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\*“Where are we in penguin physics and what are we after “

- (quark)-flavor within SM and beyond: CKM and more
- rare processes: mixing, FCNC decays

# Core questions of flavor physics

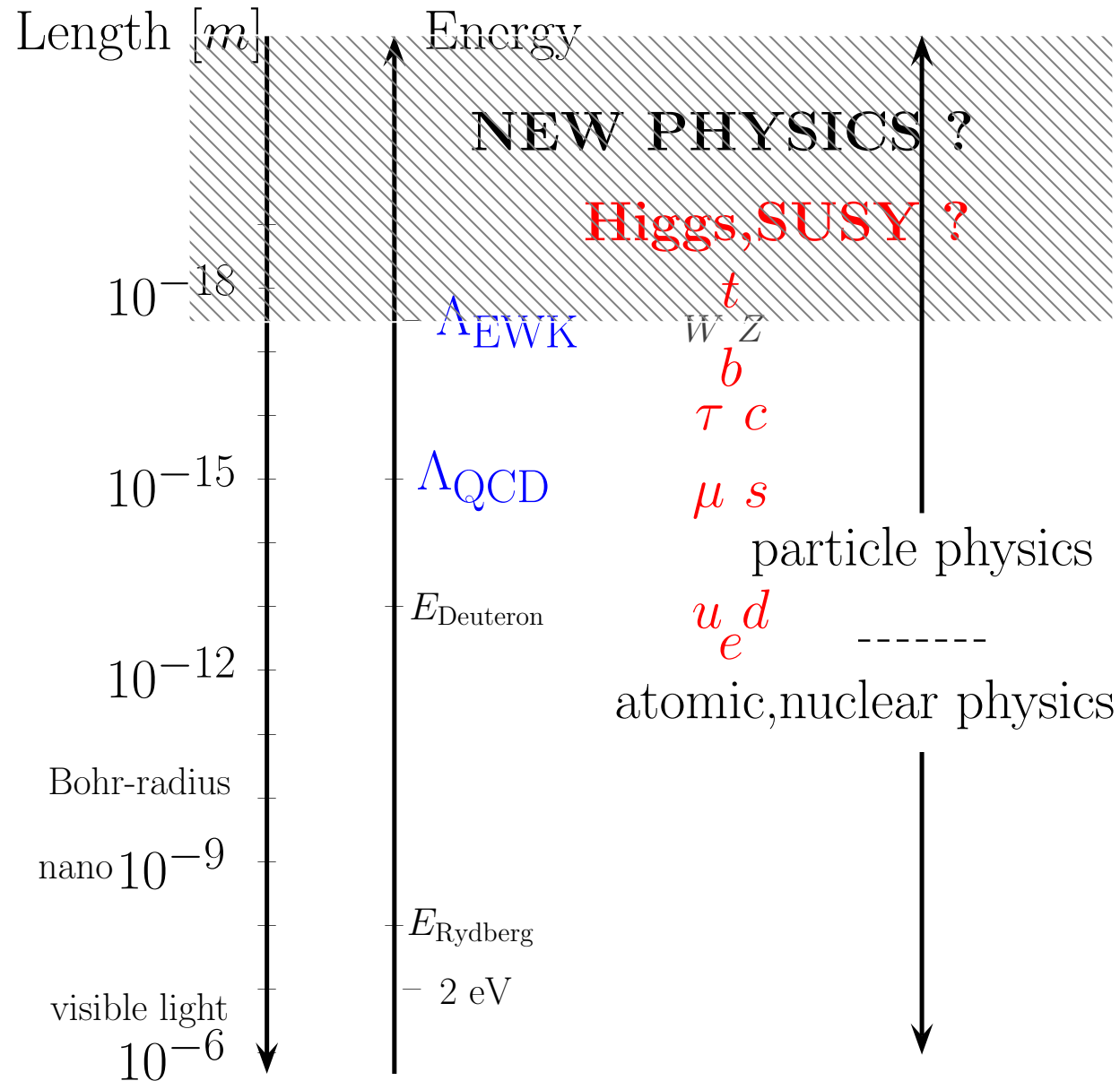
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- CP violation needed for baryogenesis  $(n - \bar{n})/s \simeq 10^{-10}$   
SM not sufficient
- strong CP problem: Why is  $\bar{\Theta} \lesssim 10^{-10}$  and  $\delta_{CKM} = \mathcal{O}(1)$ ?
- origin of flavor; explanation of peculiar masses and mixings
- neutrino masses

... are core questions of the SM, plus: unification, Higgs mass, dark matter, dark energy, gravity

despite its impressive experimental support the SM is rather viewed  
as an effective theory valid up to  $\Lambda \sim \mathcal{O}(m_W)$

# The high energy frontier



flavor and CP in SM:  $-\mathcal{L}_Y = \bar{Q}Y_u h^c U + \bar{Q}Y_d h D + \bar{L}Y_e h E + h.c.$

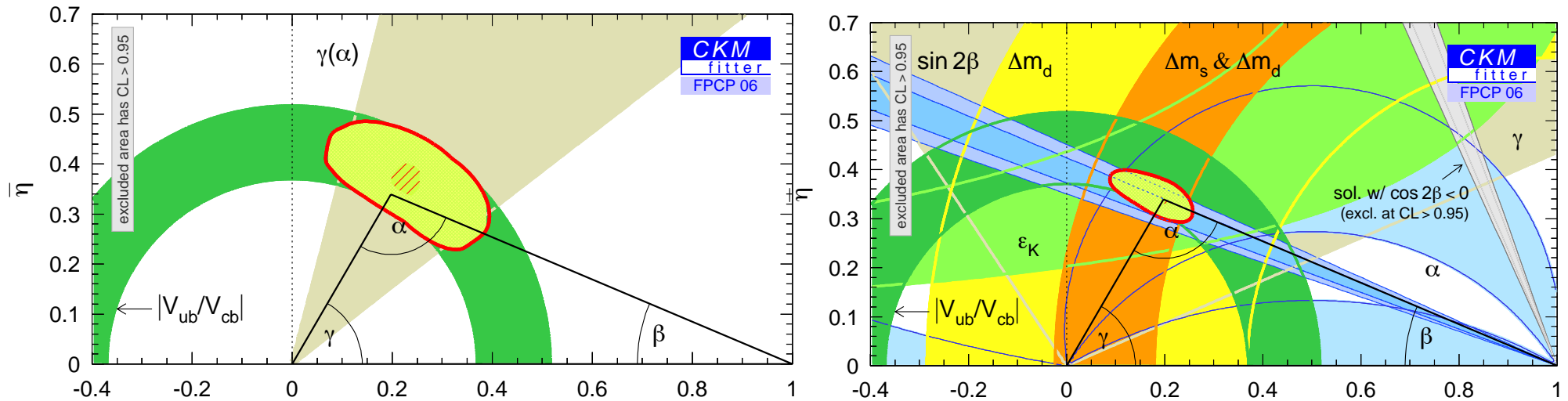
flavor symmetry:  $U(3)^5 \xrightarrow{Y} U(1)_B \times U(1)_L \times U(1)_Y$

quarks:  $Y_{u,d}$  36 real numbers  $\rightarrow$  10 physical parameters: 6 quark masses plus CKM (3 angles, 1 phase)

- determine 10 parameters  $\rightarrow$  PDG, HFAG
- how do the Yukawas look like – and why ? top-down, GUT's, textures, extra dimensions, Froggatt-Nielsen, horizontal symmetries, anarchy, ..
- Is this all to flavor/CP violation, that is, does **MFV** hold ?  
**minimal flavor violation = no further breaking of flavor than through Yukawas**
- are there deviations from the SM in rare processes ?

