}) / \text{yr}$)

$$\sigma_{pp}^{tot} \sim 110 \text{ mb} \quad (\sigma_{pp}^{inel} \sim 70 \text{ mb})$$

“Minimum Bias” rate:
 $O(10^9 \text{ Hz}) !!$



Events / sec for $\mathcal{L} = 10^{34} \text{ cm}^{-2} \text{ sec}^{-1}$

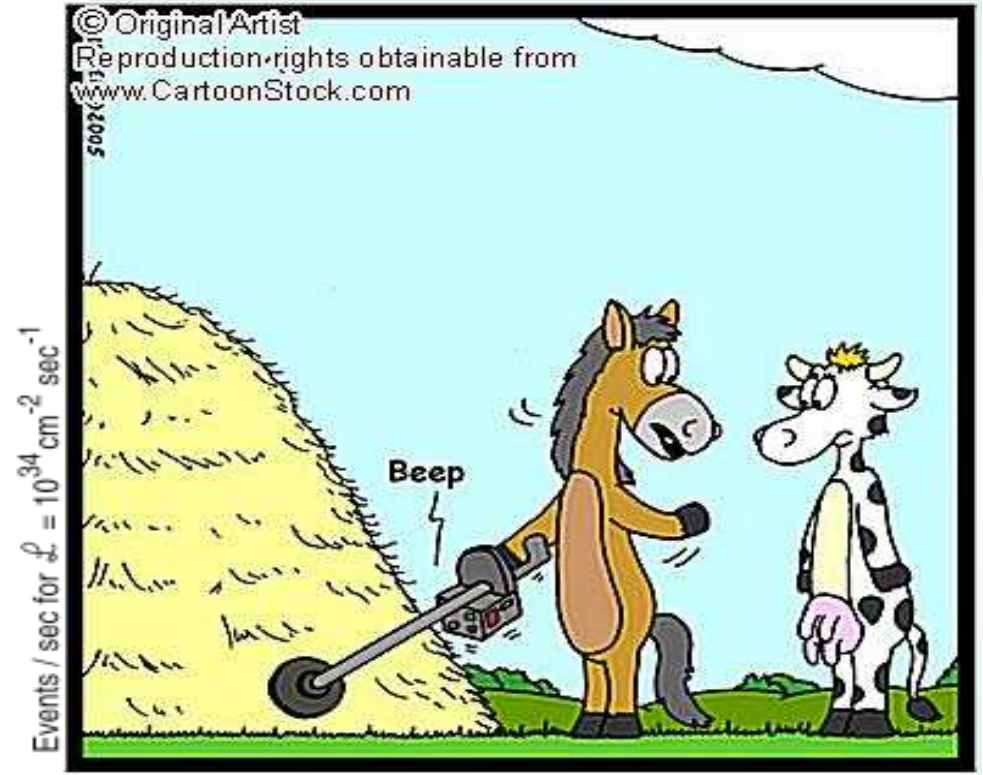
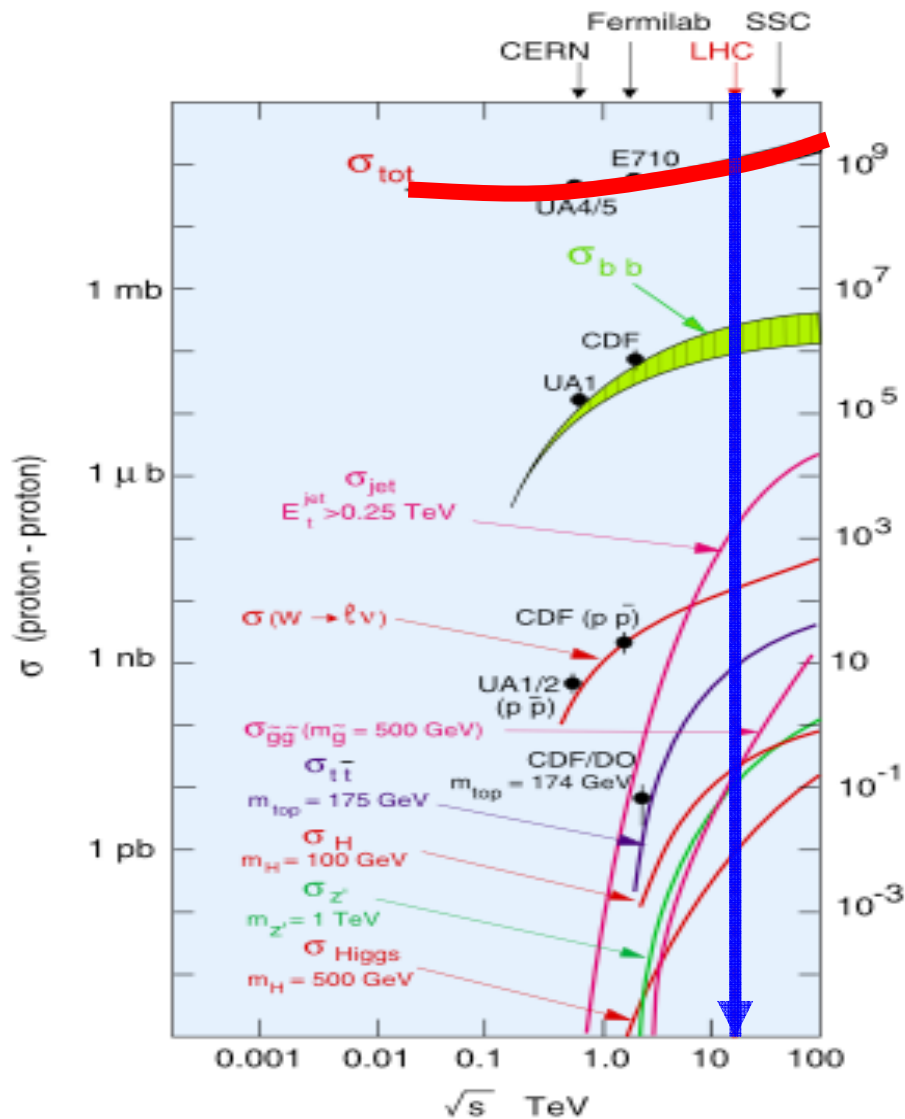
Process	σ (nb)	#evts [10 fb^{-1}]	Rates(Hz) [“high” L]
$b\bar{b}$	$500 \mu\text{b}$	5×10^{12}	5×10^6
$W \rightarrow e\nu$	15 nb	$\sim 10^8$	150
$Z \rightarrow ee$	1.5 nb	$\sim 10^7$	15
$t\bar{t}$	800 pb	$\sim 10^7$	10
$\tilde{g}\tilde{g}$ (1 TeV)	$\sim 1 \text{ pb}$	$\sim 10^4$	10^{-2}
$H(200 \text{ GeV}) \rightarrow 4\ell$	$\sim 10 \text{ fb}$	$\sim 10^2$	10^{-4}

Design luminosity:

“Low” = $(2 \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1})$ ($O(10 \text{ fb}^{-1}) / \text{yr}$)

“High” = $10^{34} \text{ cm}^{-2} \text{ s}^{-1}$ ($O(100 \text{ fb}^{-1}) / \text{yr}$)

$$\sigma_{pp}^{tot} \sim 110 \text{ mb} \quad (\sigma_{pp}^{inel} \sim 70 \text{ mb})$$



You were right: There's a needle in this haystack...

Must rely on distinctive signatures
leptons, photons, b-jets, missing E_T , ...

Detector Requirements

Excellent position and momentum resolution in central tracker

b-jets, taus

Excellent ECAL performance

electrons, photons

v. good granularity (energy and position measurements)

Good HCAL performance

jets, Emiss (neutrinos, SUSY stable LSP, etc)

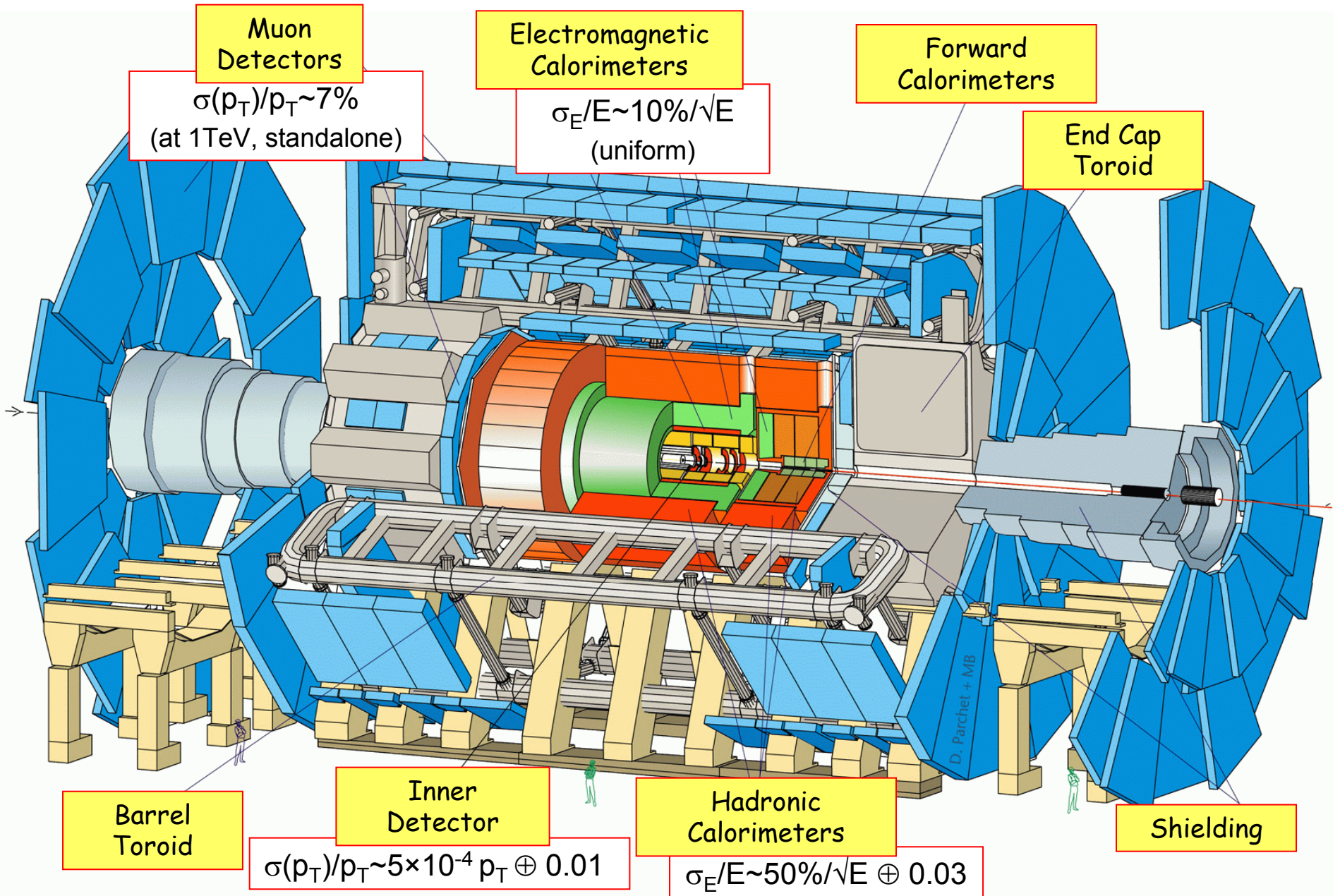
good granularity (energy and position measurements)

good η coverage (hermeticity for Emiss measurements)

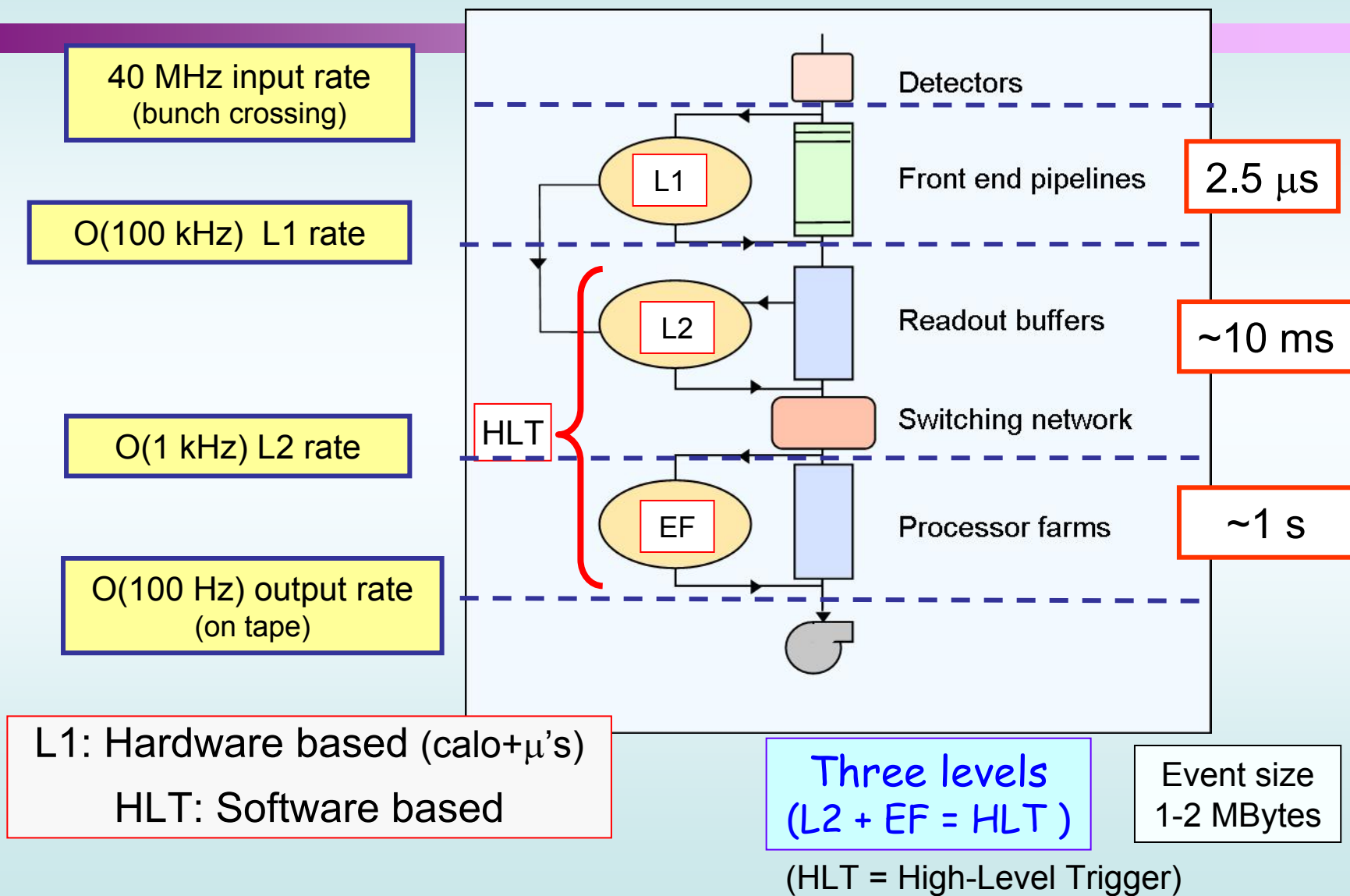
Excellent muon identification and momentum resolution

