Particle Physics

Dr Victoria Martin, Spring Semester 2013 Lecture 14: Symmetries



*Symmetries of QED and QCD
*Parity, Charge Conjugation and Time Reversal
*Parity Violation in Weak Decay
*CP and CPT

Moriond Conference



 I was webcast: <u>http://webcast.in2p3.fr/videos-</u> searches for the beh boson into fermions at atlas

Guest Seminars

• From Edinburgh University researchers on their work.

- Next Monday (18th March) in tutorial:
 - Guest seminar from Dr Greig Cowan on B-physics at LHCb
- Following Monday (25th March) in tutorial
 - Guest seminar from Dr Wahid Bhimji on Higgs physicst at ATLAS

Cabibbo-Kobayashi-Maskawa Matrix

- Mass eigenstates and weak eigenstates of quarks are not identical.
- Weak eigenstates are admixture of mass eigenstates, conventionally described using CKM matrix a mixture of the down-type quarks:

weak
eigenstates
$$\begin{pmatrix} d' \\ s' \\ b' \end{pmatrix} = \begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} \begin{pmatrix} d \\ s \\ b \end{pmatrix}$$
 mass
eigenstates

- The CKM matrix is unitary, $V_{\rm CKM}^{\dagger}V_{\rm CKM} = 1$
- Wolfenstein parameterisation

$$V_{\rm CKM} = \begin{pmatrix} 1 - \lambda^2/2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \lambda^2/2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{pmatrix} + \mathcal{O}(\lambda^4)$$

CKM Fit



 \bullet Results from BaBar, Belle and LHCb experiments used to find best values for η and ρ parameters in Wolfenstein parameterisation.

Introduction: Symmetries in Particle Physics

- Symmetries play a central role in particle physics.
- Symmetries describe operations which leave a physical system unchanged.
 - Noether's theorem: every symmetry corresponds to a conservation law.
 - Conservation laws can be experimentally verified
- We use the group theory to describe the symmetries of the three forces:
 - SU(3) symmetry of QCD \Rightarrow conservation of colour charge
 - SU(2) symmetry of weak force \Rightarrow conservation of weak isospin
 - U(1) symmetry of QED \Rightarrow conservation of electric charge
 - Translations in time and space; rotations in space ⇒ conservation of angular and fourmomentum
- Some symmetries are observed to only hold under certain conditions, e.g. conservation of quark flavour for strong and electromagnetic interactions
- Breaking of symmetries can be either "dynamical" or "spontaneous" (e.g. Higgs mechanism).

Parity (P)

- Parity, **P**, is a spatial inversion through the origin: $P\psi(\vec{r}) = \psi(-\vec{r})$ (not a mirror reflection in a plane!)
- If you act *P* twice on a state you get the original state back $\Rightarrow P^2 = 1$.
 - Eigenvalues of P are either +1 (even) or -1 (odd)
 - e.g. $\psi(x) = \sin kx$ is odd $\psi(x) = \cos kx$ is even
 - e.g. Hydrogen atom states described by Legendre polynomials: $Y_L^m = P_L^m(\cos \theta) e^{im\phi}$
 - Parity of state described by Y_L^m is $(-1)^L$ (depends on orbital angular momentum L)
 - S (L=0) and D (L=2) states are even, P (L=1) states are odd

Intrinsic Parity

- Intrinsic parity is property of the elementary particles, we **define**:
 - Fermions to be eigenstates with P = +1 (even)
 - Antifermions to be eignestates w. P = -1 (odd)

• Parity operator (**P**) can be represented by $\gamma^0 = \begin{pmatrix} \mathbf{I} & \mathbf{0} \\ \mathbf{0} & -\mathbf{I} \end{pmatrix}$

• Check this by looking at Dirac spinors (on problem sheet)