# CKM/Yukawa CP-violation from tree level, precision

the unitarity triangle  $V_{ub}V_{ud}^* + V_{cb}V_{cd}^* + V_{tb}V_{td}^* = 0$



tree fit  $V_{ub}$  ( $B \rightarrow (\bar{X}_u, \pi)\ell\nu$ ,  $B \rightarrow \tau\nu$ ),  $\gamma$  from  $B \rightarrow DK^{\bar{p}}$

loop input to full fit:  $\epsilon_K$ ,  $\Delta m_{d,s}$ ,  $\sin 2\beta(\bar{c}c)$

SM/MFV-like at least for  $b \rightarrow \bar{c}cs$ ,  $K$ -,  $B_d$ -mixing,  $\Delta m_d/\Delta m_s$

CKM=precision input within SM  $\epsilon(\alpha) = 5\%$ ,  $\epsilon(\beta) = 4\%$ ,  $\epsilon(\gamma) = 8\%$

# SM tests with indirect processes

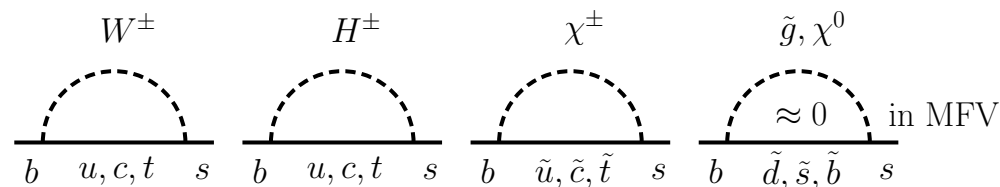
indirect loop processes:

$\Lambda \gtrsim m_W$  scale of New Physics

$$\mathcal{L}_{eff} = \sum_i c_i^{(n)} \frac{O_i^{(n)}}{\Lambda^n} \quad c_i^{(n)} \leftrightarrow f\left(\underbrace{m_j, g_l, \dots}_{colliders}; \underbrace{\varphi_{CKM}, \varphi_m, \delta_n}_{flavorphysics}\right)$$

no competition from large SM tree contributions

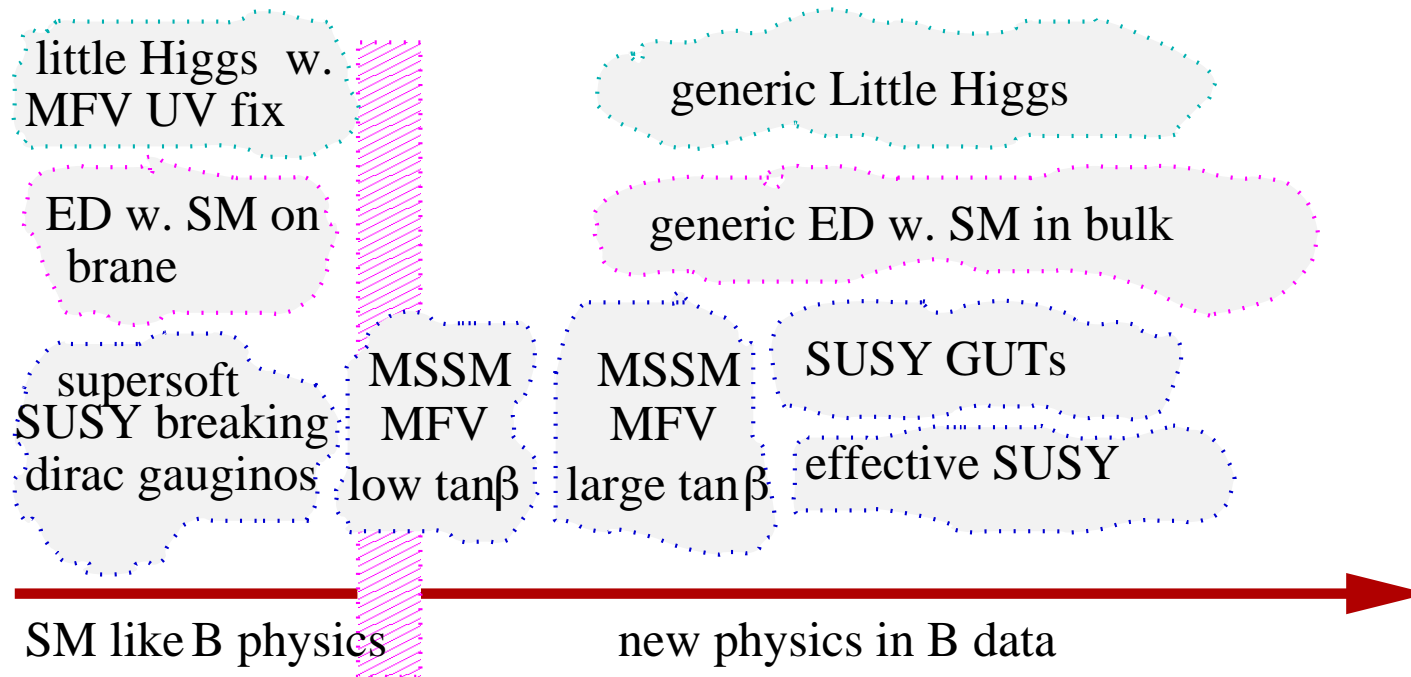
FCNC: sensitivity to SM, NP phases  $\varphi$ , flavor-breaking couplings  $\delta$



BSM-FCNC: offset rates, distort spectra, induce CP-asymmetries, V+A currents, ..

models of EWKSB with NP @ TeV

Fig from hep-ph/0207121



reach in indirect signals depends on beyond the SM flavor/CP violation (minimal=CKM), large parameters such as  $\tan\beta$  and theoretical and experimental uncertainties

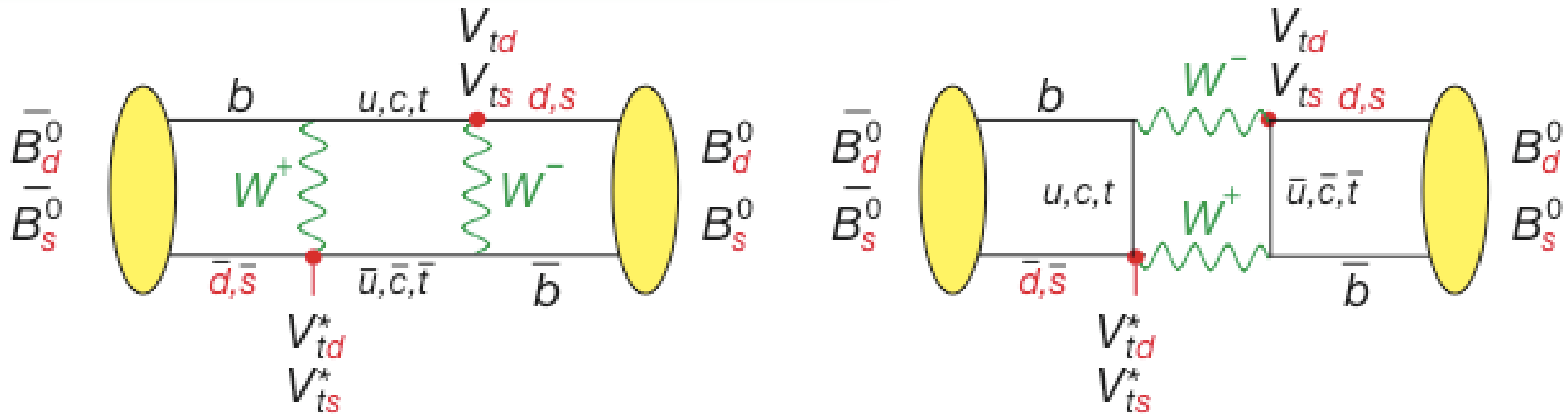
today's and future impact of

- $B_s - \bar{B}_s$  mixing (new Tevatron data from Moriond '06)
- $b \rightarrow sll$  modes