from “combined” muons in external spectrometer + central tracker

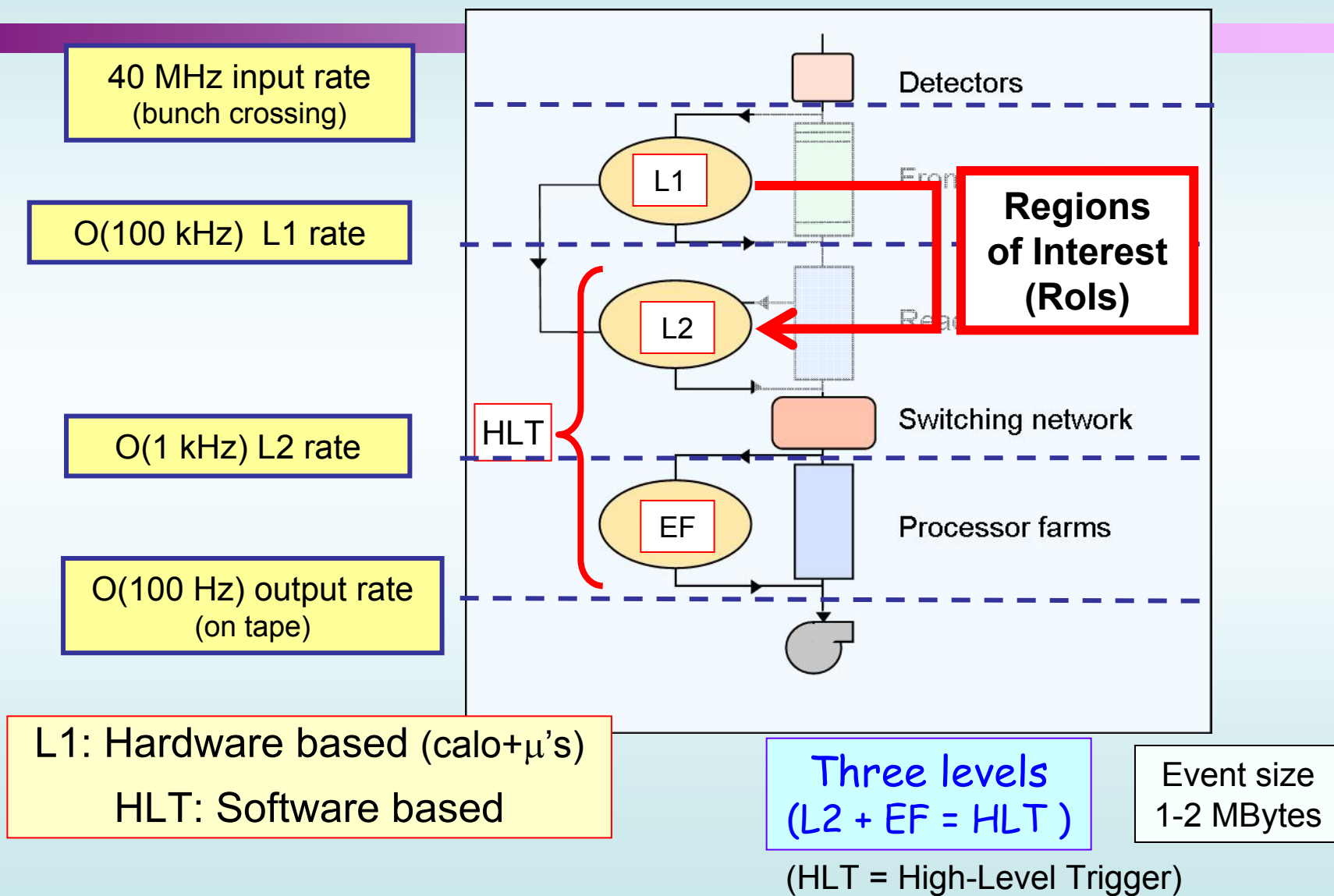
The ATLAS Detector



The ATLAS Trigger

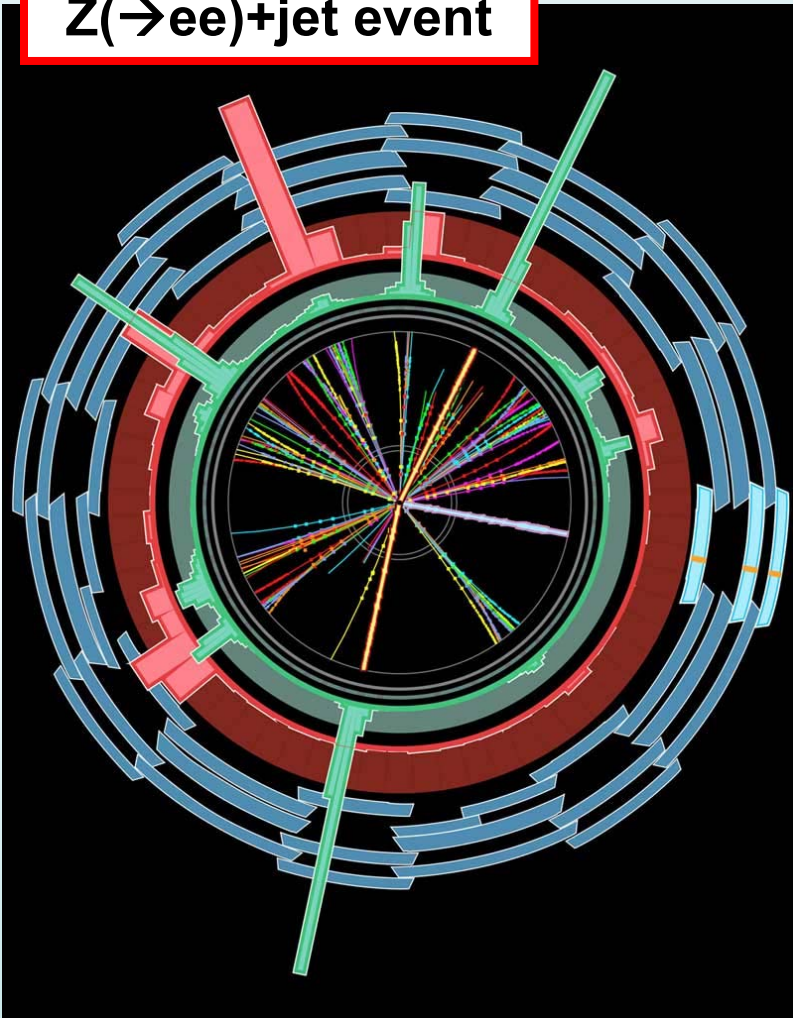


The ATLAS Trigger



The ATLAS High-Level Trigger

Z($\rightarrow ee$)+jet event



“Seeded” and “stepwise”

early rejection of uninteresting events
minimum amount of processing
maximum flexibility

Rolls “seed” trigger reconstruction chain

<10% of the event accessible at L2
EF seeded by L2 (“offline” reco)

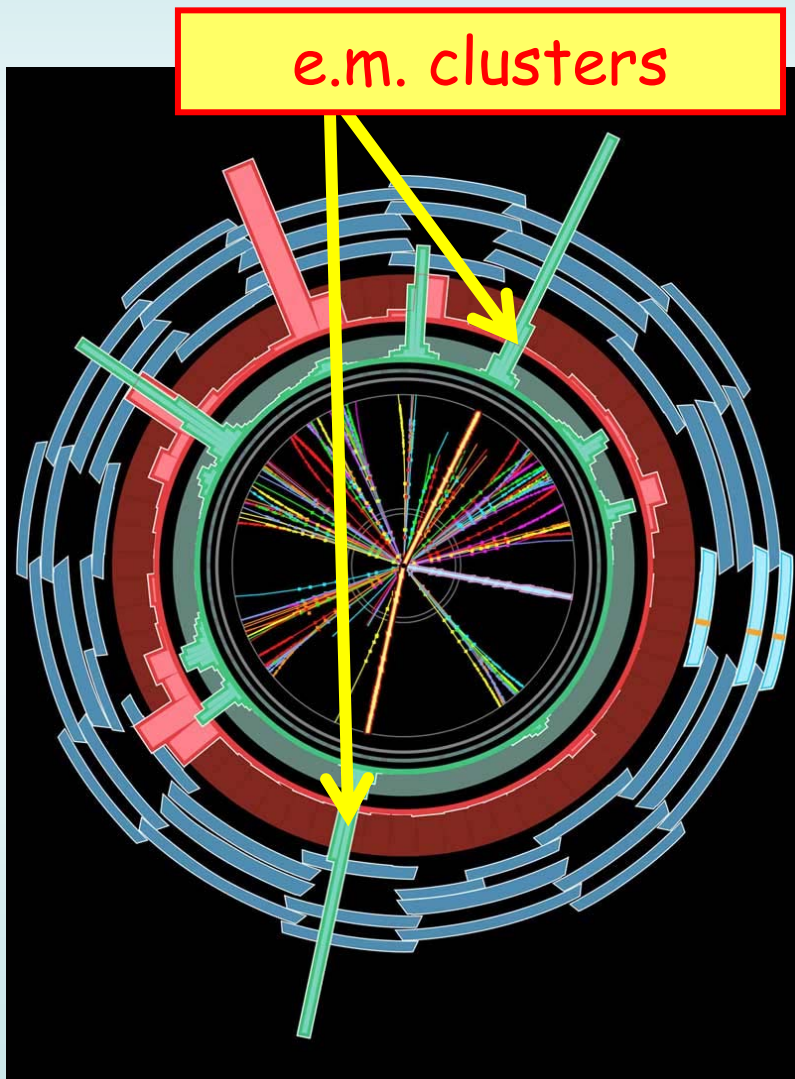
Reduced CPU/ bandwidth requirements

but increased complexity

Avoid biases but retain flexibility

to account for the “unexpected”

The ATLAS High-Level Trigger



“Seeded” and “stepwise”

early rejection of uninteresting events
minimum amount of processing
maximum flexibility

Runs “seed” trigger reconstruction chain

<10% of the event accessible at L2
EF seeded by L2 (“offline” reco)

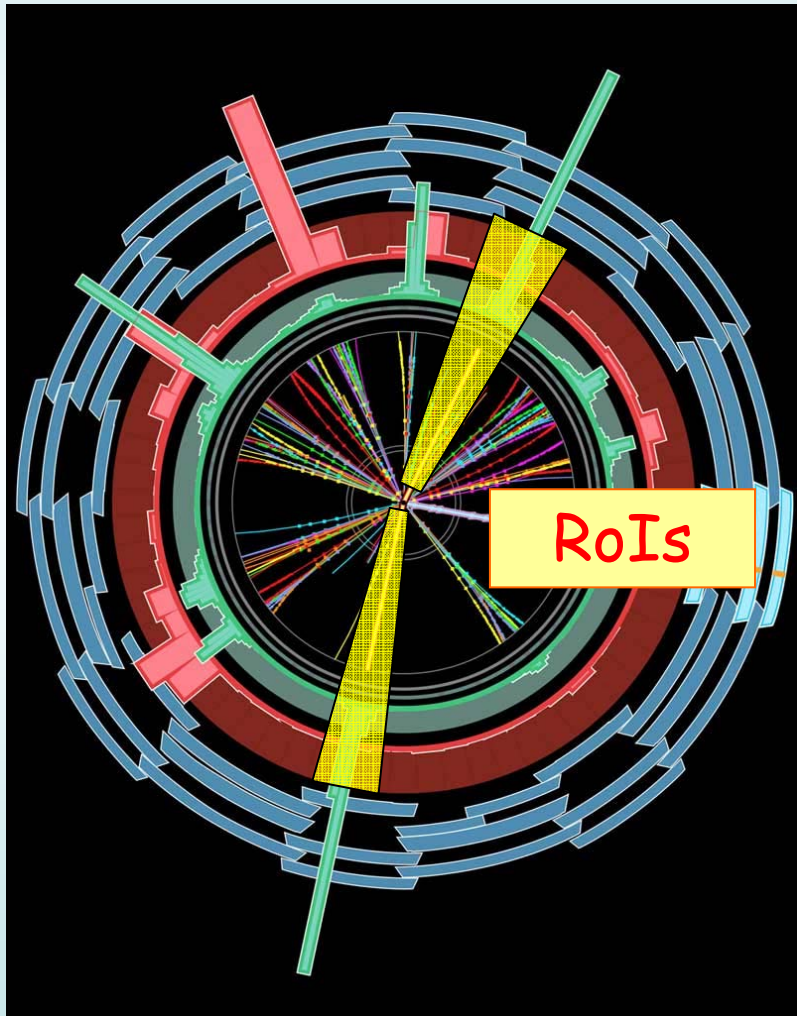
Reduced CPU/ bandwidth requirements

but increased complexity

Avoid biases but retain flexibility

to account for the “unexpected”

The ATLAS High-Level Trigger



“Seeded” and “stepwise”

early rejection of uninteresting events
minimum amount of processing
maximum flexibility

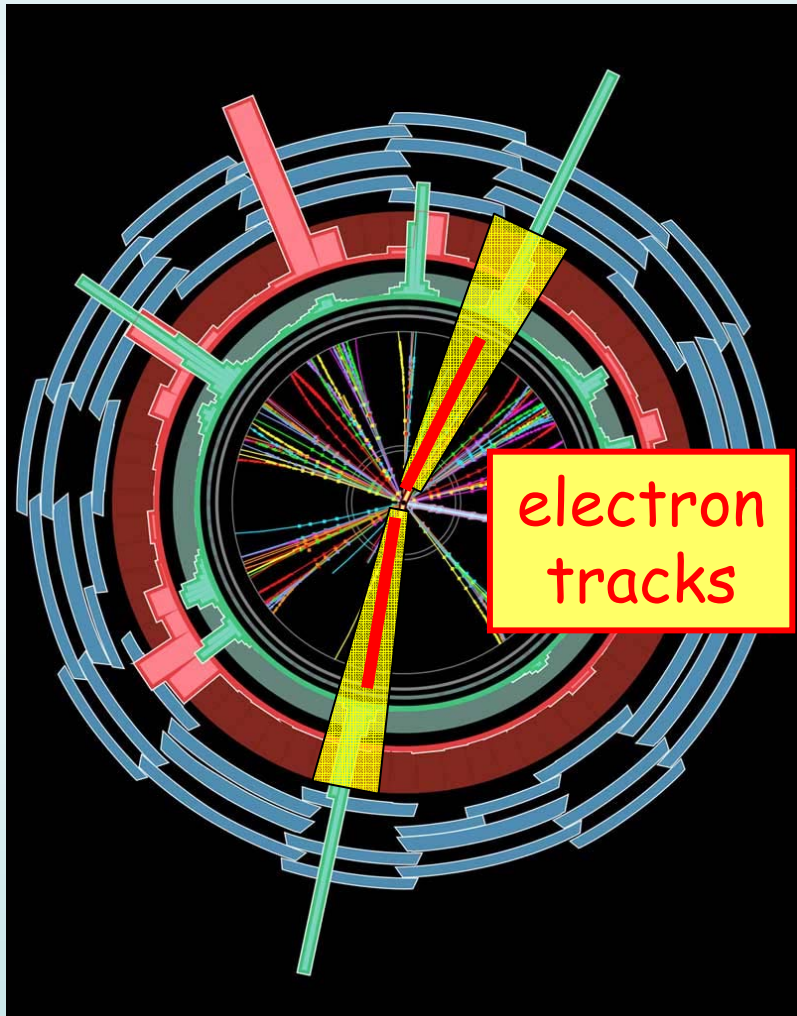
Rols “seed” trigger reconstruction chain

<10% of the event accessible at L2
EF seeded by L2 (“offline” reco)

Reduced CPU/ bandwidth requirements
but increased complexity

Avoid biases but retain flexibility
to account for the “unexpected”

The ATLAS High-Level Trigger



“Seeded” and “stepwise”

early rejection of uninteresting events
minimum amount of processing
maximum flexibility

Roles “seed” trigger reconstruction chain

<10% of the event accessible at L2
EF seeded by L2 (“offline” reco)

Reduced CPU/ bandwidth requirements

but increased complexity

Avoid biases but retain flexibility

to account for the “unexpected”

Early Physics Triggers for SUSY

Early SUSY searches will mainly concentrate on inclusive signatures

multi-jets

large missing transverse energy

possibly leptons

(see later)

Crucial to keep the trigger selection criteria as simple as possible
to minimize biases and systematic effects

Experience shows that an E_{miss} trigger takes time to become established

due to instrumental effects

E_{miss} also crucial to define control/signal regions

Strategy is to de-emphasize E_{miss} trigger in early days
rely on leptons and/or jets instead



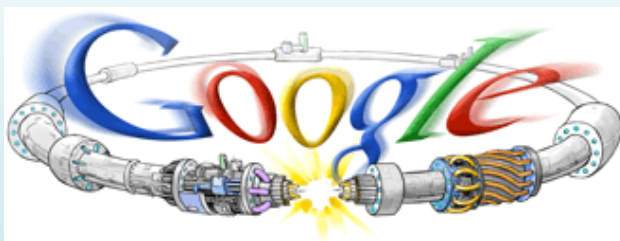
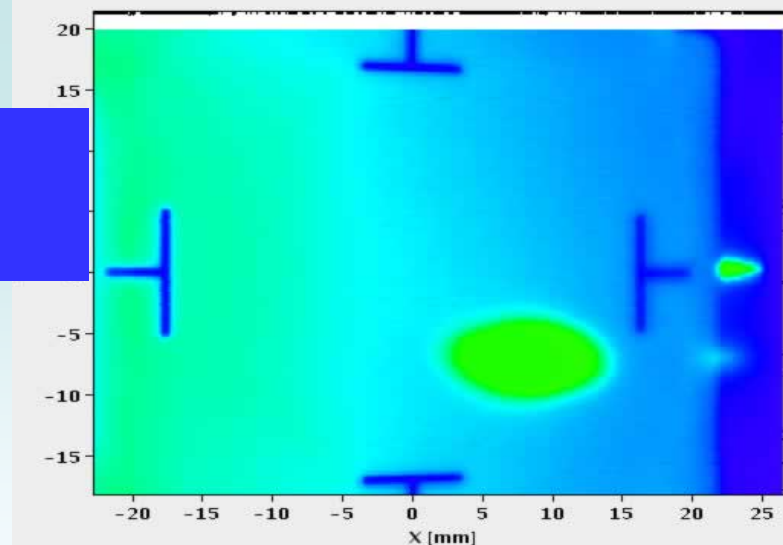
STATUS

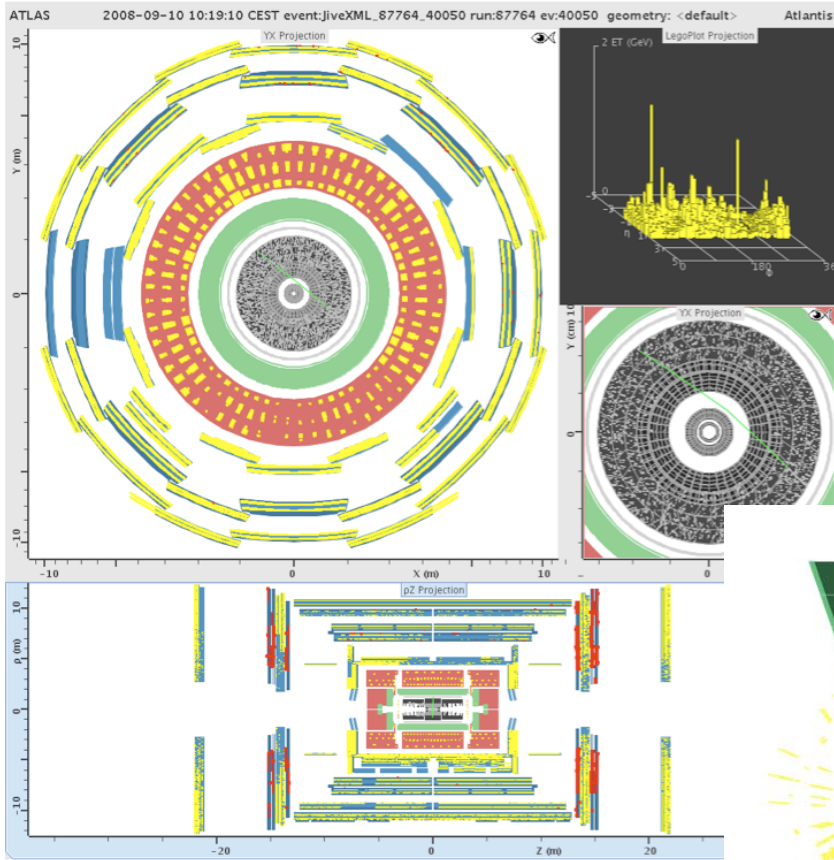
The world did not come to an end...



