

As in  $S$  we have

$$e^{i\omega(t+(x/c)\cos\alpha+(y/c)\sin\alpha)}, \quad \text{where } \vec{k} = -\frac{\omega}{c}(\cos\alpha, \sin\alpha, 0),$$

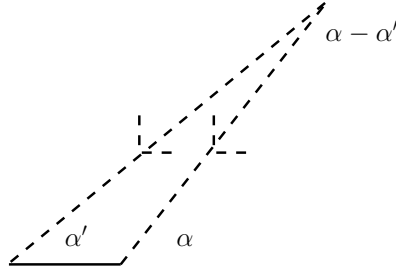
and so

$$\begin{aligned} \tan(\alpha - \alpha') &= \frac{\tan\alpha - \tan\alpha'}{1 + \tan\alpha \tan\alpha'} \\ &= \sin\alpha \frac{\gamma(\cos\alpha + v/c) - \cos\alpha}{\gamma \cos\alpha(\cos\alpha + v/c) + \sin^2\alpha} \\ &= \frac{v}{c} \sin\alpha \frac{1 + v/2c \cos\alpha + O(v^2/c^2)}{1 + v/c \cos\alpha + O(v^2/c^2)} \end{aligned}$$

or

$$\boxed{\tan(\alpha - \alpha') = \frac{v}{c} \sin\alpha \left[ 1 - \frac{v}{2c} \cos\alpha + O\left(\frac{v^2}{c^2}\right) \right]}.$$

Non-relativistic (Bradley)



As  $\tan(\alpha - \alpha') \approx \alpha - \alpha'$  then

$$\alpha - \alpha' = \frac{v}{c} \sin\alpha.$$

## 7.4 Minkowski (Space–Time) diagrams

Portray  $x, y, z, t$  as a point in a four-dimensional space-time:  $(ct, \vec{r}) \equiv (ct, x, y, z)$ .

- Every point  $P$  represents an *event* in space-time
- Under a Lorentz transformation for  $S \rightarrow S'$

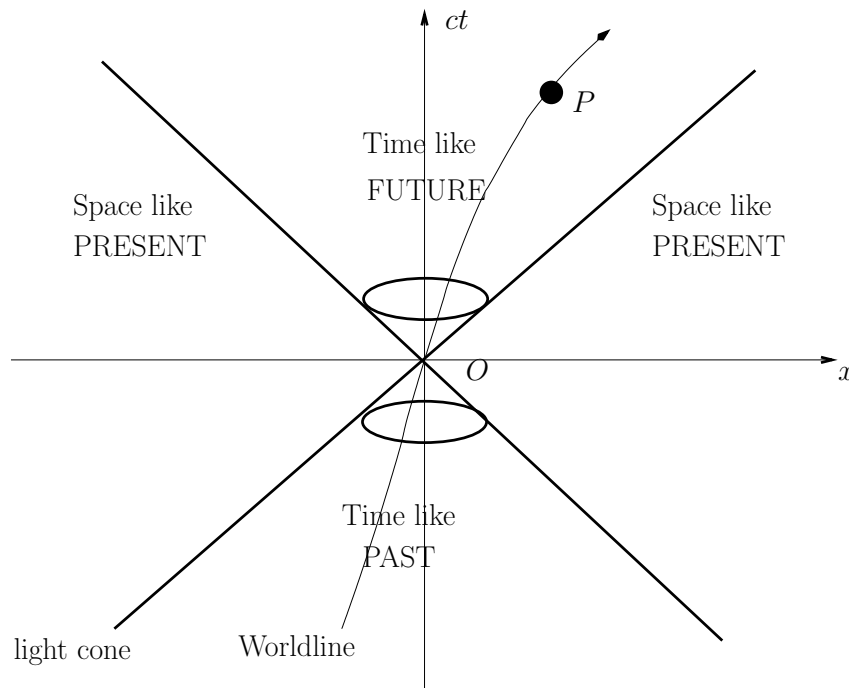
$$s^2 = c^2 t^2 - \vec{r}^2, \quad \vec{r}^2 = x^2 + y^2 + z^2,$$

is a Lorentz invariant, ie  $s^2 = s'^2$  (after a few lines of algebra). So

$$s^2 = \begin{cases} > 0 & \text{time-like} \\ = 0 & \text{light-like} \\ < 0 & \text{space-like} \end{cases}$$

This decomposition remains the same in all inertial frames (ie after any Lorentz transformation)

Minkowski or Space-time picture (wlog two dimensional)



- Light signal define a **cone** (at  $45^\circ$  as choose  $ct$  for  $y$ -axis)
- All time-like points lie *within* light-cone
- All space-like points lie *outside* light-cone
- World-lines are motion of a particle (starting at  $t = 0, x = 0$ ). World-lines of a photon lie *on* the light cone

Consider two space-time points  $P_1(ct_1, \vec{r}_1)$  and  $P_2(ct_2, \vec{r}_2)$ . Then

$$s_{12}^2 = c^2(t_1 - t_2)^2 - |\vec{r}_1 - \vec{r}_2|^2,$$

also a Lorentz invariant.