SM:  $W - (u, c, t)$ -box  $\Delta m_s^{\text{SM}} \sim (V_{tb}^* V_{ts})^2 \frac{g^4}{16\pi^2} \frac{(\bar{b}\Gamma_s)(\bar{b}\Gamma'_s)}{m_W^2}$

top dominated;  $V_{tb} \simeq 1$ ,  $V_{ts} = -A\lambda^2$  independent of CKM-triangle-fit



$\Delta m_d^{\text{world ave}} = 0.507 \pm 0.004 \text{ ps}^{-1}$  (HFAG '05)

$\Delta m_s^{\text{SM}} = 18.3_{-1.5}^{+6.5} \text{ ps}^{-1}$  (CKMfitter/EPS'05)  $\Delta m_s^{\text{SM}} = 20.0 \pm 1.8 \text{ ps}^{-1}$  (UTfit hep-ph/0501199)

DØ:  $17 < \Delta m_s < 21 \text{ ps}^{-1}$  at 90% C.L. hep-ex/0603029

CDF:  $\Delta m_s = 17.31_{-0.18}^{+0.33} \pm 0.07 \text{ ps}^{-1}$  hep-ex/0606027

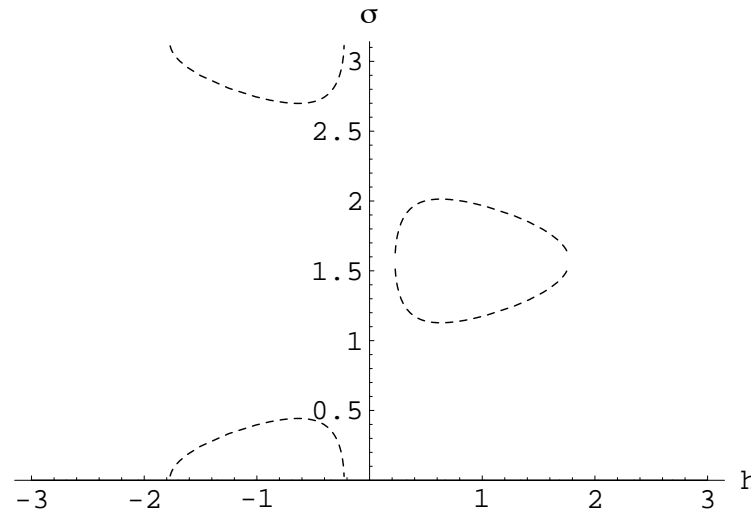
# generic NP constraints from lower bound

assume in plots:  $\Delta m_s^{\text{SM}} = \Delta m_s^{\text{data}} = 19.0 \pm 2 \text{ ps}^{-1}$  @ 90 % C.L.

mixing with NP:  $\Delta m_s = \Delta m_s^{\text{SM}} \cdot |1 + h e^{2i\sigma}|$  see also hep-ph/0509117 for nicer plots

only lower bound near SM  $\Delta m_s > 16.6 \text{ ps}^{-1}$  @ 95 % C.L. (HFAG '05)

$\Delta m_s > \Delta m_s^{\text{SM}}(1 - \epsilon)$ :



dotted:  $\epsilon = 0.23$

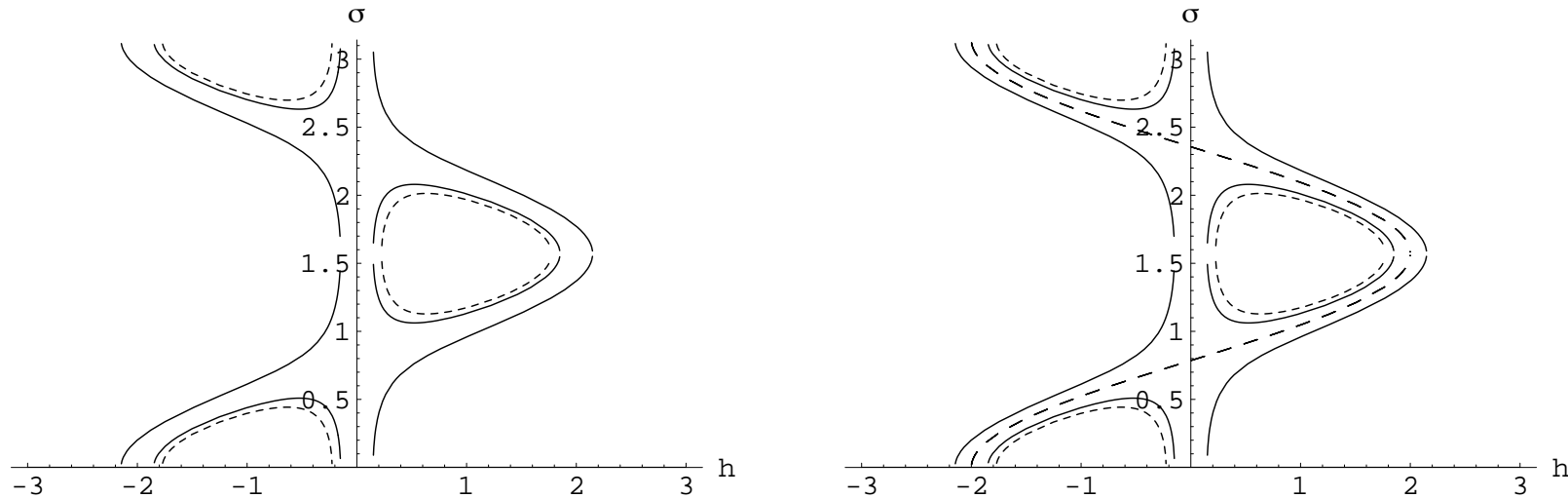
gaps excluded around  $h \simeq +1, \sigma \simeq \pi/2$  and  $h \simeq -1$  with  $\sigma \simeq 0, \pi$

size of NP-amplitude  $h$  unconstrained

# generic NP constraints: measurement

$$\Delta m_s = \Delta m_s^{\text{SM}} \cdot |1 + h e^{2i\sigma}|; \text{ with } \Delta m_s \simeq \Delta m_s^{\text{SM}} (1 \pm \epsilon)$$

$$\Delta m_s^{\text{data}} = 19.0 \pm 2 \text{ ps}^{-1} \text{ @ 90 \% C.L.}$$



dotted: lower bound only, solid:  $\epsilon = 0.15$ , dashed (and  $h = 0$ ):  $\epsilon = 0$

size of NP-amplitude  $|h| < 2 + \epsilon$  limited

$\mathcal{O}(1)$  NP-amplitude possible if NP-phase  $\sigma$  cooperates

large  $|h|$  is fine-tuned to some degree for small errors  $\epsilon \rightarrow 0$

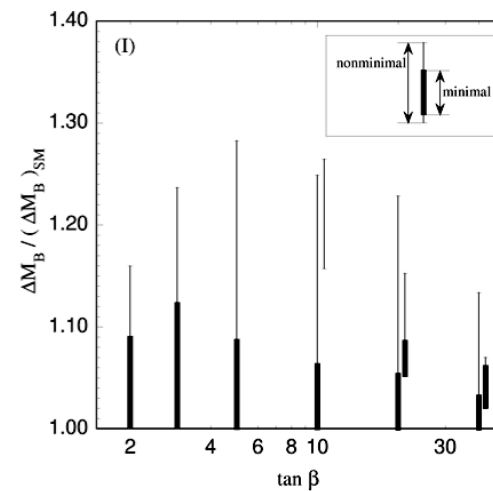
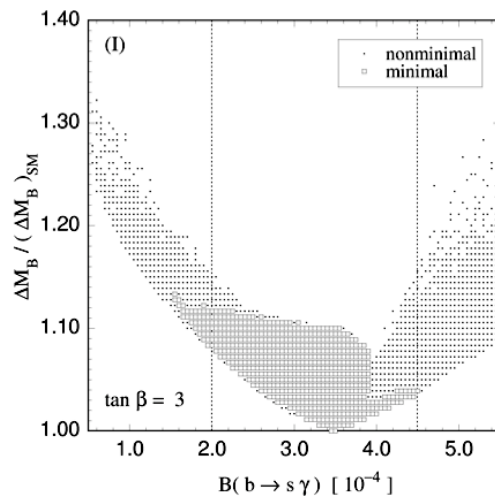
I. low  $\tan \beta$ : CKM-ology: hep-ph/9903535, SUGRA: hep-ph/9908499