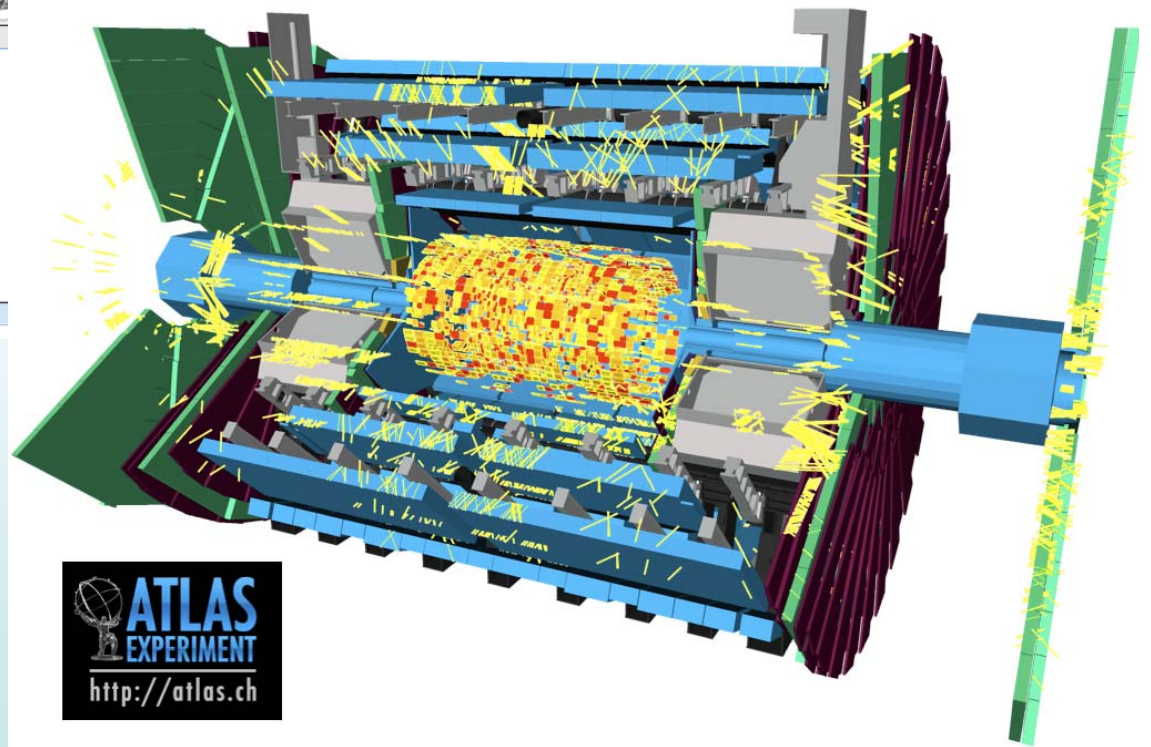
September 10, 2008

A Great Start...





**First beam event
in ATLAS**



Edinburgh, 3rd Oct 2008

Unfortunately...

Incident in LHC sector 34

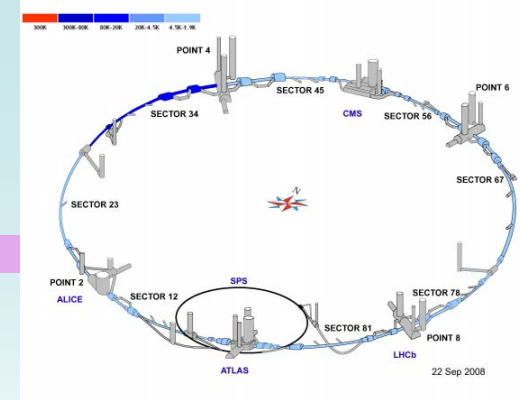
Geneva, 20 September 2008.

During commissioning (without beam) of the final LHC sector (sector 34) at high current for operation at 5 TeV, an incident occurred at mid-day on Friday 19 September resulting in a large helium leak into the tunnel. Preliminary investigations indicate that the most likely cause of the problem was a faulty electrical connection between two magnets, which probably melted at high current leading to mechanical failure. CERN's strict safety regulations ensured that at no time was there any risk to people.

A full investigation is underway, but it is already clear that the sector will have to be warmed up for repairs to take place. This implies a minimum of two months down time for LHC operation. For the same fault, not uncommon in a normally conducting machine, the repair time would be a matter of days.

Further details will be made available as soon as they are known.

(R.Aymar, CERN DG)



Unfortunately...

Incident in LHC sector 34

Geneva, 20 September 2008.

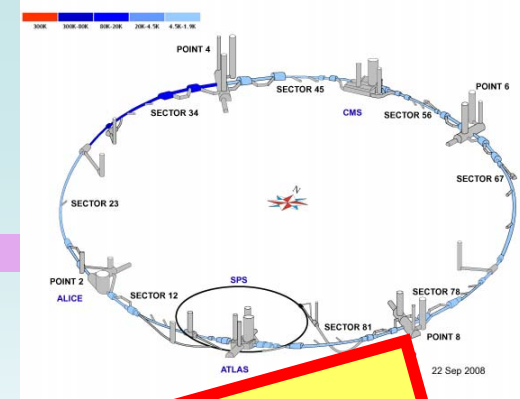
During commissioning (without beam) of the final LHC high current for operation at 5 TeV, an incident occurred on 19 September resulting in a large helium leak. Investigations indicate that the fault was caused by an electrical connection problem in the current leading to a short circuit that caused the helium leak.

First collisions expected in 2009

The sector will have to be repaired, which implies a minimum of two months. This is the same fault, not uncommon in a new machine. Repair time would be a matter of days.

Further information will be made available as soon as they are known.

(R.Aymar, CERN DG)



SUPERSYMMETRY

Why go Beyond the Standard Model?

Despite its many successes, the Standard Model is widely believed to be only an effective theory, valid up to a scale $\Lambda \ll M_{\text{Planck}}$

Gravity not included in SM

Hierarchy/naturalness problem:

$$M_{\text{EW}} \ll M_{\text{Planck}}$$

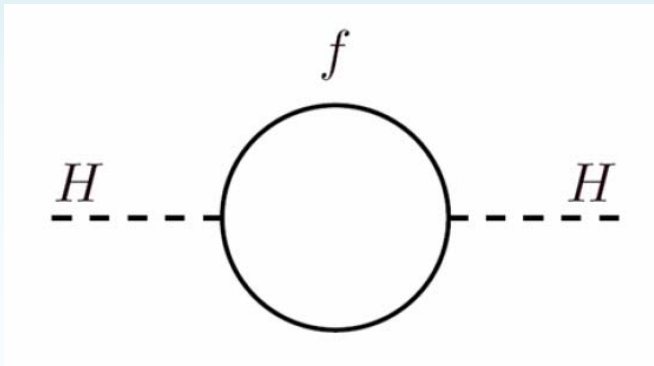
Fine-tuning

Unification of couplings

Need a more fundamental theory of which the SM is only a low-energy approximation

Hierarchy Problem and Naturalness

In SM, loop corrections to Higgs boson mass:



$$\delta m_H^2 = O\left(\frac{\alpha}{\pi}\right) \Lambda^2$$

Theory cut-off

Natural scale of scalar mass is very large!

These corrections, which are large, give rise to fundamental problems when requiring that:

$m_H \ll$ fundamental mass scale (i.e. M_{Planck})

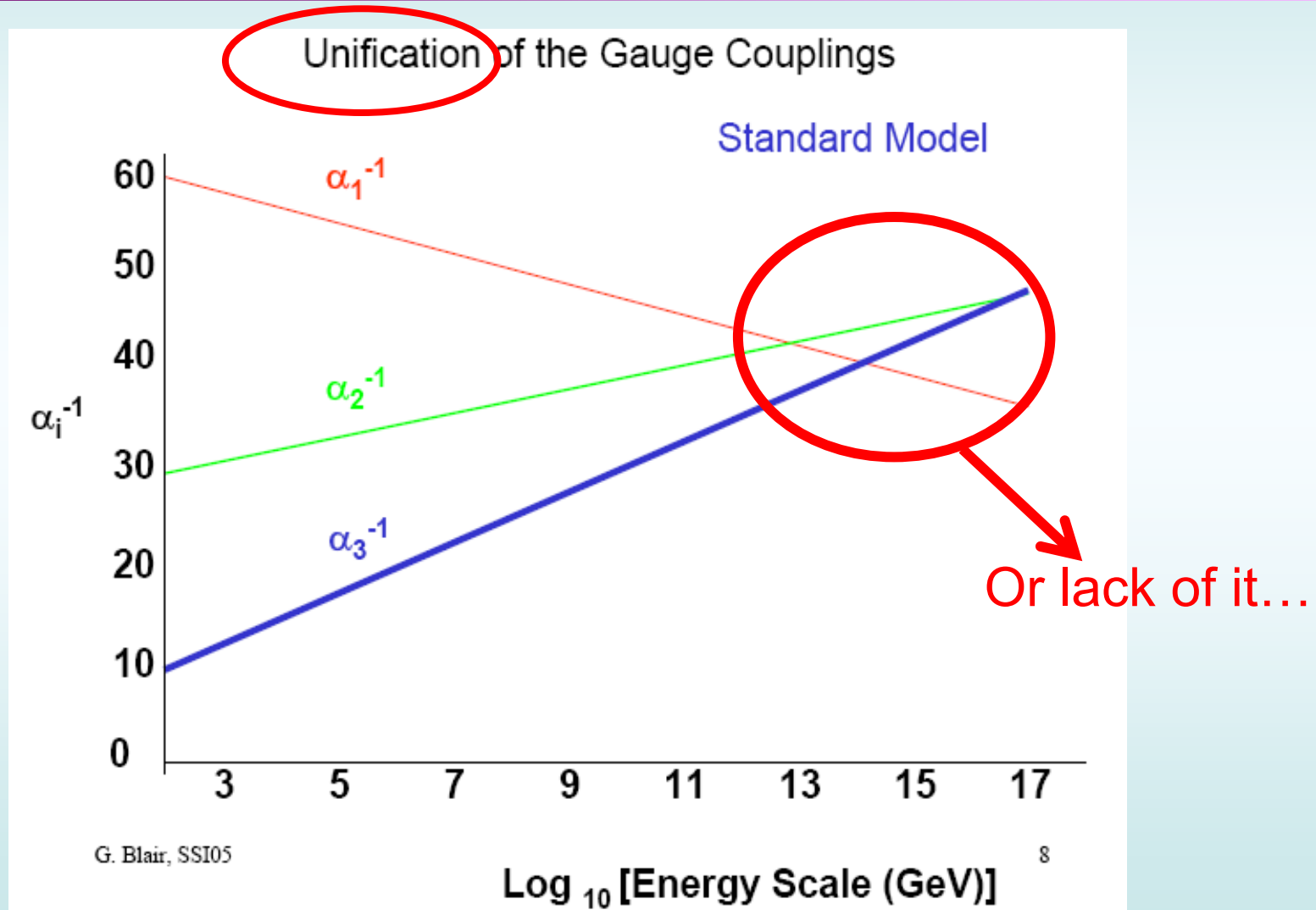
(hierarchy problem)

Corrections δm_H^2 to Higgs mass should not be $\gg m_H^2$

(naturalness)

Need either fine-tuning or protective symmetry!

Unification of Coupling Constants in the SM



Supersymmetry

Space-time symmetry that relates fermions (matter) and bosons (interactions)

$$Q|boson\rangle = |fermion\rangle \quad \text{and} \quad Q|fermion\rangle = |boson\rangle$$

Further doubling of the particle spectrum

Every SM field has a “superpartner” with the same mass

Spin differs by 1/2 between SUSY and SM partners

Identical gauge numbers

Identical couplings

Superpartners have not been observed

SUSY must be a broken symmetry

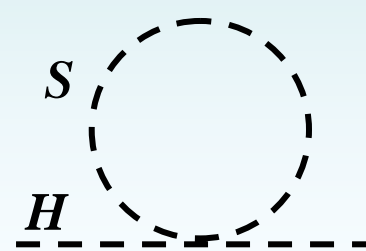
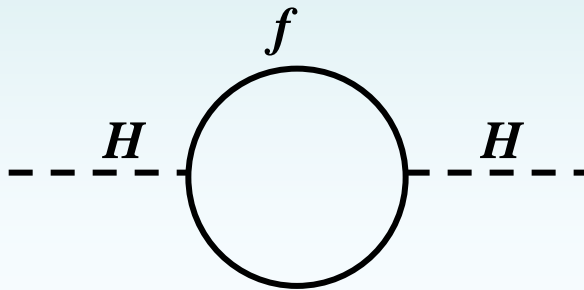
But SUSY-breaking terms in Lagrangian must not re-introduce quadratic divergences in theory !

Minimal Supersymmetric Standard Model

Standard Model Particles and Fields		Supersymmetric Partners			
		Interaction Eigenstates		Mass Eigenstates	
Symbol	Name	Symbol	Name	Symbol	Name
$q = u, d, c, s, t, b$	quark	\tilde{q}_L, \tilde{q}_R	squark	\tilde{q}_1, \tilde{q}_2	squark
$l = e, \mu, \tau$	lepton	\tilde{l}_R, \tilde{l}_L	slepton	\tilde{l}_1, \tilde{l}_2	slepton
$l = \nu_e, \nu_\mu, \nu_\tau$	neutrino	$\tilde{\nu}$	sneutrino	$\tilde{\nu}$	sneutrino
g	gluon	\tilde{g}	gluino	\tilde{g}	gluino
W^\pm	W-boson	\tilde{W}^\pm	wino	$\tilde{\chi}_{1,2}^\pm$	chargino
H_u^+, H_d^-	charged Higgs boson	$\tilde{H}_u^+, \tilde{H}_d^-$	charged higgsino		
B	B-field	\tilde{B}	bino	$\tilde{\chi}_{1,2,3,4}^0$	neutralino
W^0	W^0 -field	\tilde{W}^0	wino		
H_u^0, H_d^0	neutral Higgs boson	$\tilde{H}_u^0, \tilde{H}_d^0$	neutral higgsino		

Supersymmetric Solution to Divergences

Bosonic and fermionic diagrams now give **equal and opposite** contributions:



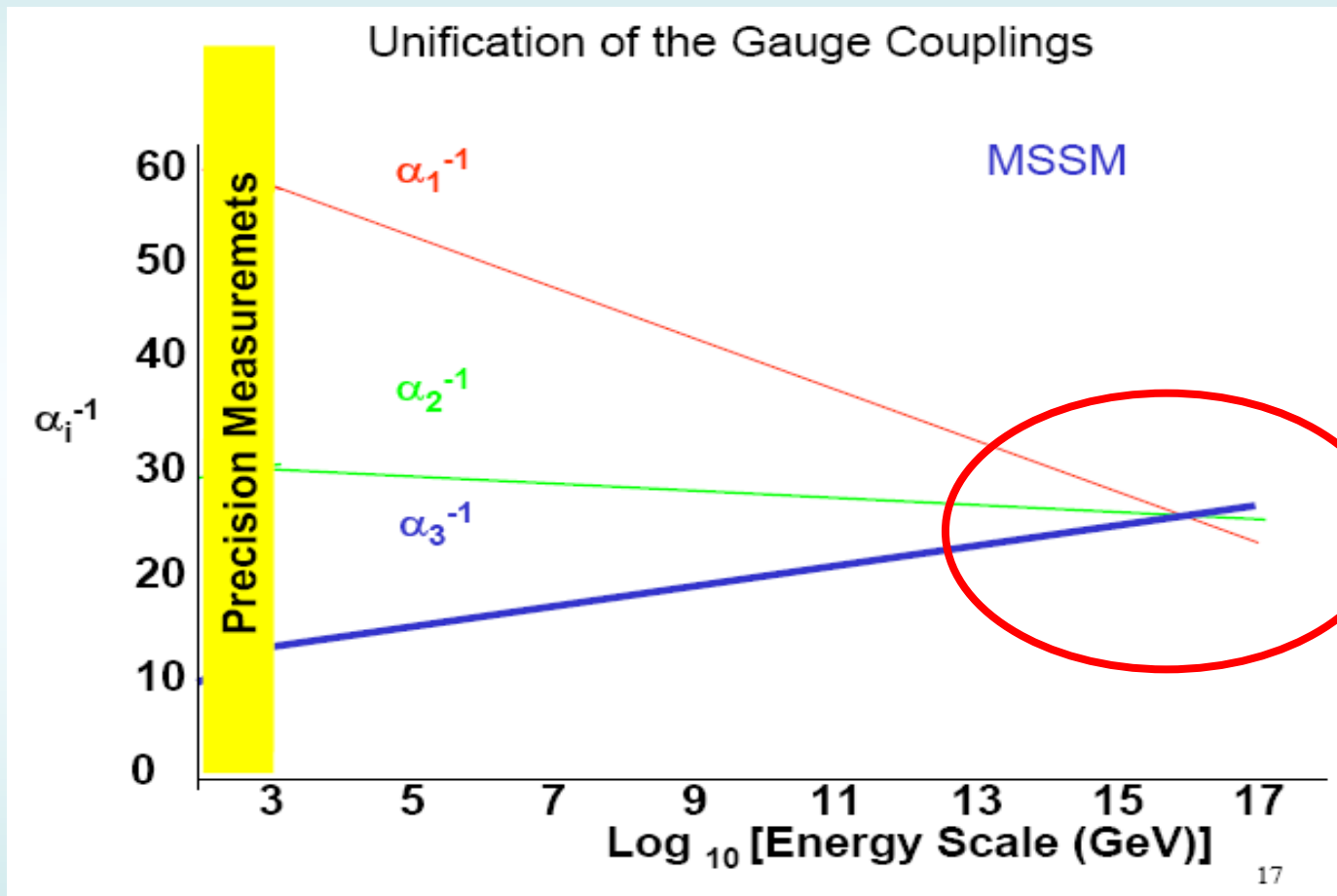
$$\delta m_H^2 = \left(\frac{g_f^2}{16\pi^2} \right) (\Lambda^2 + m_f^2) - \left(\frac{g_s^2}{16\pi^2} \right) (\Lambda^2 + m_s^2) = \mathcal{O}\left(\frac{\alpha}{4\pi}\right) |m_s^2 - m_f^2|$$