- Photons have odd parity, P = -1
 - Parity operation changes the direction of Electric fields.
 - Check this from transformation of electromagnetic fields (A^{μ})
- In general:
 - Scalar ($\overline{\psi} \psi$) and axial-vector ($\overline{\psi} \gamma^5 \gamma^{\mu} \psi$) quantities have P = +1 (even)
 - Pseudoscalar ($\overline{\psi} \gamma^5 \psi$) and vector ($\overline{\psi} \gamma^{\mu} \psi$) quantities have P = -1 (odd)

parity

right handed

left handed

Charge Conjugation (C)

- A change from particle to antiparticle: $C\bar{f} = f$; $Cf = \bar{f}$
- $C^2 = 1 \Rightarrow$ eigenvalues of C are either +1 (even) or -1 (odd)
 - (*C* can be represented by $i\gamma^2$)
- Fermions and antifermions are not eigenstates of C !



- Photons have C = -1. C changes sign of the electric charges, and therefore of the electromagnetic field. Similarly gluons have C = -1.
- C and P observed to be conserved in electromagnetic and strong interactions.
- Mesons $(q\overline{q})$ have $C = (-1)^{L+S}$ and $P = (-1)^{L+1}$
- Lightest mesons pseudoscalars $(\uparrow\downarrow)$ with $J^{PC} = 0^{-+}$ Second-lightest states are vectors $(\uparrow\uparrow \text{ or }\downarrow\downarrow)$ with $J^{PC}=1^{--}$



• e.g. J/ψ (cc) meson has $J^{PC}=1^{--}$. EM and strong decays via odd number of photons or gluons to conserve C and P

Time Reversal (T)

• A reversal of the arrow of time: $T\psi(t) = \psi(-t)$

- $T^2 = 1$ so eigenvalues of T are either +1 (even) or -1 (odd)
- *T* corresponds to interchanging the initial and final states in a Feynman diagram.
- *T* reverses the sign of time-dependent variables including momentum and angular momentum (spin)
- It is evident that the Universe is not *T* symmetric!
 - Big Bang, Hubble expansion, increase in Entropy ...

Discrete Symmetries of Physical Quantities

Quantity	Notation	P	C	T
Position	\vec{r}	-1	+1	+1
Momentum (Vector)	\vec{p}	-1	+1	-1
Spin (Axial Vector)	$\vec{\sigma} = \vec{r} \times \vec{p}$	+1	+1	-1
Helicity	$\vec{\sigma} \cdot \vec{p}$	-1	+1	+1
Electric Field	$ $ \vec{E} $ $	-1	-1	+1
Magnetic Field	$ $ \vec{B} $ $	+1	-1	-1
Magnetic Dipole Moment	$\vec{\sigma} \cdot \vec{B}$	+1	-1	+1
Electric Dipole Moment	$\vec{\sigma} \cdot \vec{E}$	-1	-1	-1
Transverse Polarization	$\left \vec{\sigma} \cdot (\vec{p_1} \times \vec{p_2}) \right $	+1	+1	-1

Parity Violation in Weak Decays

- First observed by Chien-Shiung Wu in 1957 through β decay of polarised ⁶⁰Co nuclei: ${}^{60}Co \rightarrow {}^{60}Ni + e^- + v_e$
 - Recall: under parity changes momentum sign but not spin $\vec{p} \to -\vec{p}$ $\vec{S} \to \vec{S}$
- Electrons were observed to be emitted to opposite to nuclear spin direction
- Particular direction is space is preferred!
- *P* is found to be violated maximally in weak decays
- *P* is a good symmetry in strong and QED.
- Recall the vertex term for the weak force is $g_w \gamma^\mu (1-\gamma^5)/\sqrt{8}$
- Fermion currents are proportional to $\overline{\psi}(e) (\gamma^{\mu} \gamma^{\mu} \gamma^5) \psi(v_e)$ "vector axial vector"
- Under parity, $P(\overline{\psi}(\gamma^{\mu} \gamma^{\mu}\gamma^{5})\psi) = \overline{\psi}(\gamma^{\mu} + \gamma^{\mu}\gamma^{5})\psi$
- Compare to QED and QCD vertices: **P** ($\overline{\psi} \gamma^{\mu} \psi$) = $\overline{\psi} \gamma^{\mu} \psi$