$h = h_{H^\pm} + h_{\chi^\pm}$ ;  $\sigma = 0$  (no BSM CP-violation)

$h > 0$  in whole parameter space  $h$  equal for  $B_d$  and  $B_s$

ok with  $b \rightarrow s\gamma$ ,  $m_\chi > 91$  GeV, all other charged SUSY-partners above 80 GeV:  $0 < h \lesssim 0.75$ ;  $h$  decreases for heavy  $m_{\tilde{t}}, m_{\tilde{\chi}}, m_{H^\pm}$

SUGRA:  $h \lesssim 0.4$ ; calc. only valid up to moderate  $\tan \beta$  (see next slide)

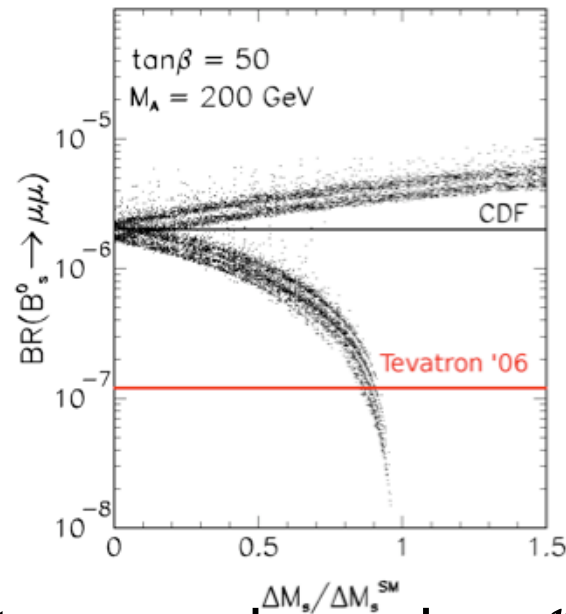


II. large  $\tan \beta$ : [hep-ph/0210145](#)  $h = \underbrace{h_{H^\pm}}_{<0} + \underbrace{h_{\chi^\pm}}_{>0} + \underbrace{h_{DP}}_{<0}, \sigma = 0$

$h < 0$  in most of the parameter space;  $h(B_d) \neq h(B_s)$  due to DP

DP: double penguin from neutral Higgses:  $DP(B_s) \propto m_s \tan \beta^4$  big!

$\Delta m_s$  correlated with  $\mathcal{B}(B_s \rightarrow \mu^+ \mu^-) < 1.0 \cdot 10^{-7}$  [hep-ph/0207241](#)



$\Delta m_s$  is lowered w.r.t. SM; predicts upper bound on  $\mathcal{B}(B_s \rightarrow \mu^+ \mu^-)$

$W = QY_u H_u U + QY_d H_d D + LY_e H_d E + \lambda H_d H_u N - \frac{1}{3}kN^3$   $N$ :singlet  
at large  $\tan \beta$ : naturally light pseudoscalar  $A_1^0$ , radiatively stable

$A_1^0$  masses as low as  $\mathcal{O}(10\text{MeV})$  viable [hep-ph/0404220](#)

iff very light and weakly coupled,  $A_1^0$  becomes missing E

$h^0 \rightarrow A^0 A^0$  very important for Higgs searches [hep-ph/0005308](#)

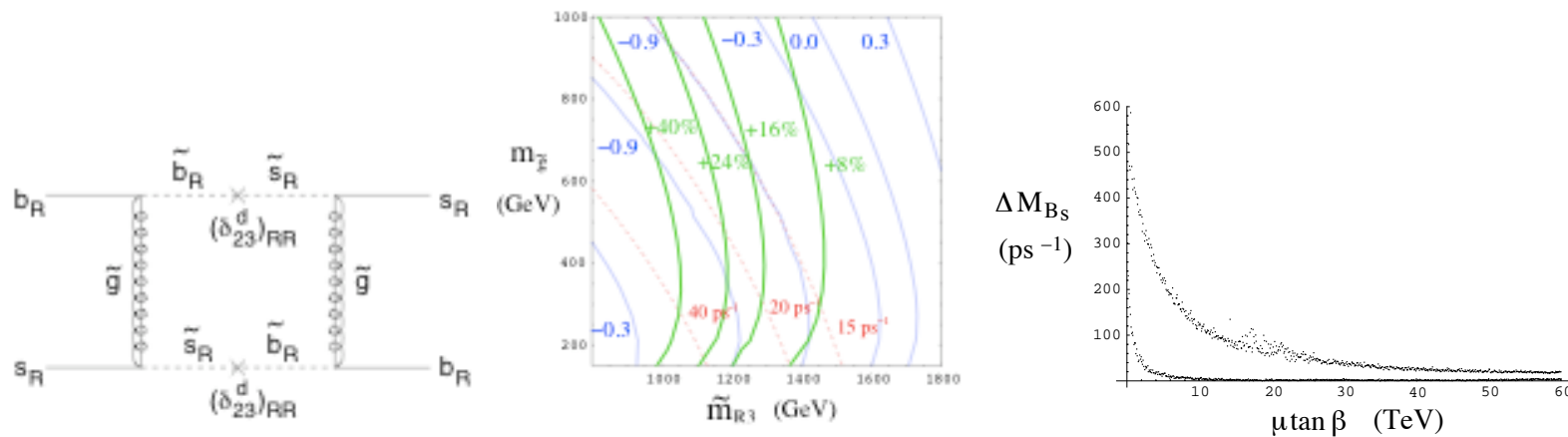
NMSSM at large  $\tan \beta$ :  $h = h_{MSSM}^{large \tan \beta} + h_{A_1^0}$ ;  $h_{A_1^0} \propto \tan \beta^2 / m_{A_1^0}^2$

$$h_{A_1^0}(B_d) = h_{A_1^0}(B_s)$$

from  $\Delta m_d$ -measurement:  $|h| \lesssim 0.4$

unlike MSSM, no correlation with  $\mathcal{B}(B_s \rightarrow \mu^+ \mu^-)$  [hep-ph/0404220](#)

large  $\nu_\mu - \nu_\tau$  mixing in  $SO(10)$  GUT models implies large mixing between right handed  $\tilde{s} - \tilde{b}$ :  $(\delta_{23}^D)_{RR}$  large and complex Figs from [hep-ph/0212180](http://hep-ph/0212180)