If #boson = #fermions and they have equal masses and couplings, the quadratic divergences cancel

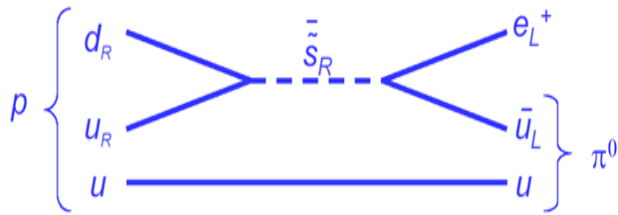
Higgs mass correction $\delta m_H^2 < m_H^2$ if $|m_s^2 - m_f^2| < \sim \text{TeV}^2$

Gauge boson contribution cancelled by gaugino contribution

Unification of Coupling Constants in MSSM



Now unification of strong, weak and e.m. forces achieved at $\sim M_{\text{GUT}}$



R-parity

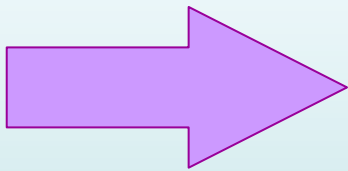
A new symmetry is introduced to cure unwanted effects (e.g. proton decay) from lepton and baryon number violating terms in MSSM:

$$R_p = (-1)^{3(B-L)+2S}$$

$$R_p = +1 \text{ (SM)}$$

$$R_p = -1 \text{ (SUSY)}$$

R_p conservation



Stable LSP (=Lightest SUSY Particle)

typically the lightest neutralino (good DM candidate)

Pair-production of sparticles

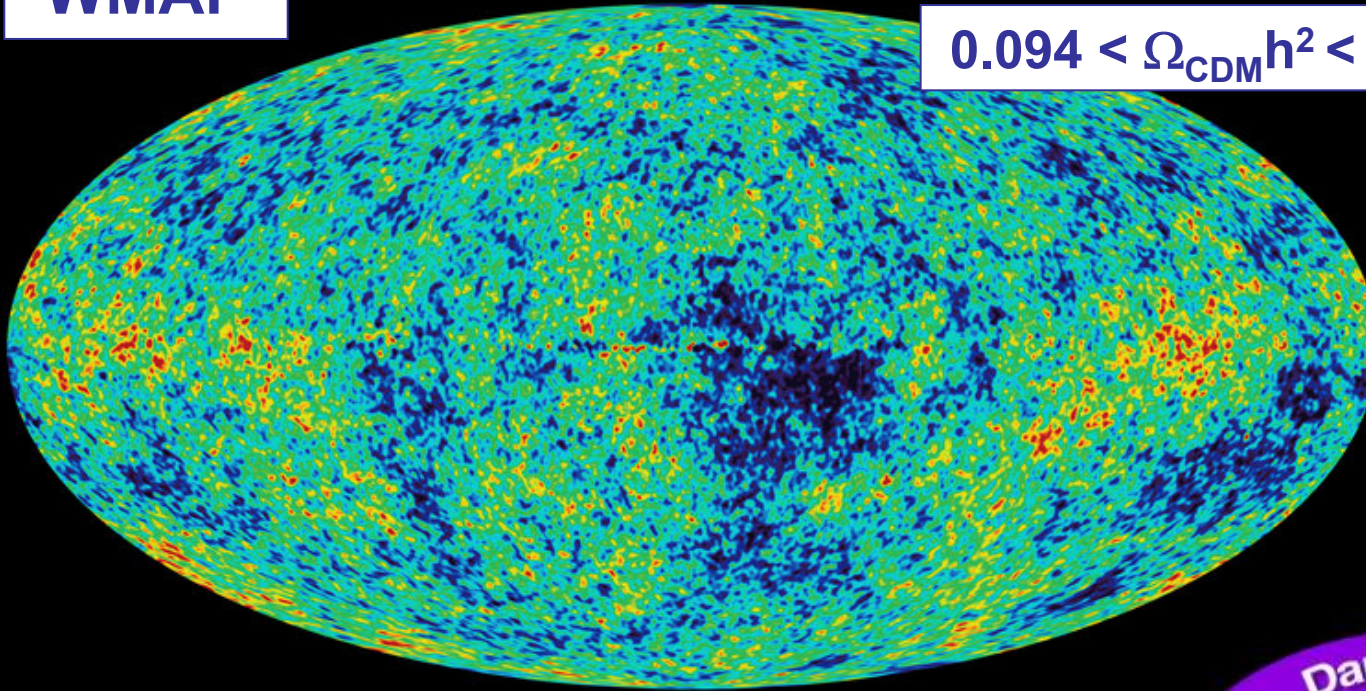
in scattering of SM particles (e.g. pp at the LHC)

In the following, will assume R-parity conservation

Neutralino as Dark Matter Constituent

WMAP

$$0.094 < \Omega_{\text{CDM}} h^2 < 0.136 \quad (95\% \text{ CL})$$

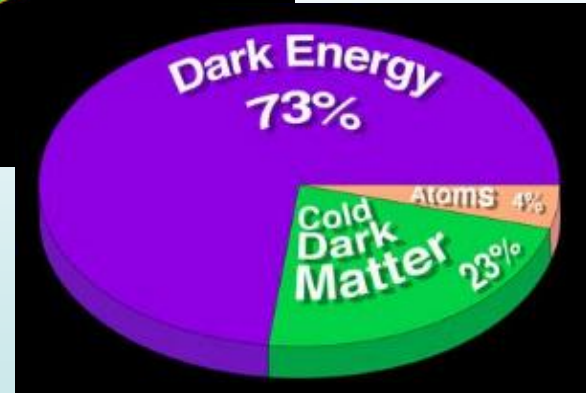


Neutralino LSP is a good DM candidate

stable

electrically neutral

weakly and gravitationally interacting



Benchmarks and Strategy

Baseline paradigm is R_p -conserving mSUGRA:

- m_0 universal scalar mass
- $m_{1/2}$ universal gaugino mass
- A_0 trilinear soft breaking parameter at GUT scale
- $\tan\beta$ ratio of Higgs vevs
- $\text{sgn}(\mu)$ sign of SUSY Higgs mass term

$$R_p = (-1)^{3(B-L)+2S}$$

Stable LSP
(good DM candidate)

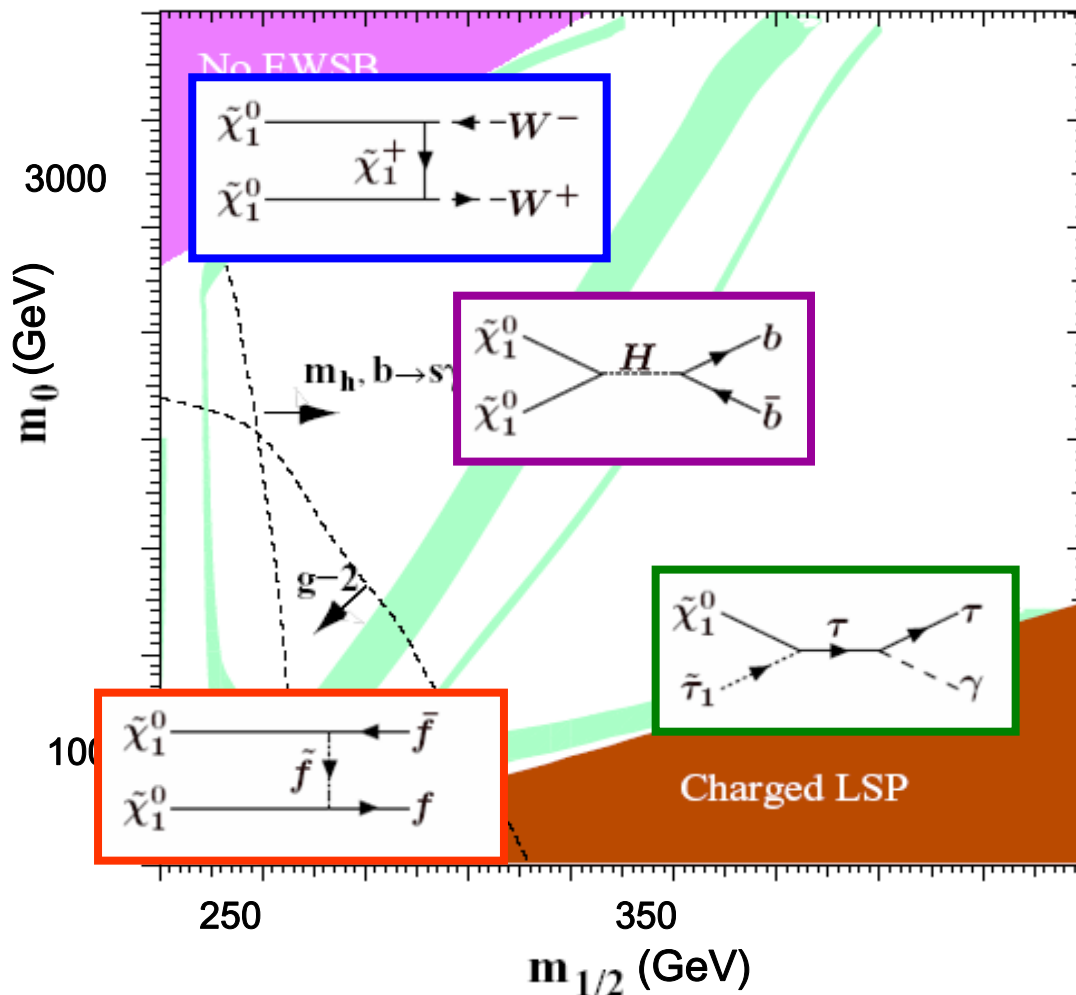
Other scenarios are also considered (GMSB, AMSB, ...)

Full coverage of signatures and topologies in the entire phase space

If it's there, don't miss it !!

mSUGRA Benchmarks

Four WMAP-compatible regions, with different mechanisms for neutralino annihilation and rather different observable phenomenology



bulk

neutralino mostly bino, annihilation to ff via sfermion exchange

focus point

neutralino has strong higgsino component, annihilation to WW, ZZ

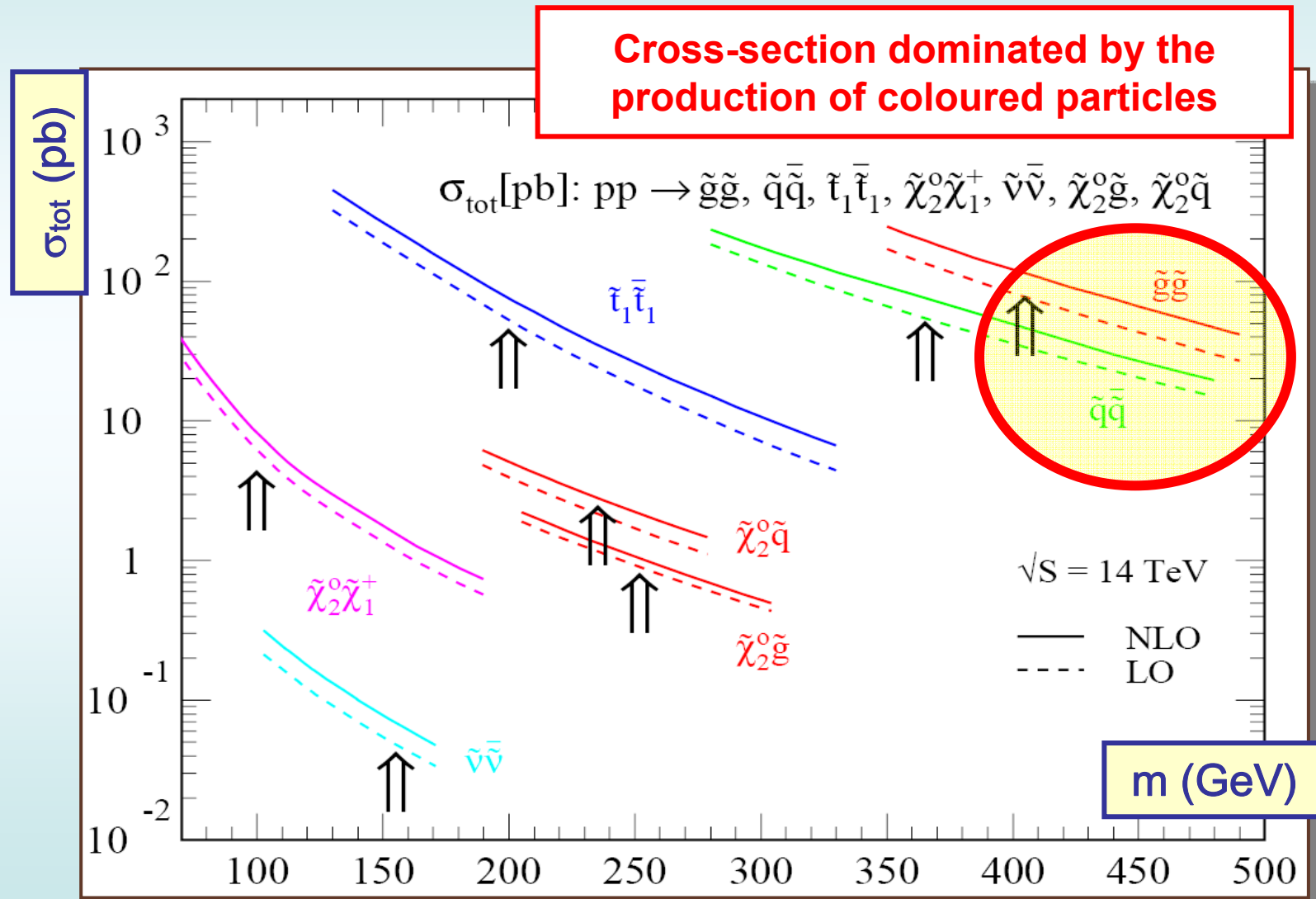
co-annihilation

pure bino, small NLSP-LSP mass difference, typically coannihilation with stau

Higgs funnel

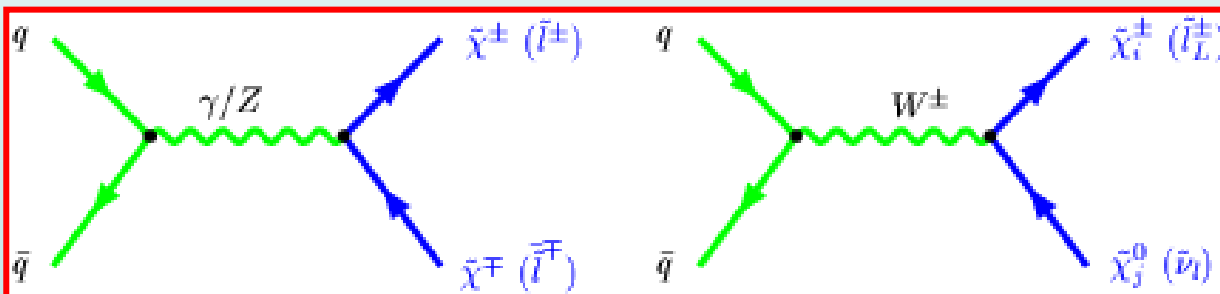
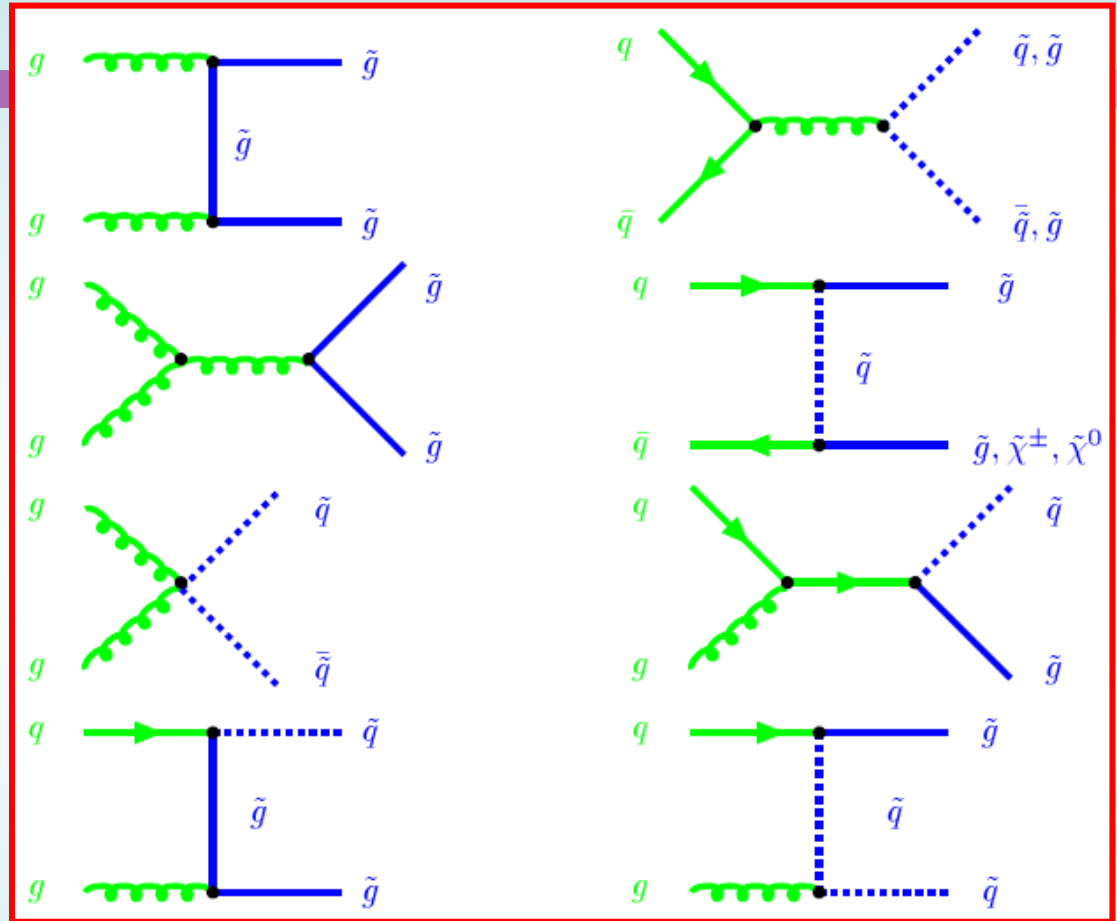
decay to fermion pair through resonant A exchange – high $\tan\beta$