### Space-like separation

From  $s_{12}^2 < 0$ , then  $x_1 - x_2 > c(t_1 - t_2)$ , ie events not connected by a light signal

#### no causal connection

Can find a Lorentz transformation to  $S'$  where both events  $P_1$  and  $P_2$  are at the same time, as

$$c(t'_1 - t'_2) = \gamma \left( c(t_1 - t_2) - \frac{v}{c}(x_1 - x_2) \right).$$

Choose

$$v = c \frac{c(t_1 - t_2)}{x_1 - x_2} < c, \quad \text{so } S' \text{ possible}$$

Hence can always find Lorentz transformations so that the order of space-like events is interchanged.

### Time-like separation

From  $s_{12}^2 > 0$ , then  $c(t_1 - t_2) > x_1 - x_2$ , ie events connected by a light signal

#### causal connection possible

Cannot find a Lorentz transformation to  $S'$  where both events  $P_1$  and  $P_2$  are at the same time, as would need

$$v = c \frac{c(t_1 - t_2)}{x_1 - x_2} > c, \quad \text{so } S' \text{ not possible}$$

So cannot interchange cause and effect. However as

$$x'_1 - x'_2 = \gamma ((x_1 - x_2) - v(t_1 - t_2)).$$

can find a frame  $S'$  where  $x'_1 = x'_2$ , events happen at the same place

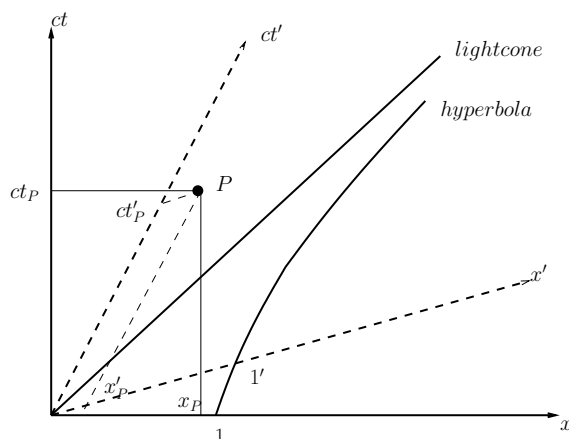
Define **proper time** by

$$\boxed{c^2 \tau^2 = s^2}, \quad \text{or} \quad \tau = \sqrt{t^2 - \frac{\vec{r}^2}{c^2}}.$$

So here we have

$$\tau_{P_1} - \tau_{P_2} = \sqrt{(t_1 - t_2)^2 - \frac{|\vec{r}_1 - \vec{r}_2|^2}{c^2}}.$$

ie time measured on a clock moving with uniform speed  $|\vec{v} = c\vec{r}_1 - \vec{r}_2|/(t_1 - t_2)$  between the two events,  $P_1$  and  $P_2$ .



### 7.4.1 Diagrammatic relation between $S$ and $S'$

We now briefly consider the relationship between  $S$  and  $S'$  on a Minkowski diagram

Co-ordinate systems same at  $t = 0 = t'$

- $S'$  time-axis defined by  $x' = 0 = \gamma(x - vt)$  or

$$ct = \frac{1}{v/c} x,$$

ie straight line, gradient  $> 1$ , ie between  $S$  time-axis and light signal

- $S'$  space-axis defined by  $t' = 0 = \gamma(t - v/c^2 x)$  or

$$ct = v/c x,$$

ie straight line, gradient  $< 1$ , ie between  $S$  space-axis and light signal

#### Scaling of axes

The Lorentz transformation  $S \rightarrow S'$  changes the scale of the axes.

Choose a length unit by  $\bar{s}^2 = -1$ . Then

$$x^2 = (ct)^2 + 1,$$

– ie **hyperbola**, all points on it have length  $\bar{s}^2 = -1$ . In particular  $t = 0$  cuts  $x$ -axis at  $x = 1$ .

But as  $s^2$  is a Lorentz invariant, then this is equivalent to

$$x'^2 = (ct')^2 + 1,$$