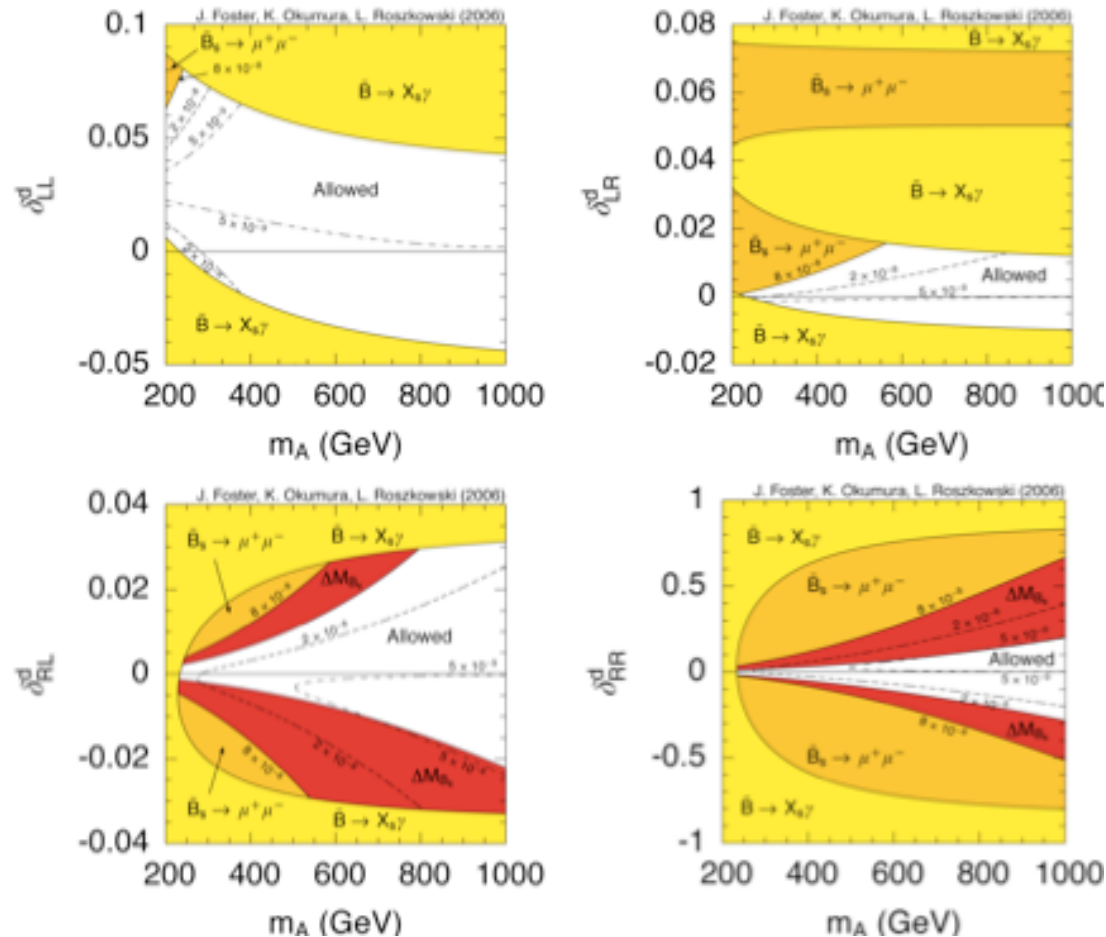


implications:  $\Delta m_s$  can be huge  $\sim 100 \text{ ps}^{-1}$  (range in right fig) but can be also SM-like (middle fig):

needs heavy superpartners  $m_{Q2,Q3,D2} \sim 2\text{TeV}$   
or small couplings  $(\delta_{23}^D)_{RR}$

(green: percent increase in  $B(b \rightarrow s\gamma)$ , blue:  $S_{\Phi K_S}$ )

# Bounds on beyond MFV flavor mixing



$m_{\tilde{g},q} = 1 \text{ TeV}$ ,  $\tan \beta = 40$ ,  $\mu = -A_u = 500 \text{ GeV}$  Figs from [hep-ph/0604121](https://arxiv.org/abs/hep-ph/0604121)

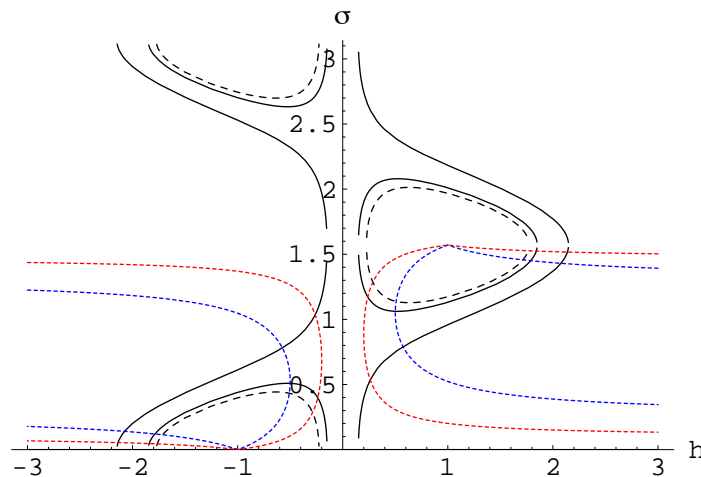


# Further information on $\Delta B = 2$ : getting the phase

Time-dependent CP-asymmetries  $B_s \rightarrow J/\Psi\Phi$  (predom. SM  $b \rightarrow c\bar{c}s$  amplitude)

$$S = \eta_{CP} \cdot \sin [2\beta_s - \arg(1 + h e^{2i\sigma})] \quad \beta_s = \arg(-V_{ts}V_{tb}^*/(V_{cs}V_{cb}^*)) \simeq \eta\lambda^2$$

cf.  $B_d \rightarrow J/\Psi K_S$ :  $S(J/\Psi K_S) = \sin [2\beta + \arg(1 + h_d e^{2i\sigma_d})]$



$S = 0.2$ ,  $S = 0.5$ , **SM**:  $S/\eta_{CP} = \sin 2\beta_s = 0.038 \pm 0.003$  SM CKM fit

$B_s \rightarrow J/\Psi\Phi$  vs.  $B_s \rightarrow (\eta_c, \chi_{0,1,2})\Phi$ : **NP on decay amplitude** hep-ph/0307251

$$S(B_s \rightarrow \Phi\Phi) = \eta_{CP} \sin [2\beta_s - \arg(1 + h e^{2i\sigma}) + \Delta\varphi(b \rightarrow s\bar{s}s\text{-amplitude})]$$

asymmetry into wrong sign leptons:

$$A_{SL}^q \equiv \frac{\Gamma(\bar{B}_q^0(t) \rightarrow \ell^+ X) - \Gamma(B_q^0(t) \rightarrow \ell^- X)}{\Gamma(\bar{B}_q^0(t) \rightarrow \ell^+ X) + \Gamma(B_q^0(t) \rightarrow \ell^- X)} = \text{Im} \frac{\Gamma_{12}^q}{M_{12}^{q SM} (1 + h_q \exp 2i\sigma_q)}$$

$A_{SL} = 0$  means no CP violation in  $\Delta B = 2$  mixing

data:  $A_{SL}^d = 0.0011 \pm 0.0055$  CLEO,Belle,BaBar SM:  $A_{SL}^d \simeq -1 \cdot 10^{-3}$  hep-ph/0202010

collider:  $A_{SL} \equiv \frac{\Gamma(b\bar{b} \rightarrow \mu^+ \mu^+ X) - \Gamma(b\bar{b} \rightarrow \mu^- \mu^- X)}{\Gamma(b\bar{b} \rightarrow \mu^+ \mu^+ X) + \Gamma(b\bar{b} \rightarrow \mu^- \mu^- X)} = -0.0027 \pm 0.0029$  OPAL,ALEPH,DØ

$A_{SL} \simeq 0.6 A_{SL}^d + 0.4 A_{SL}^s$  extract  $A_{SL}^s = -0.008 \pm 0.011$  hep-ph/0604112, 0605028