SUSY Cross-Sections



Production Mechanisms

Squark/Gluino Production



Direct Gaugino Production

Possible Final States

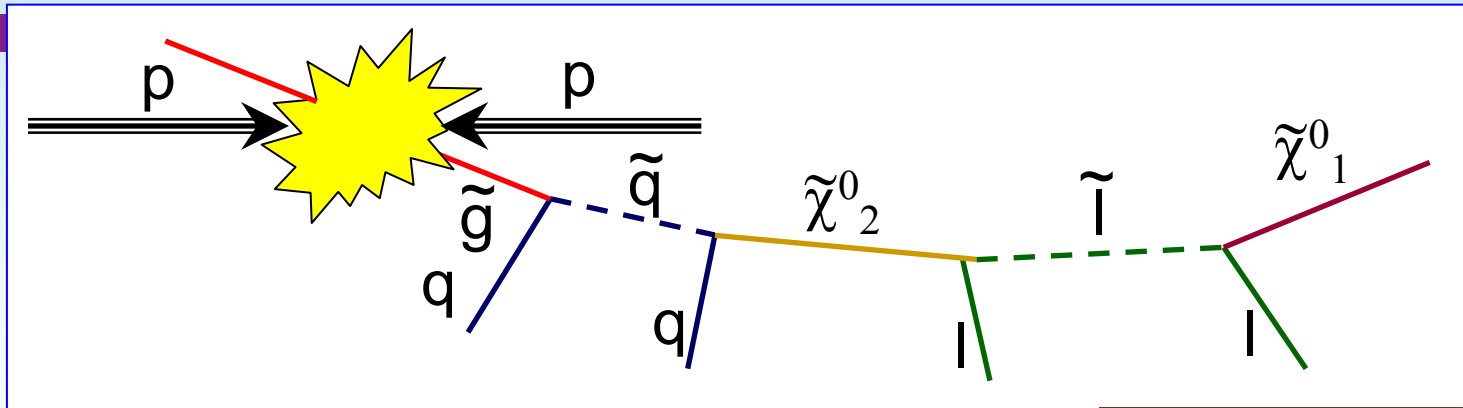
<u>Production</u>	<u>Key Decay Modes</u>	<u>Signatures</u>
<ul style="list-style-type: none"> $\tilde{g}\tilde{g}, \tilde{q}\tilde{q}, \tilde{g}\tilde{q}$ 	$\left. \begin{array}{l} \tilde{g} \rightarrow q\bar{q}\tilde{\chi}_1^0 \\ q\bar{q}'\tilde{\chi}_1^\pm \\ g\tilde{\chi}_1^0 \\ \tilde{q} \rightarrow q\tilde{\chi}_i^0 \\ \tilde{q} \rightarrow q'\tilde{\chi}_i^\pm \end{array} \right\} \begin{array}{l} m_{\tilde{q}} > m_{\tilde{g}} \\ m_{\tilde{g}} > m_{\tilde{q}} \end{array}$	$\cancel{E}_T + \text{multijets}$ (+leptons)

Inclusive Searches

<ul style="list-style-type: none"> $\tilde{\chi}_1^\pm \tilde{\chi}_2^0$ 	$\tilde{\chi}_1^\pm \rightarrow \tilde{\chi}_1^0 l^\pm \nu, \tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 ll$	Trilepton + \cancel{E}_T
	$\tilde{\chi}_1^\pm \rightarrow \tilde{\chi}_1^0 q\bar{q}', \tilde{\chi}_2^0 \rightarrow \tilde{\chi}_1^0 ll$	Dilepton + jet + \cancel{E}_T
<ul style="list-style-type: none"> $\tilde{\chi}_1^+ \tilde{\chi}_1^-$ 	$\tilde{\chi}_1^+ \rightarrow l\tilde{\chi}_1^0 l^\pm \nu$	Dilepton + \cancel{E}_T
<ul style="list-style-type: none"> $\tilde{\chi}_i^0 \tilde{\chi}_i^0$ 	$\tilde{\chi}_i^0 \rightarrow \tilde{\chi}_1^0 X, \tilde{\chi}_i^0 \rightarrow \tilde{\chi}_1^0 X'$	$\cancel{E}_T + \text{Dilepton} + (\text{jets}) + (\text{leptons})$
<ul style="list-style-type: none"> $\tilde{t}_1 \tilde{t}_1$ 	$\tilde{t}_1 \rightarrow c\tilde{\chi}_1^0$	2 acollinear jets + \cancel{E}_T
	$\tilde{t}_1 \rightarrow b\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow \tilde{\chi}_1^0 l^\pm \nu, \tilde{\chi}_1^\pm \rightarrow \tilde{\chi}_1^0 q\bar{q}'$	single lepton + $\cancel{E}_T + b's$
	$\tilde{t}_1 \rightarrow b\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow \tilde{\chi}_1^0 l^\pm \nu, \tilde{\chi}_1^\pm \rightarrow \tilde{\chi}_1^0 l^\pm \nu$	Dilepton + $\cancel{E}_T + b's$
<ul style="list-style-type: none"> $\tilde{l}\tilde{l}, \tilde{l}\tilde{\nu}, \tilde{\nu}\tilde{u}\tilde{\nu}$ 	$\tilde{l}^\pm \rightarrow l^\pm \tilde{\chi}_i^0, \tilde{l}^\pm \rightarrow \nu l \tilde{\chi}_i^\pm$	Dilepton + \cancel{E}_T
	$\tilde{\nu} \rightarrow \nu \tilde{\chi}_1^0$	Single lepton + $\cancel{E}_T + (\text{jets})$
		\cancel{E}_T

Exclusive Searches

Inclusive Searches



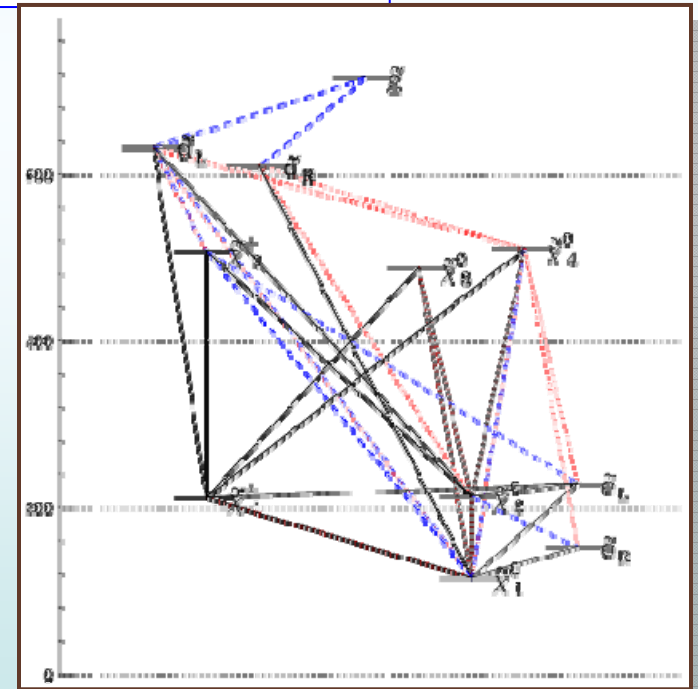
Complex long decay chains to undetected $\tilde{\chi}_1^0$

Inclusive search:

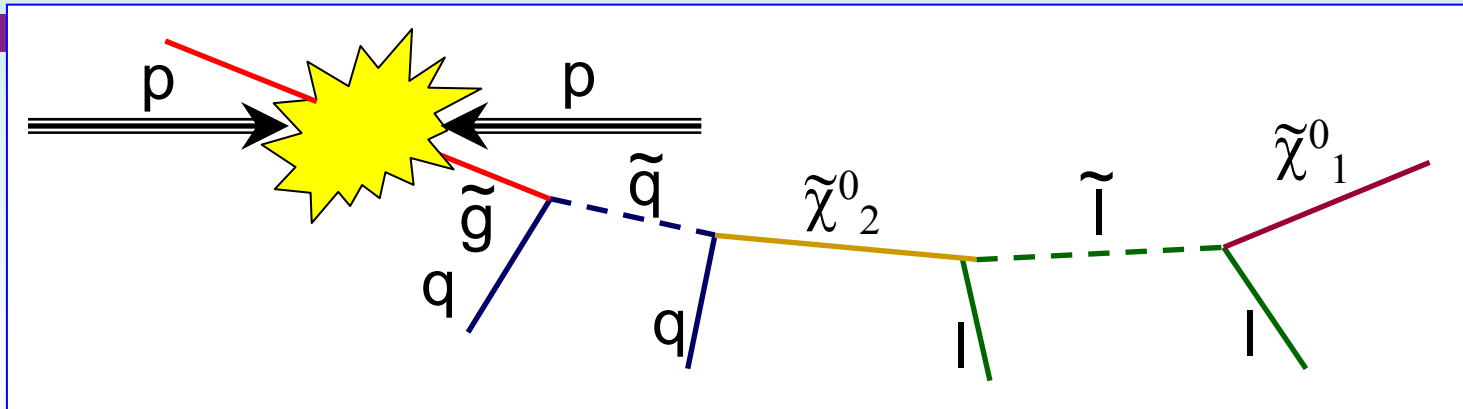
high multiplicity of high- p_T jets

large E_T^{miss} (from escaping LSP)

≥ 0 (high- p_T) leptons



Inclusive Searches



Complex long decay chains to undetected $\tilde{\chi}_1^0$

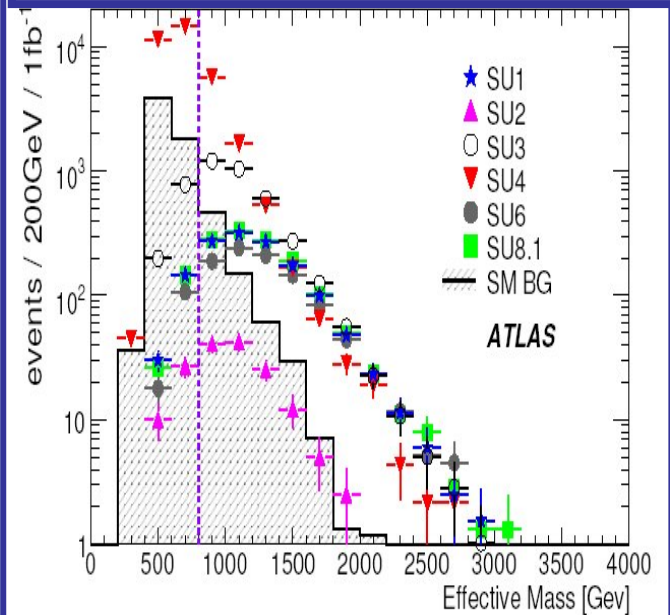
Inclusive search:

high multiplicity of high- p_T jets

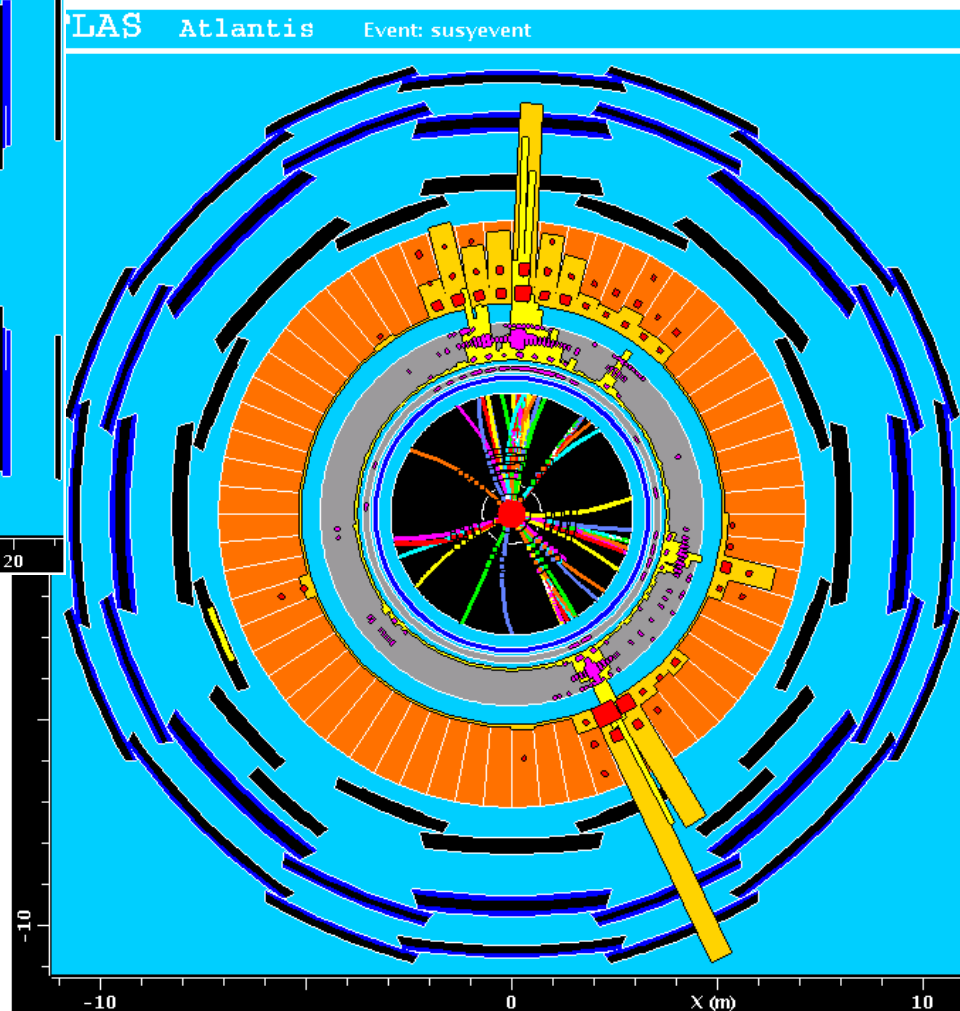
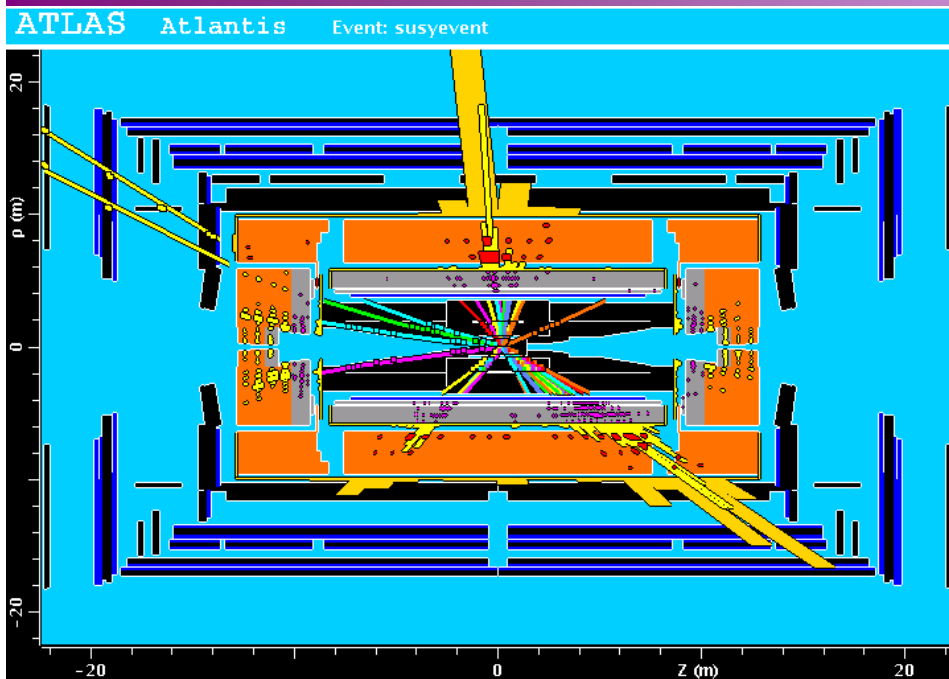
large E_T^{miss} (from escaping LSP)

≥ 0 (high- p_T) leptons

$$M_{\text{eff}} = E_T^{\text{miss}} + \sum_{\text{jet}=1}^4 E_T^{(\text{jet})} (+ p_T(\text{lep}))$$