(pure eigenstates - no *P* violation)

(mixture of eigenstates)



Parity Violation in Pion Decays

- Consider charged pion decay at rest, $\pi^- \rightarrow \mu^- + \overline{\nu_{\mu}}$
- Charged pion has $S = 0 \Rightarrow$ muon and neutrino produced with equal & opposite spin
 - \rightarrow The muon and neutrino will have identical helicities.
 - Experiments observe muon helicity is always right-handed
 - Conclusion: Only right-handed anti-neutrinos exist! (or only right-handed) anti-neutrinos interact)
- Similarly for $\pi^+ \rightarrow \mu^+ + v_{\mu}$ anti-muon helicity is always left-handed only left-handed neutrinos exist!
- This observation also explains why $\pi^- \rightarrow \mu^- + \overline{v}_{\mu}$ is preferred over $\pi^- \rightarrow e^- + \overline{v}_e$
 - The weak interaction only acts on LH chiral states, pion must decay to RH helicity electron or muon. $\frac{m_\ell}{m_\pi + m_\ell}$
 - \rightarrow The change from LH chiral to RH helicity for lepton (ℓ) mass: \propto

 $\overline{\nu}_{\mu}$ \longleftarrow μ^{-}

Neutrino States



CPT Theorem

- *CPT* is the combination of *C*, *P* and *T*. Turns a forward-going particle with LH helicity into backward-going antiparticle with RH helicity.
- CPT Theorem states:

"All interactions described by a local Lorentz invariant gauge theory must be invariant under the combined operation *CPT* "

- CPT violation would imply non-locality and/or loss of Lorentz invariance
- Impossible to write down relativistic quantum field theories?
- Impossible to describe interactions in terms of Feynman diagrams?
- CPT conservation implies that CP violation is equivalent to T violation
- The Universe needs *CP* violation for the matter-antimatter asymmetry and it needs *T* violation for the arrow of time



Tests of CPT Invariance

• CPT invariance implies particles and antiparticles must have equal masses:

$$\frac{M(\mathbf{K}^{0}) - M(\overline{\mathbf{K}}^{0})}{\frac{1}{2}[M(\mathbf{K}^{0}) + M(\overline{\mathbf{K}}^{0})]} < 10^{-18} \text{ GeV}$$

• Particle and antiparticles must have equal lifetimes:

$$\frac{\Gamma(\mathbf{K}^0) - \Gamma(\overline{\mathbf{K}}^0)}{\frac{1}{2}[M(\mathbf{K}^0) + M(\overline{\mathbf{K}}^0)]} < \times 10^{-17}$$



 Particle and antiparticles must have equal and opposite charges and magnetic moments

$$Q(p) + Q(\bar{p}) < 10^{-21} e$$

$$\frac{g(e^+) - g(e^-)}{\frac{1}{2}[g(e^+) + g(e^-)]} < 2 \times 10^{-12}$$

• Hydrogen and anti-hydrogen atoms have identical spectra

Summary

- Symmetries play a key role in describing interactions in particle physics.
- QED and QCD obey Gauge symmetries in the Lagrangian corresponding to symmetry groups. These lead to conservation of electric and colour charge.
- Three important discreet symmetries: Charge Conjugation (C), Parity (P) and Time reversal (T).
 - C: changes the sign of the charge
 - P: spatial inversion, reserves helicity. Fermions have P=+1, antifermions P=-1
 - T: changes the initial and final states
 - Gluons and photons have C = -1, P = -1
- C and P are conserved in QED and QCD, maximally violated in weak
- Only LH neutrinos and RH anti-neutrinos are found in nature.
- CPT is thought to be absolutely conserved.

CKM Fit



- All measurements, including results from BaBar and Belle experiments, penguin diagrams, neutral kaon mixing used to find
- \bullet Best values for η and ρ parameters in Wolfenstein parameterisation.