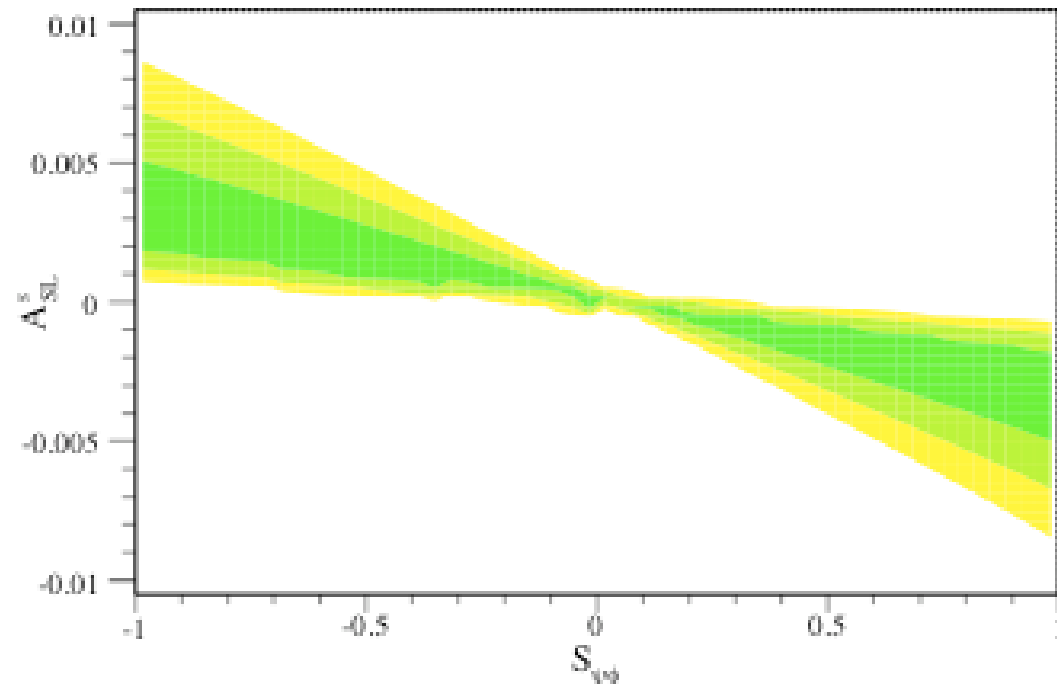
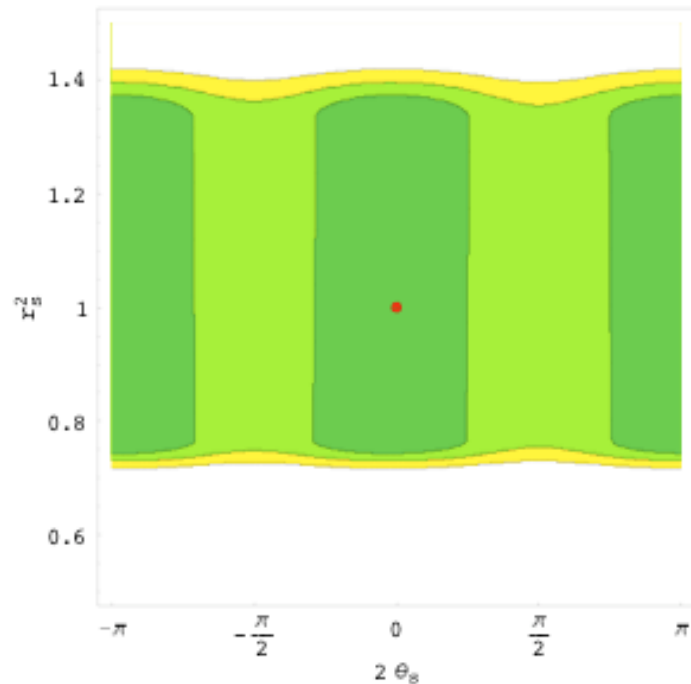
other parametrization:  $M_{12}^s / M_{12}^{s SM} = r_s^2 \exp(2i\vartheta_s) = 1 + h \exp(2i\sigma)$

$$A_{SL}^s = - \underbrace{\text{Re}\left(\frac{\Gamma_{12}^s}{M_{12}^s}\right)^{SM}}_{-0.0040 \pm 0.0016} \frac{\sin 2\vartheta_s}{r_s^2} \rightarrow 2\vartheta_s \simeq \pi/2 \text{ disfavored}$$

$$\Delta\Gamma_{CP}^s = -0.22 \pm 0.08 ps^{-1} \text{ (CDF,ALEPH,D}\emptyset) \text{ vs. } \Delta\Gamma_{CP}^s = \underbrace{(\Delta\Gamma^s)^{SM}}_{-0.07 \pm 0.03 ps^{-1}} \cos^2 2\vartheta_s$$

$\rightarrow 2\vartheta_s \simeq \pm\pi/2$  disfavored

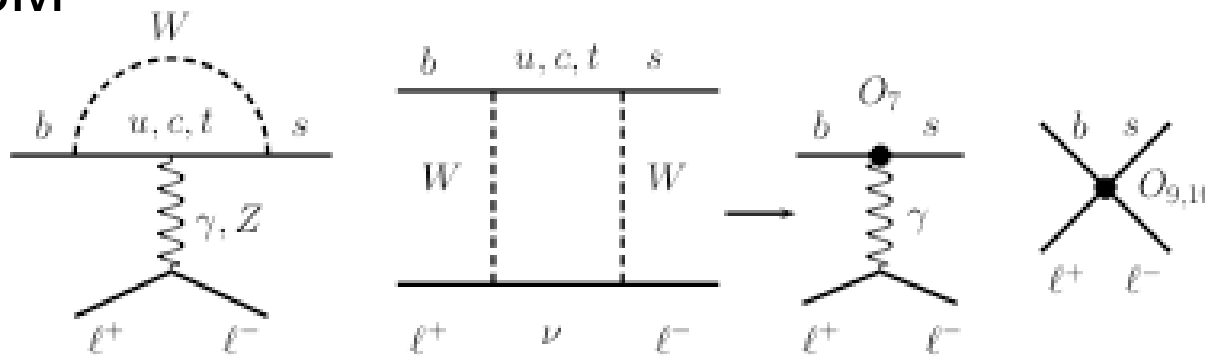
Figs from hep-ph/0605028,0604112



$S, A_{SL}, \Delta\Gamma, \Delta m$  correlated; constraints are still weak

# SM tests with $b \rightarrow s\gamma, b \rightarrow sl^+\ell^-$ decays

diagrams in SM



$$\mathcal{H}_{\text{eff}} = -4 \frac{G_F}{\sqrt{2}} V_{tb} V_{ts}^* \sum C_i(\mu) O_i(\mu)$$

dipole operators  $O_7 \propto \bar{s}_L \sigma_{\mu\nu} b_R F^{\mu\nu}$

$O_8 \propto \bar{s}_L \sigma_{\mu\nu} b_R G^{\mu\nu}$

4-Fermi operators  $O_9 \propto (\bar{s}_L \gamma_\mu b_L)(\bar{\ell} \gamma^\mu \ell)$

$O_{10} \propto (\bar{s}_L \gamma_\mu b_L)(\bar{\ell} \gamma^\mu \gamma_5 \ell)$

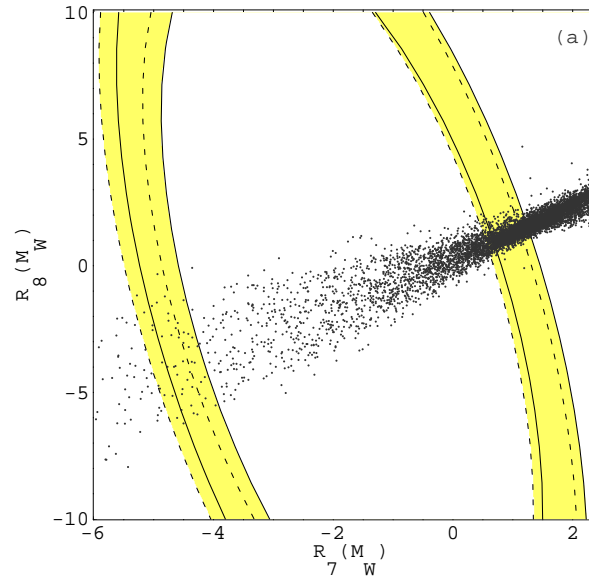
NP in Wilson coefficients  $C_i = C_i^{SM} + C_i^{NP}$  or new operators

model-independent analysis:  $Br$ 's,  $A_{CP}$ ,  $A_{FB} = f(C_i) \rightarrow \text{fit!}$  [hep-ph/9408213](https://arxiv.org/abs/hep-ph/9408213)