A (simulated!) SUSY event in ATLAS

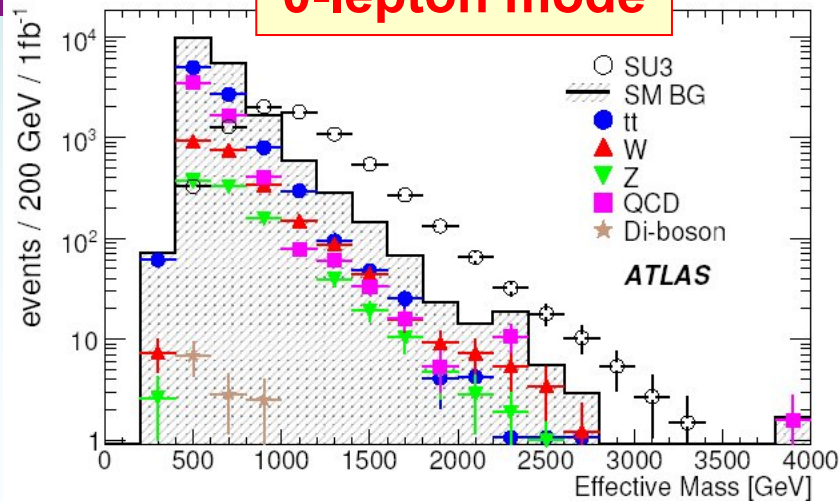


Multi-jet event in
Bulk Region

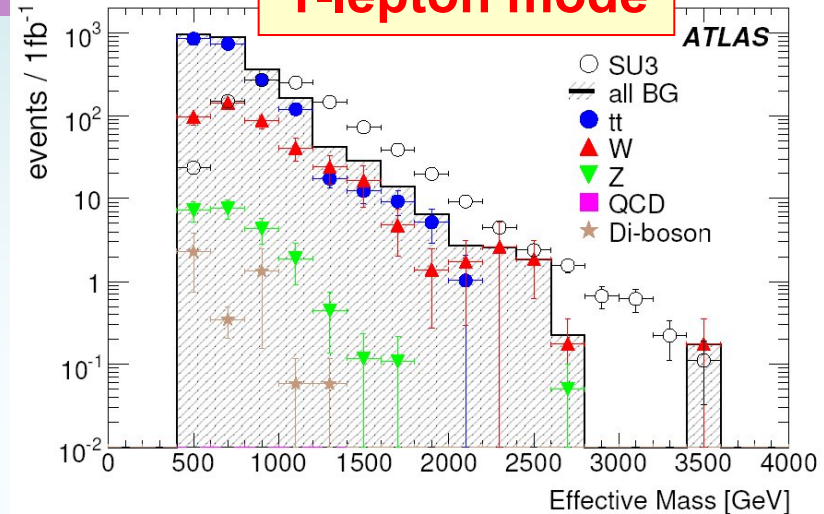
6 jets
2 high-pt muons
Large missing E_T

Standard Model Backgrounds

0-lepton mode



1-lepton mode



Dominant SM backgrounds:

multi-jet QCD

W/Z+jets

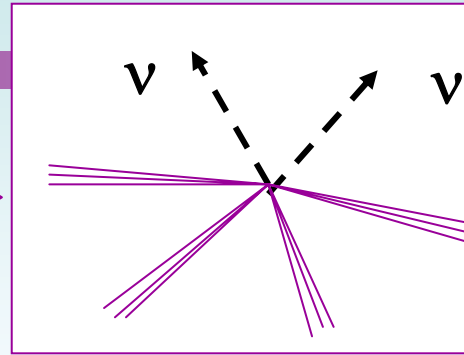
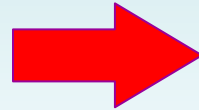
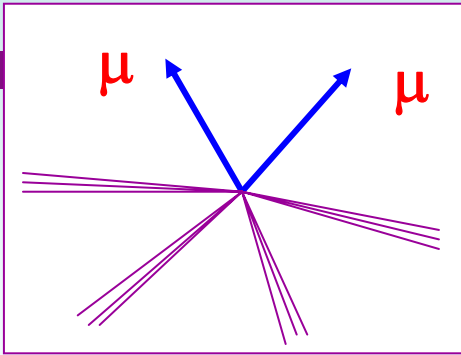
ttbar

dibosons

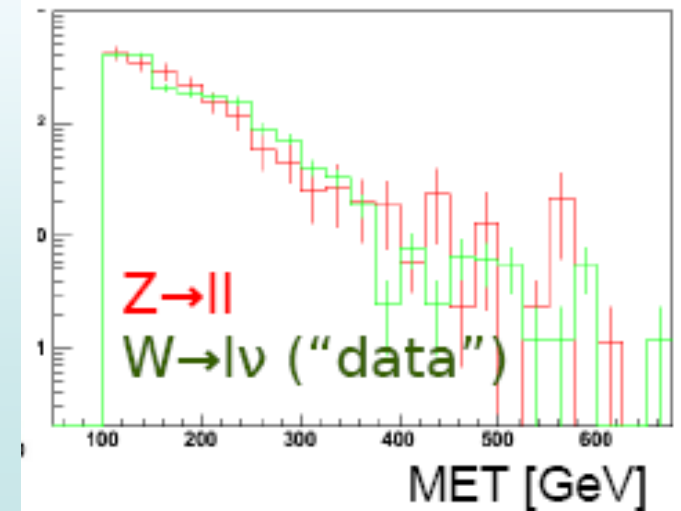
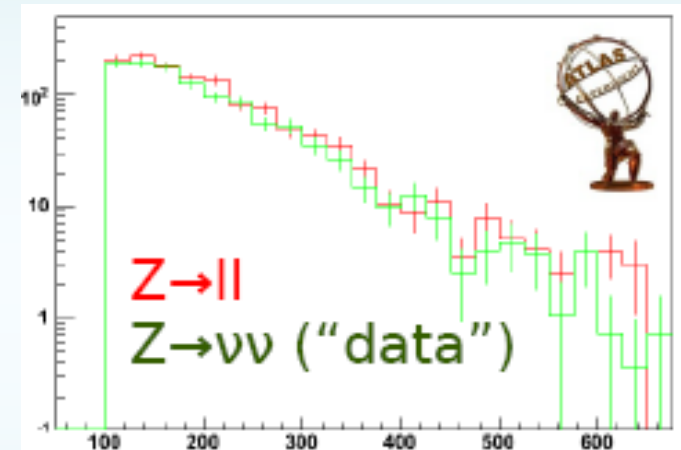
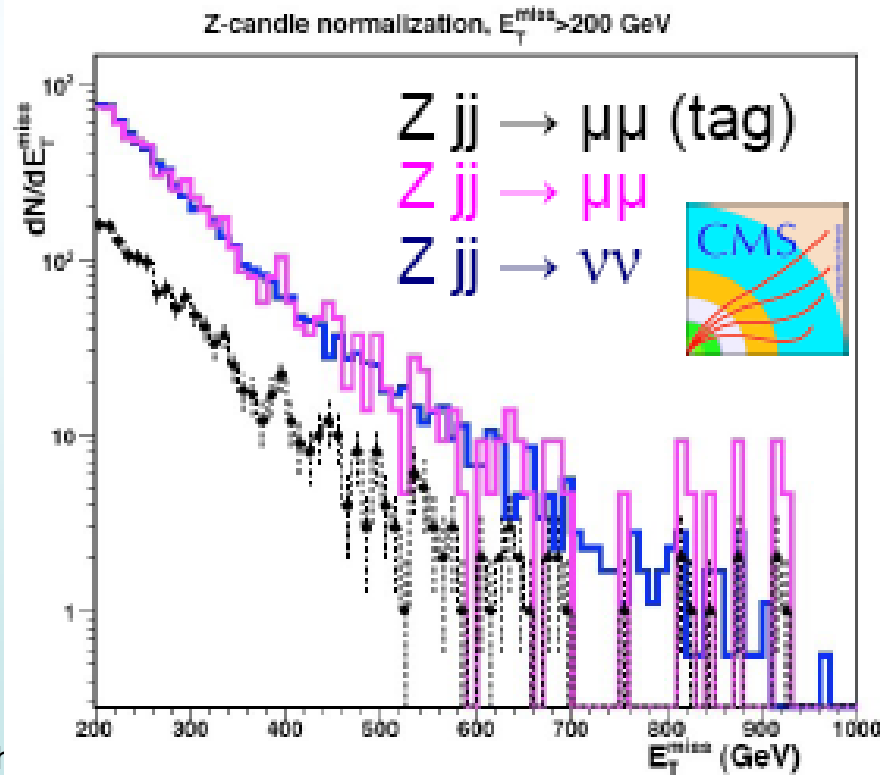
Need robust, data-driven background predictions (as well as predictions from reliable Monte Carlo simulations)

Before we can claim a discovery, we must achieve **complete confidence** in our understanding of the **detectors**, the **trigger**, the **reconstruction** and **particle identification algorithms** and, above all, of the **SM backgrounds**

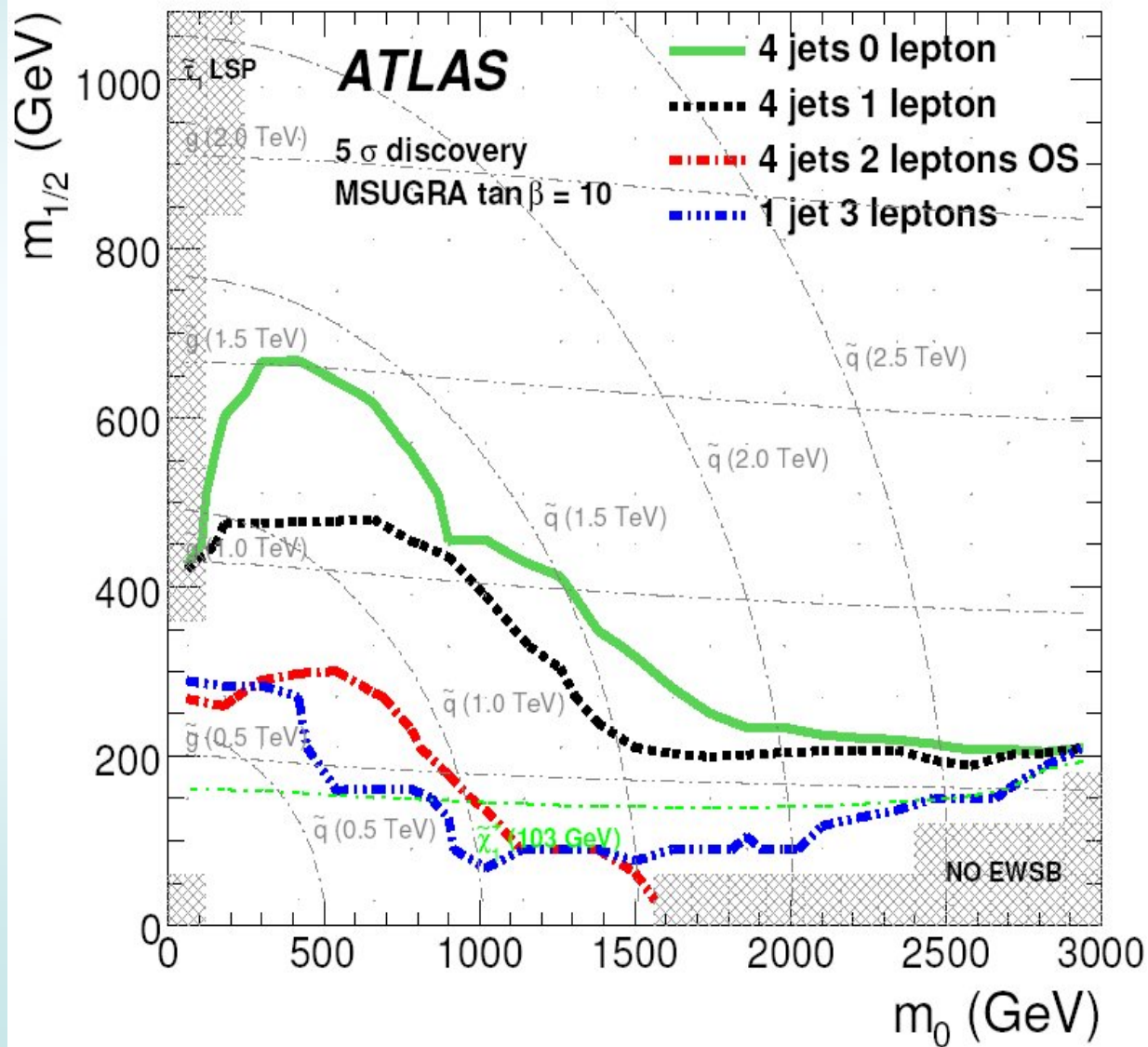
Data-Driven Background Estimation – An Example



Charged lepton replacement in $(Z \rightarrow ll) + \text{jets}$ events

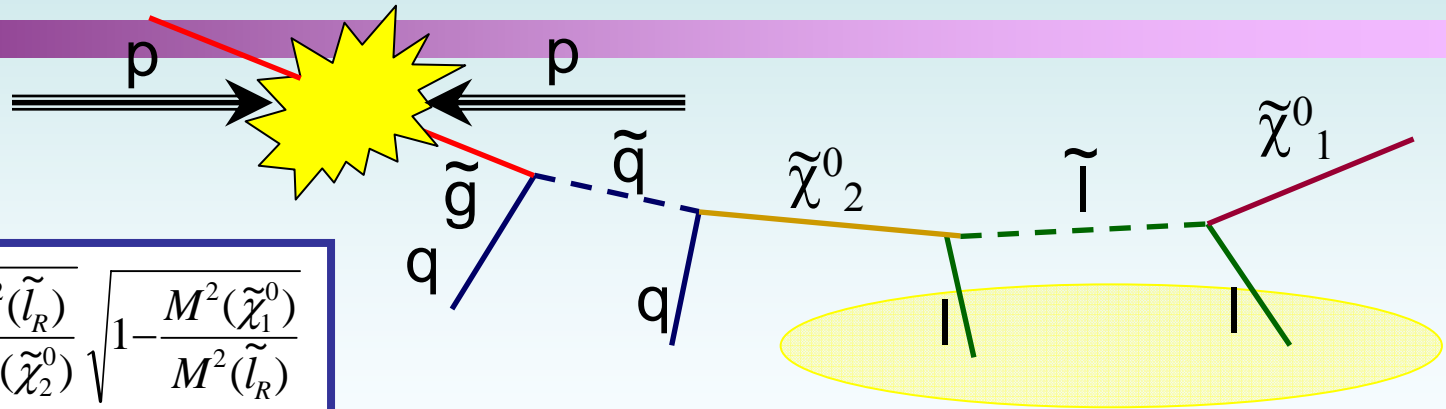


Discovery potential with 1 fb^{-1}



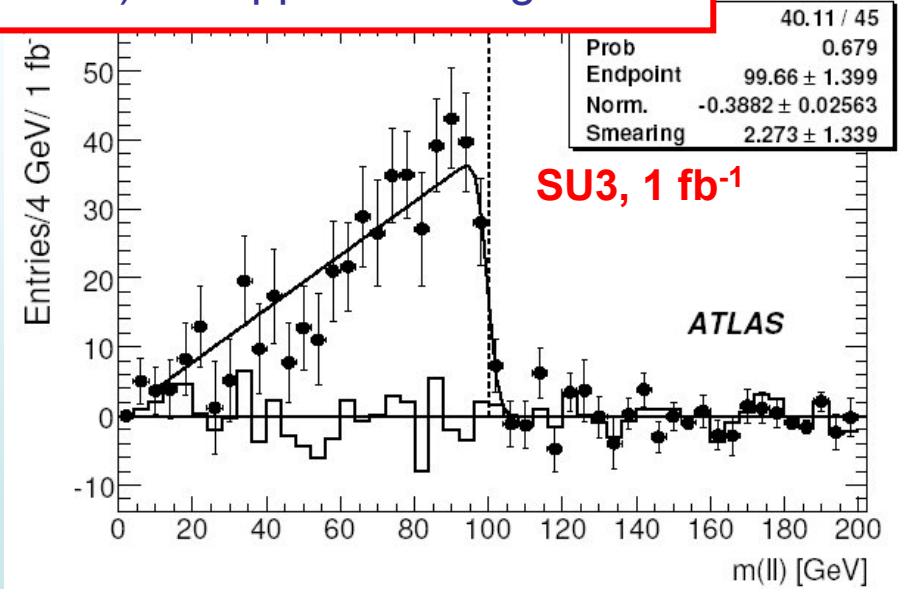
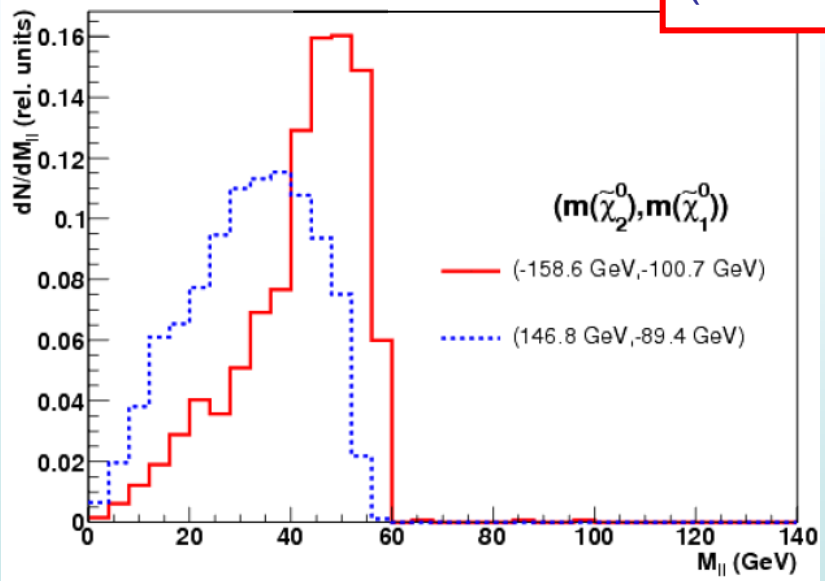
Discovery potential largely dominated by cross-sections and backgrounds