# Constraints from $b \rightarrow s\gamma$ branching ratio

$$\mathcal{B}(b \rightarrow s\gamma)_{LO} \sim |C_7(m_b)|^2 \quad \text{at NLO } R \equiv \frac{C^{SM} + C^{NP}}{C^{SM}} \quad \text{hep-ph/0112300}$$

$R_8 - R_7$ :



theory errors renorm. scale and charm mass **solid:pole, dashed:MS**

scatter points: MSSM with MFV:  $C_7 = \underbrace{C_7^{SM} + C_7^{H^\pm}}_{<0} + C_7^{\chi^\pm}$

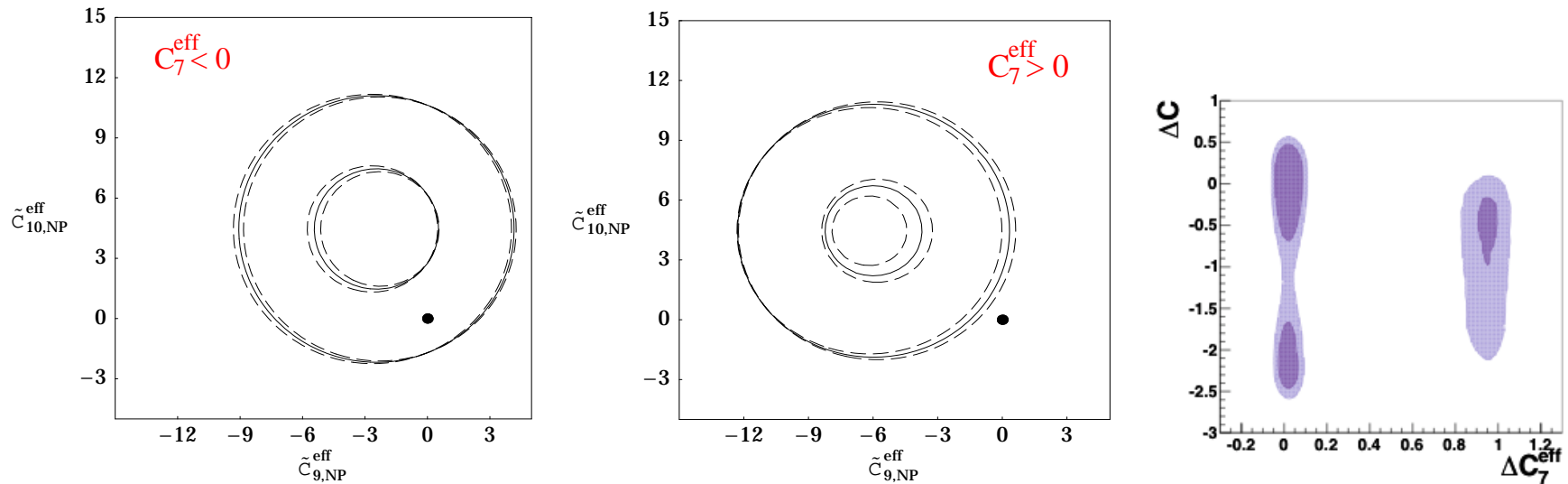
$C_7^{\chi^\pm} \propto \mu A_t \tan \beta f(m_{\tilde{t}_i}, m_{\tilde{\chi}_j}) m_b / (v(1 + \epsilon \tan \beta)); \epsilon \propto (\alpha_s/\pi) \mu m_{\tilde{g}} \tan \beta$

beyond MFV: gluino loops with down squark-mixing  $\delta_{23}^D$  e.g. hep-ph/0212397

# Combined $b \rightarrow s\ell^+\ell^-$ and $b \rightarrow s\gamma$ data

$$\frac{d\Gamma(B \rightarrow X_s \ell^+ \ell^-)}{d\hat{s}} = \left(\frac{\alpha}{4\pi}\right)^2 \frac{G_F^2 m_b^5 |V_{ts}^* V_{tb}|^2}{48\pi^3} (1 - \hat{s})^2 \left[ (1 + 2\hat{s}) \left( |C_9^{\text{eff}}|^2 + |C_{10}^{\text{eff}}|^2 \right) f_1 + 4(1 + 2/\hat{s}) |C_7^{\text{eff}}|^2 f_2 + 12\text{Re} \left( C_7^{\text{eff}} C_9^{\text{eff}*} \right) f_3 + f_c \right] \quad f_i: 1/m_{c,b}^2 \text{ corr.}$$

$$\mathcal{B}(B \rightarrow X_s \mu\mu)_{SM} = 4.3 \pm 0.7 \cdot 10^{-6} \text{ vs data} = 4.3 \pm 1.2 \cdot 10^{-6} \text{ (full } q^2)$$



non-SM sign  $C_7^{\text{eff}} > 0$  disfavored iff no BSM ops [ph/0410155, 0505110 Inami-Lim-study](#)

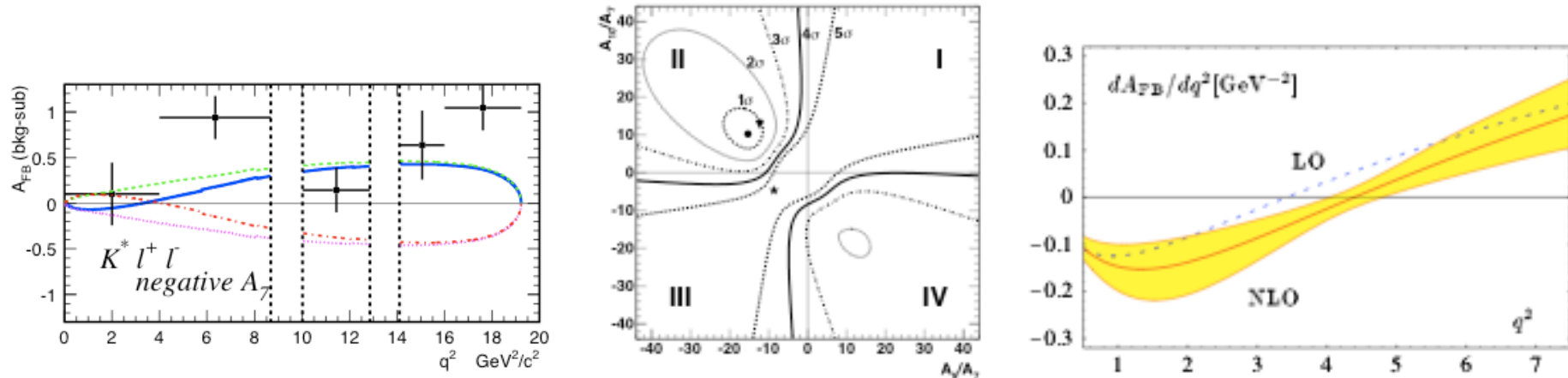
MSSM+MFV:  $C_9, C_{10}$  near SM, not  $\tan\beta$  enhanced [hep-ph/0112300](#)

check for BSM operators: (pseudo)scalar (neutral Higgs),  $L \leftrightarrow R$

# NP sensitivity of $b \rightarrow sl^+l^-$ spectra, perspectives

$A_{FB}$ : # forward - # backward  $l^+$  in dilepton CMS w.r.t.  $\bar{B}$  (CP-odd)