Constraining SUSY – Dilepton Edges

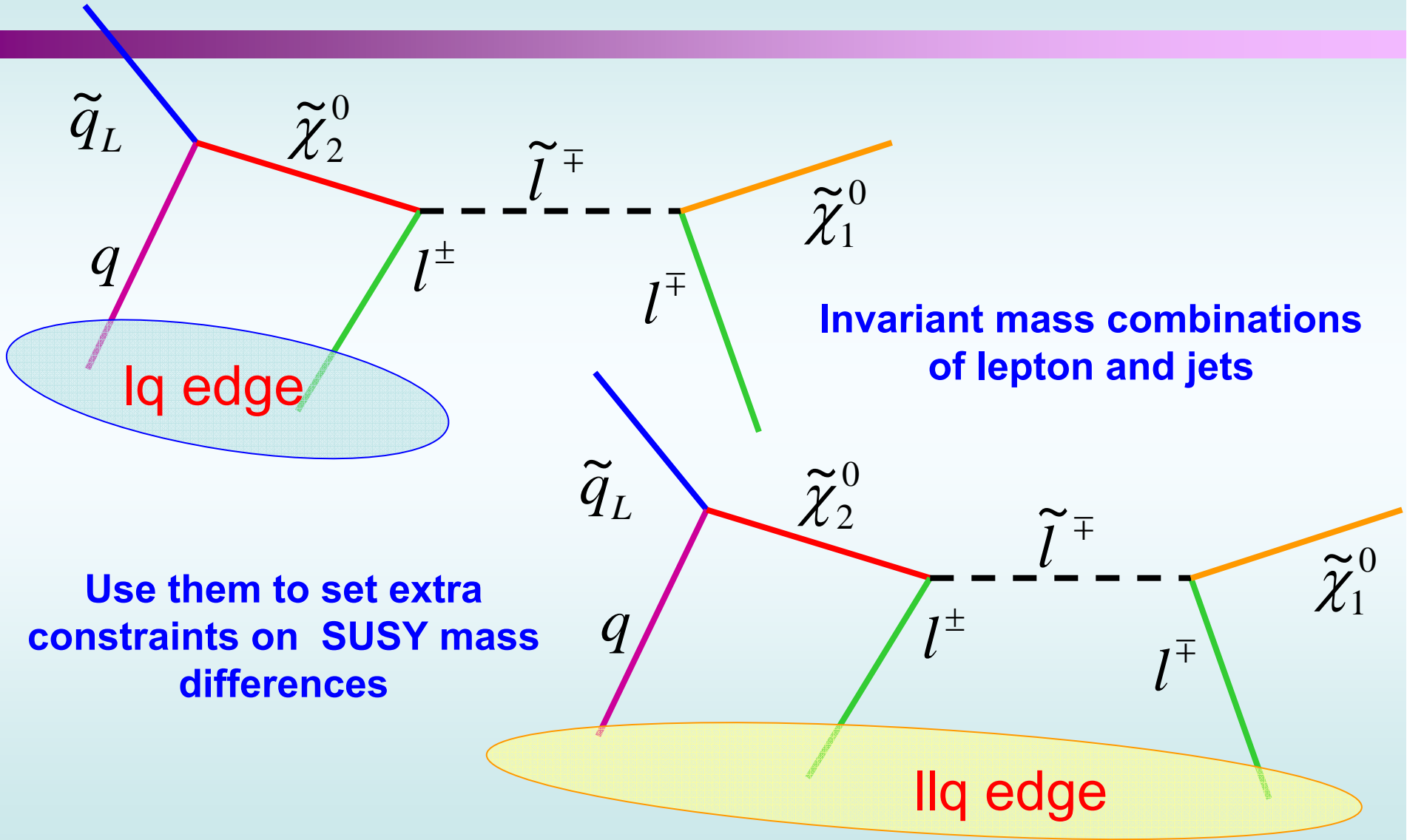


$$M_{ll}^{\max} = M(\tilde{\chi}_2^0) \sqrt{1 - \frac{M^2(\tilde{l}_R)}{M^2(\tilde{\chi}_2^0)}} \sqrt{1 - \frac{M^2(\tilde{\chi}_1^0)}{M^2(\tilde{l}_R)}}$$

Use flavour-subtracted combination (SFOS-OFOS) to suppress backgrounds



More Edges

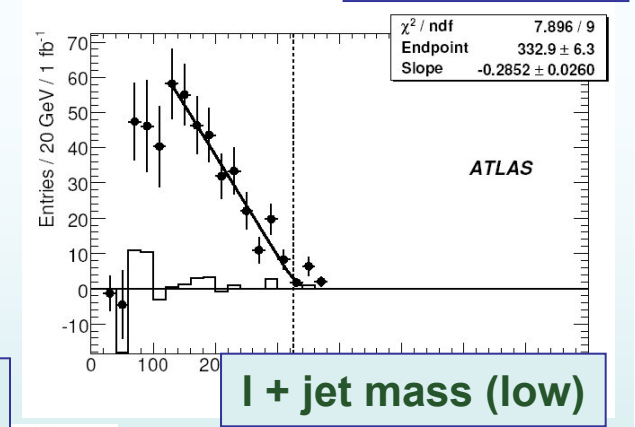
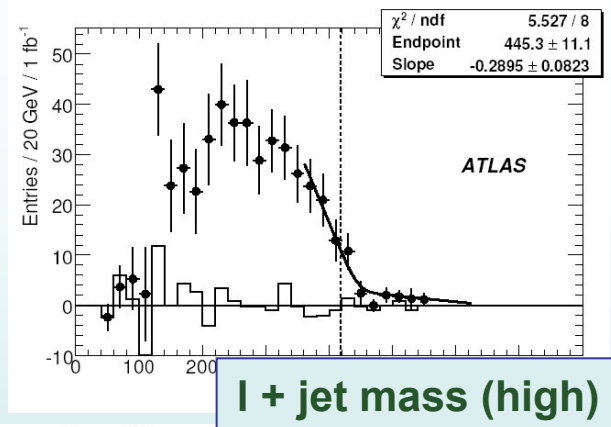
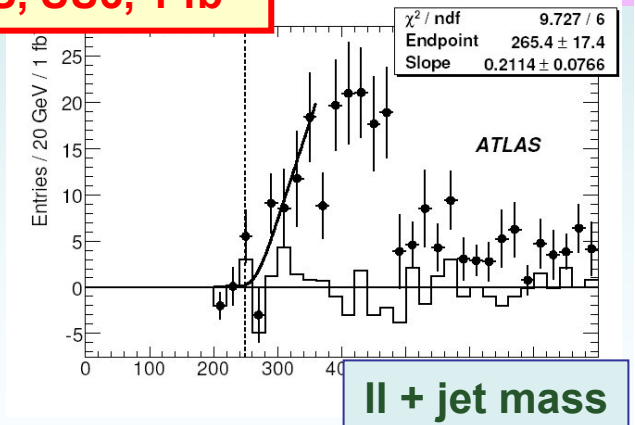
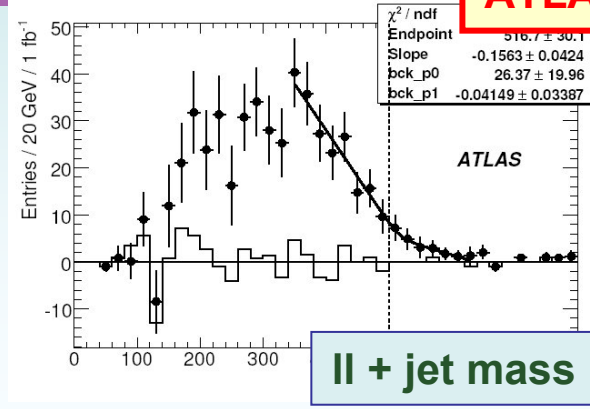


Fits to Mass Edges

ATLAS, SU3, 1 fb⁻¹

Already with O(1 fb⁻¹), the fits to these edges can give reasonably good estimates of SUSY mass differences, and can be used to constrain SUSY

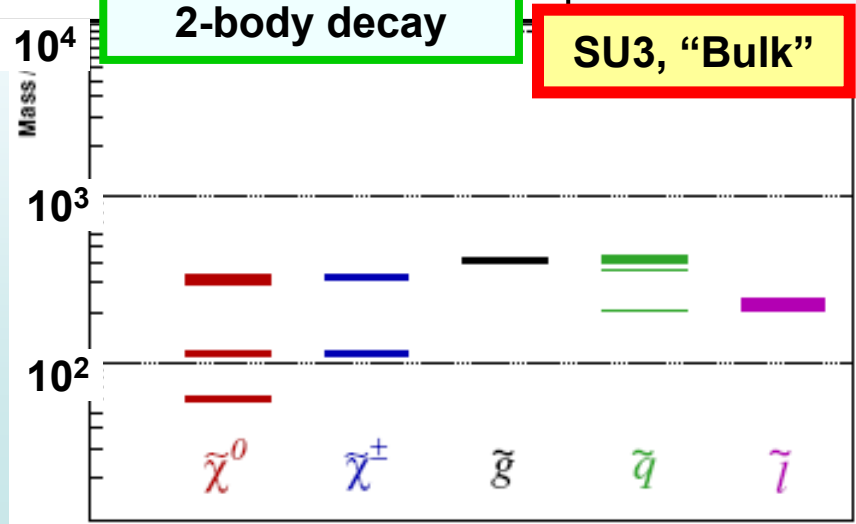
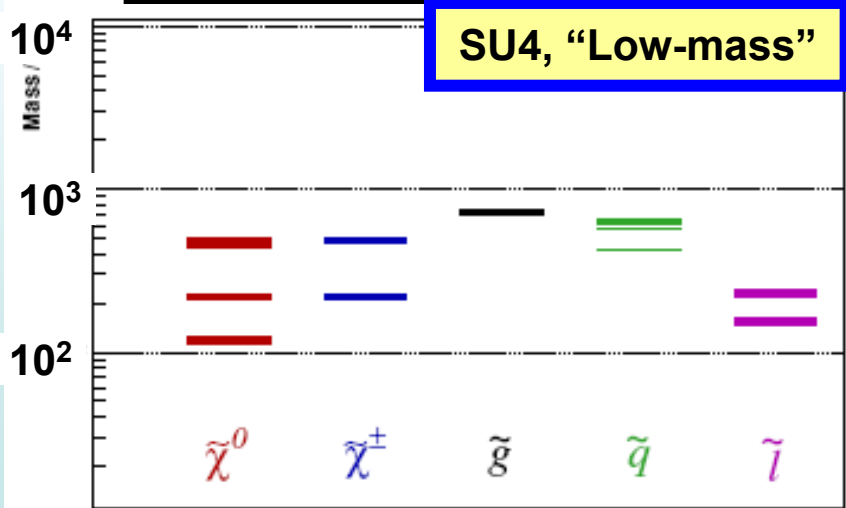
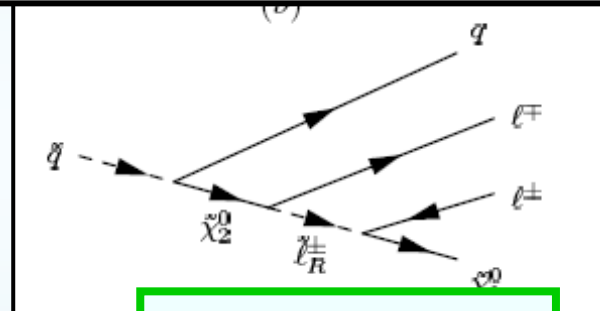
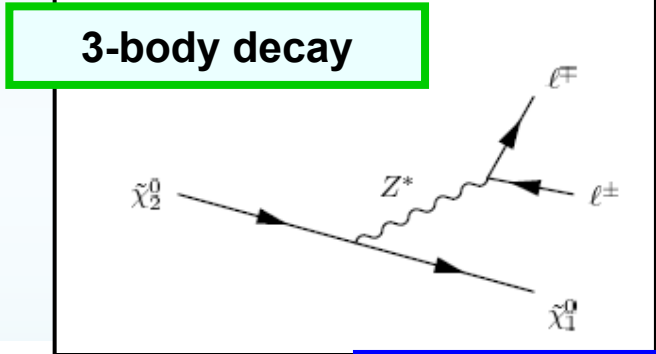
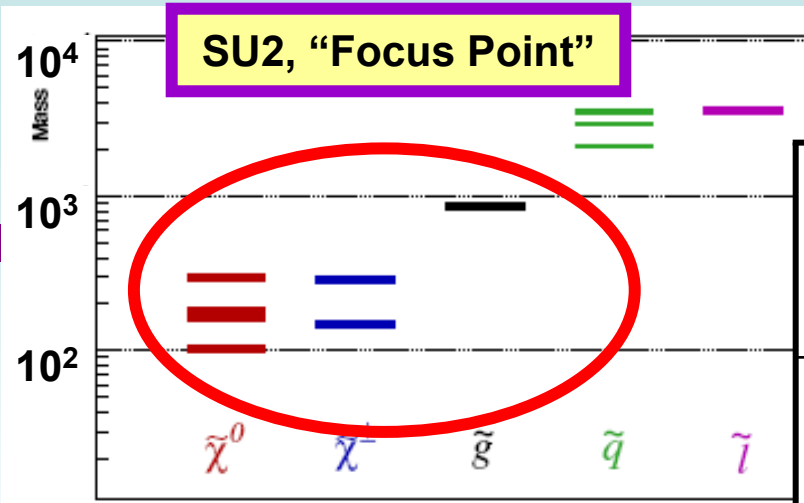
Under the assumption that the considered decay chain is the one actually realised in data



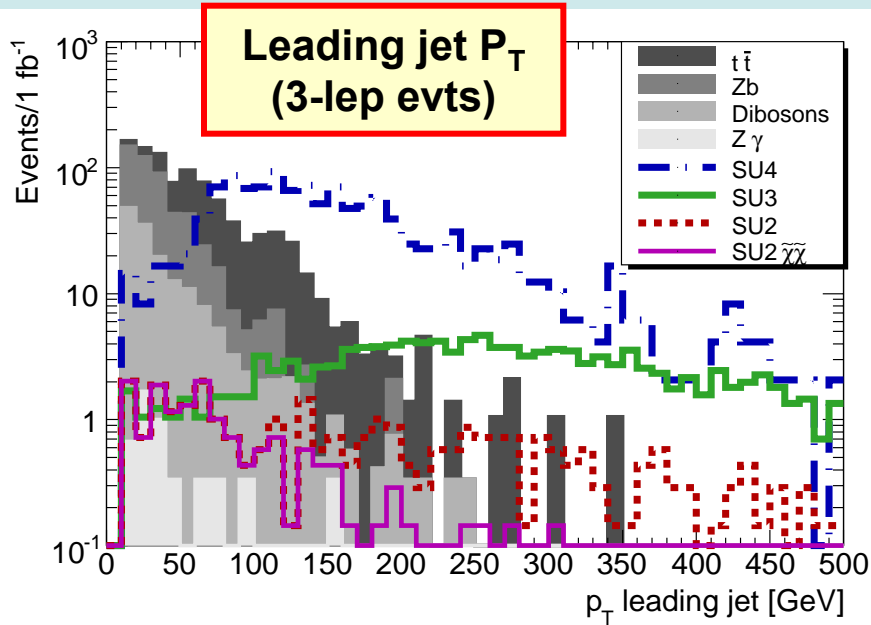
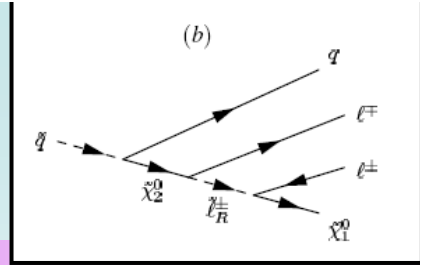
Endpoint	SU3 truth	SU3 measured	SU4 truth	SU4 measured
$m_{\ell\ell q}^{\max}$	501	$517 \pm 30 \pm 10 \pm 13$	340	$343 \pm 12 \pm 3 \pm 9$
$m_{\ell\ell q}^{\min}$	249	$265 \pm 17 \pm 15 \pm 7$	168	$161 \pm 36 \pm 20 \pm 4$
$m_{lq(\text{low})}^{\max}$	325	$333 \pm 6 \pm 6 \pm 8$	240	$201 \pm 9 \pm 3 \pm 5$
$m_{lq(\text{high})}^{\max}$	418	$445 \pm 11 \pm 11 \pm 11$	340	$320 \pm 8 \pm 3 \pm 8$

Trileptons in ATLAS

Point	mSUGRA Parameters					Cross Sections (pb)	
	m_0	$m_{1/2}$	A_0	$\tan\beta$	$\text{sgn}(\mu)$	Tot.	3-lep
SU2	3550	300	0	10	+	7.2	0.07
SU3	100	300	-300	6	+	27.7	0.30
SU4	200	160	-400	10	+	402.2	2.5



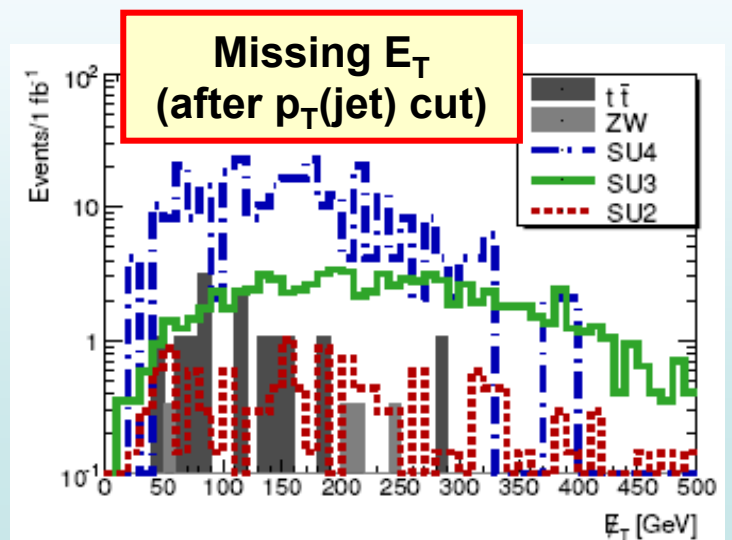
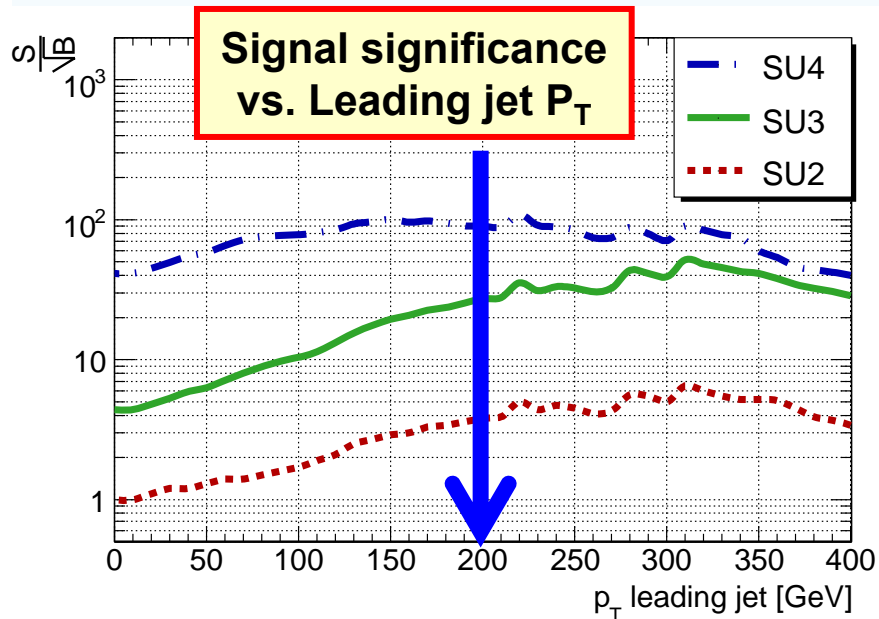
Inclusive Trilepton Analysis – I



Crucially simple analysis flow:

- 3 isolated leptons (e, μ) (taus are next step!)
- At least 1 high-pt jet

Don't need large missing E_T cut – good !!

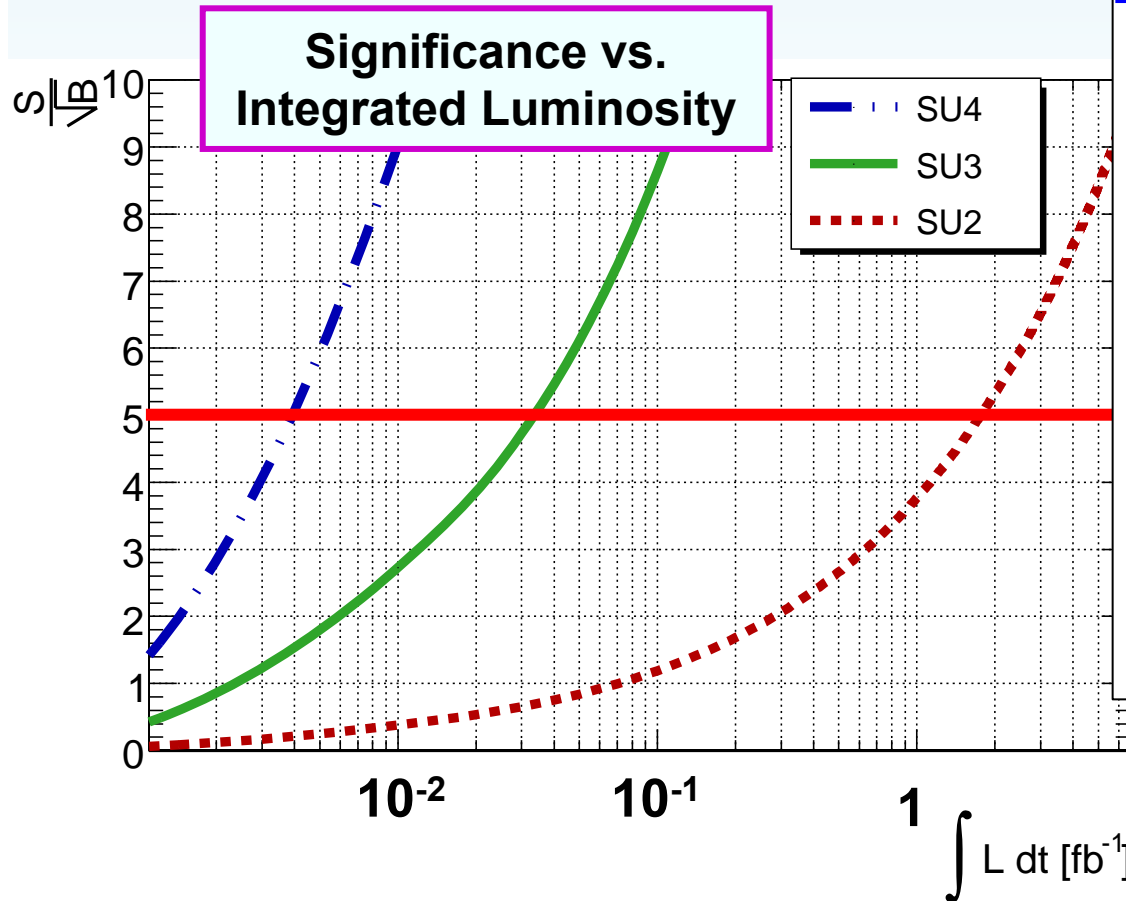


Inclusive Trilepton Analysis

A .De Santo et al., Trilepton signatures at ATLAS, ATL-PHYS-INT-2008-037

A. De Santo et al., Expected performance of the ATLAS experiment, CERN-OPEN-2008-020

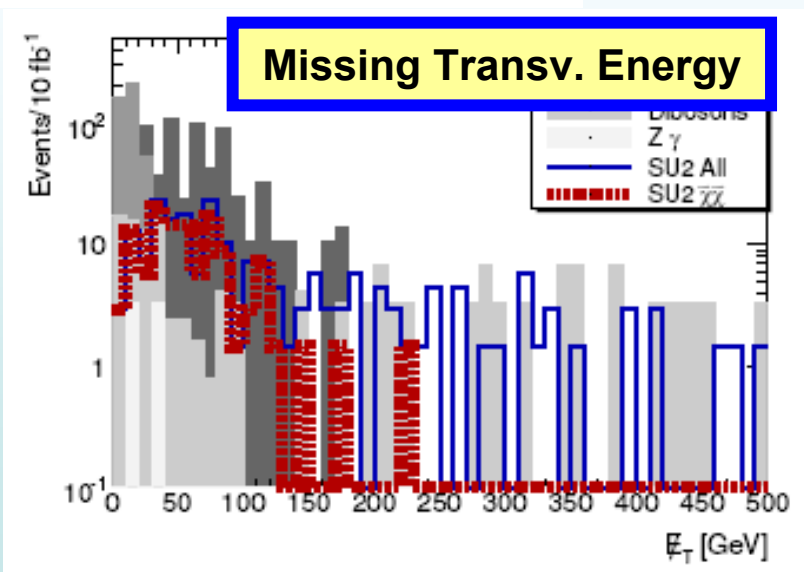
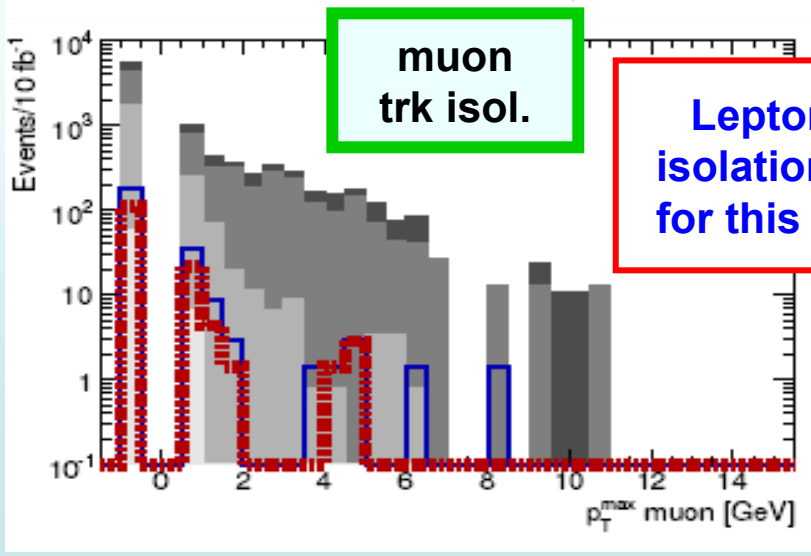
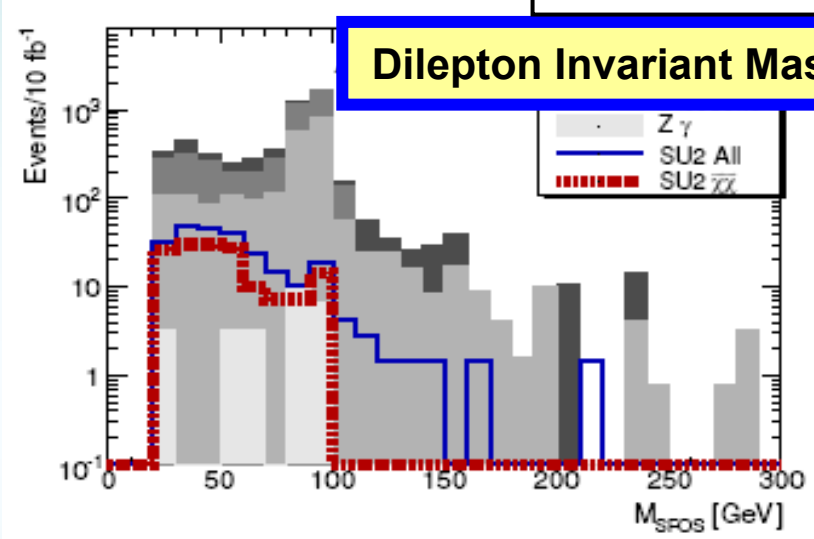
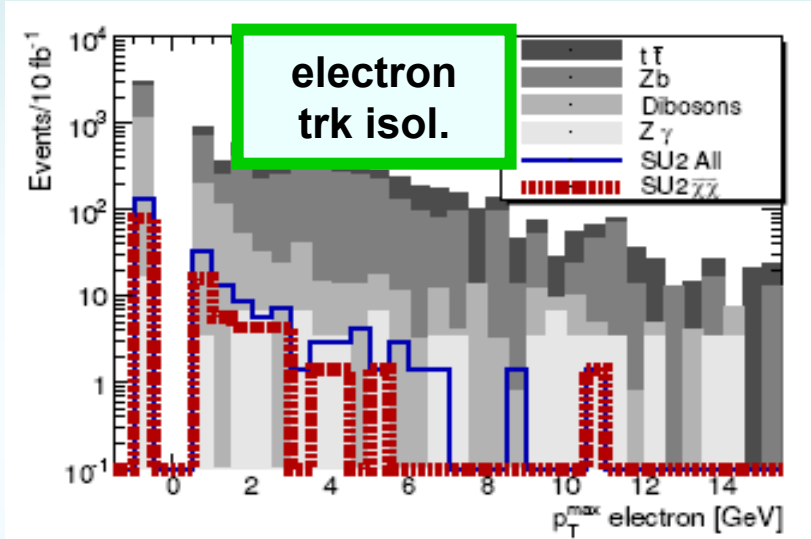
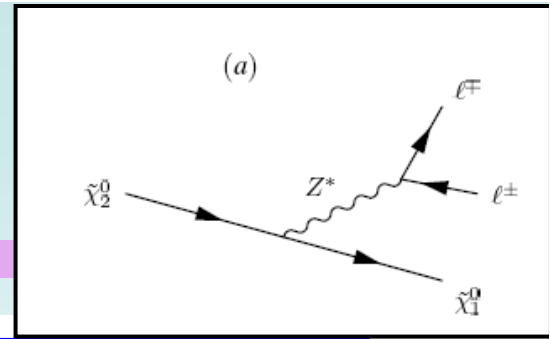
Potentially an early SUSY discovery channel!!



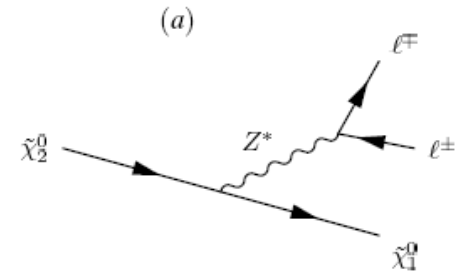
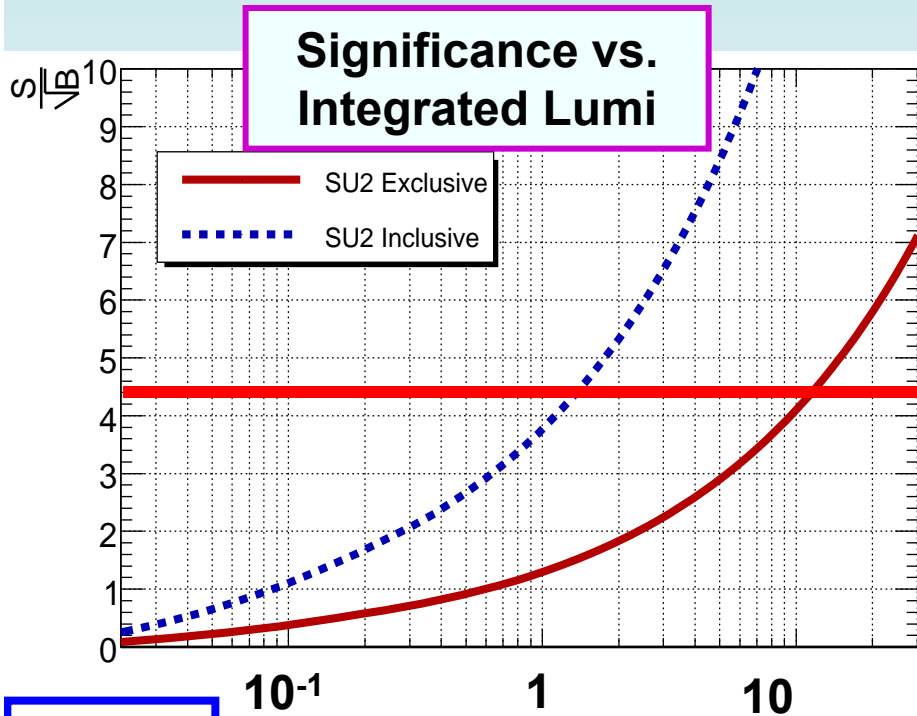
1 fb⁻¹

	None	N_ℓ	p_T^{jet1}
SU2	7111.7	35.0	13.0
SU3	27304.3	139.3	94.3
SU4	396445.5	1283.9	311.7
$t\bar{t}$	440657.9	444.0	10.6
Zb	159115.7	661.6	0.0
ZW	15672.0	192.7	1.3
ZZ	3820.2	58.9	0.0
WW	40051.7	3.3	0.0
Z γ	3283.2	9.4	0.0
$\frac{S}{\sqrt{B}}$	SU2	0.9	3.8
	SU3	3.8	27.3
	SU4	34.7	90.3
$\frac{S}{B}$	SU2	0.0	1.1
	SU3	0.1	7.9
	SU4	0.9	26.2

Exclusive 3-lep Analysis at Focus Point



Exclusive 3-lep Analysis at Focus Point



Not an early physics channel
 But potentially the only way to see SUSY if both scalars and gluino are too heavy!

A. De Santo et al., ditto

10 fb⁻¹

	No cut	N_ℓ	SFOS	Track Isol.	Z-window	\cancel{E}_T	Jet Veto	IP_N	
SU2 $\tilde{\chi}\tilde{\chi}$	64036.6	186.4	177.7	153.2	119.9	98.3	86.7	80.9	
SU2 non- $\tilde{\chi}\tilde{\chi}$	7080.5	163.3	127.2	95.4	85.3	83.8	0.0	0.0	
$t\bar{t}$	4406578.8	4440.2	2812.2	634.3	507.5	475.7	327.7	179.7	
Zb	1591156.7	6616.1	6562.8	2422.8	386.0	0.0	0.0	0.0	
ZW	156719.6	1926.8	1910.1	1682.2	321.7	217.8	214.5	204.4	
ZZ	38201.8	589.4	579.9	475.8	56.8	13.4	11.8	11.0	
WW	400516.9	32.7	24.5	8.2	8.2	8.2	8.2	0.0	
Z γ	32832.3	93.9	90.6	26.8	6.7	3.4	3.4	3.4	
(i)	$\tilde{\chi}\tilde{\chi} = S$	S/\sqrt{B}	1.6	1.6	2.1	3.2	3.5	3.6	4.1
	SM+non- $\tilde{\chi}\tilde{\chi} = B$	S/B	0.0	0.0	0.0	0.1	0.1	0.2	0.2
(ii)	All SU2 = S	S/\sqrt{B}	3.0	2.8	3.4	5.7	6.8	3.6	4.1
	SM = B	S/B	0.0	0.0	0.0	0.2	0.3	0.2	0.2

CONCLUSIONS

Conclusions

With the LHC turn on, new physics at the TeV scale is becoming accessible experimentally for the first time at an accelerator

SUSY provides a well motivated extension to the Standard Model

Very excitingly, we might have the possibility to observe “dark matter in a laboratory”

LET THE SHOW BEGIN !!