$A_{FB}(\hat{s}) \sim -\text{Re} [C_{10}^*(C_7^{\text{eff}} + \beta(\hat{s})C_9^{\text{eff}})]$  also  $B \rightarrow K^*l^+l^-$  Belle 0508009,0603018



shape sensitive to sign  $C_7$ ;  $A_{FB} \propto C_{10}$

zero  $X_S$ :  $\hat{s}_{SM}^{NNLL} = 0.162 \pm 0.002(8)$  hep-ph/0208088,0209006

$K^*$ :  $s_{SM}^{NLO} = 4.4 \pm 0.4 \text{ GeV}^2$  hep-ph/0106067

CP  $A_{FB}^{CP} \equiv \frac{A_{FB} + \bar{A}_{FB}}{A_{FB} - \bar{A}_{FB}} \sim \arg C_{10} \arg C_9^{\text{eff}}$ ;  $A_{FB}^{CP}|_{SM} \lesssim 10^{-3}$  hep-ph/0006136

full angular analysis  $B \rightarrow K^*(\rightarrow K\pi)l^+l^-$  hep-ph/9907386

$$d\Gamma^4 \sim I(s, \Theta_l, \Theta_{K^*}, \Phi) ds d\cos\Theta_l d\cos\Theta_{K^*} d\Phi$$

# More model-independent studies, neutral Higgses

Higgses split between  $\mu^+\mu^-$  and  $e^+e^-$  in  $b \rightarrow sl^+\ell^-$  ratios with SAME cut on dilepton mass [hep-ph/0310219](http://hep-ph/0310219)

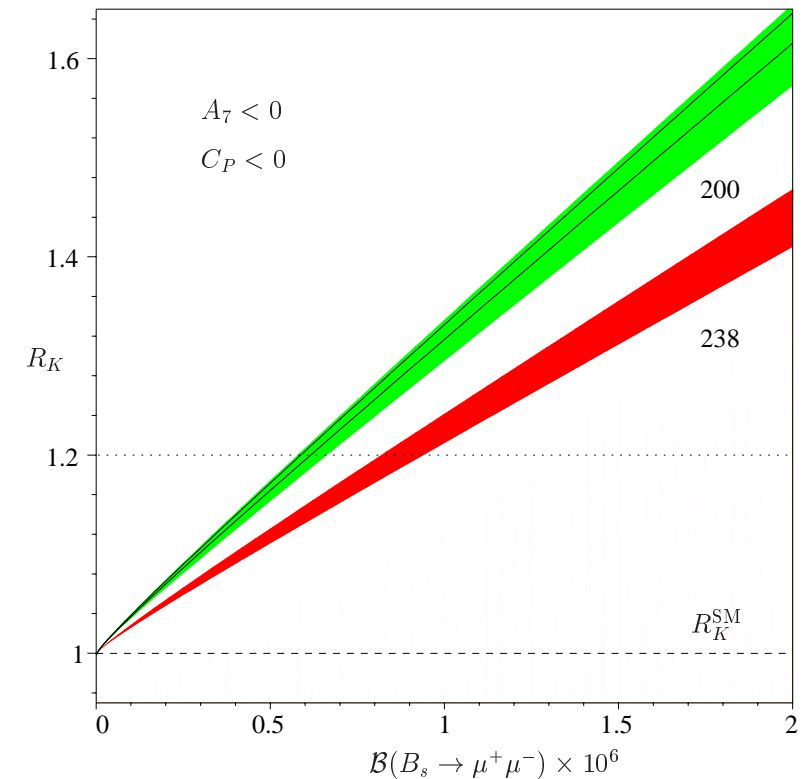
$$R_H \equiv \frac{\int_{4m_\mu^2}^{q_{\max}^2} dq^2 \frac{d\Gamma(B \rightarrow H\mu^+\mu^-)}{dq^2}}{\int_{4m_\mu^2}^{q_{\max}^2} dq^2 \frac{d\Gamma(B \rightarrow He^+e^-)}{dq^2}} \quad H = X_s, K^{(*)}$$

$$R_H^{SM} = 1 + \mathcal{O}(m_\mu^2/m_b^2) \quad \text{very clean}$$

$$R_K = 1.06 \pm 0.48 \pm 0.05 \quad \text{BaBar 0507005}$$

$$R_K = 1.38_{-0.41-0.07}^{+0.39+0.06} \quad \text{Belle 0410006}$$

constrain (pseudo)-scalar couplings



$R_K$  constrains  $C_{S,P} + C'_{S,P}$ ,  $\mathcal{B}(B_s \rightarrow \mu^+\mu^-)$  constrains  $C_{S,P} - C'_{S,P}$   
correlation breaks down if  $C_{S,P} \not\propto m_\ell$



- CKM@tree: input for SM tests and flavor model building thanks to tremendous exp (Belle, BaBar) and th (loops, HQET, lattice, ..) efforts
- $b \rightarrow s\ell\ell$  modes under th and exp investigation; model independent analysis (w.  $b \rightarrow s\gamma$ )  $\rightarrow$  do  $e$  and  $\mu$  separately long term goal: angular analysis in  $B \rightarrow K^*(\rightarrow K\pi)\ell\ell$
- tool in penguin-physics: **multi-observable analyses and fits**
- $b \rightarrow d$  FCNCs beginning to be probed
- $B_s, B_c, \Lambda_b$ -physics coming up CDF&D0; much more from LHC(b)

- data on  $B_s - \bar{B}_s$  mixing touch unknown territory
  - with  $\Delta m_s$  near  $\Delta m_s^{\text{SM}}$ :
- generically,  $O(1)$  NP allowed with some amount of fine-tuning and CP-phase; further constraints from  $\Delta\Gamma$ ,  $A_{SL}$  (starting) and CP-asymmetries (to be done)
- models with MFV, such as CMSSM can accommodate SM-like  $B_s - \bar{B}_s$ -mixing within errors  $|h| < \epsilon$
- models beyond MFV are constrained significantly
- so far no significant conflict with SM/MFV; many FCNCs only weakly or just un-constrained  $V + A, \tau^\pm, \nu\bar{\nu}, b \rightarrow d, B_s$ -physics

# Some (further) SM tests with $b$ -physics

	experiment	SM	comments
$a_{CP}(b \rightarrow s\gamma)$	$0.4 \pm 3.6 \%$	$0.42 \pm 0.17 \%$ <a href="#">hep-ph/0312260</a>	CPX in $bs\gamma, g$
$a_{CP}(b \rightarrow d/s\gamma)$	$-0.110 \pm 0.116$ <a href="#">BaBar'05</a>	$10^{-9}$ <a href="#">hep-ph/0312260</a>	test MFV
$S_{K_S\pi^0\gamma}$	$0.00 \pm 0.28$ <a href="#">Belle/BaBar'05</a>	$-2m_s/m_b$	V+A FCNCs
spin $\Lambda_b \rightarrow \Lambda\gamma$	—	$\sim m_s/m_b$	V+A FCNCs
$\mathcal{B}(B \rightarrow X_s g)$	$< 9\%$ <a href="#">CLEO'97</a>	$5.0 \pm 1.0 \cdot 10^{-3}$	NP in $bsg$
TDCPA $b \rightarrow s\bar{s}s$	$S_{ave} = 0.50 \pm 0.06$	$\sin 2\beta + \Delta S$	CPX
$\mathcal{B}(B \rightarrow X_s \bar{\mu}\mu)$	$4.3 \pm 1.2 \cdot 10^{-6}$	$4.3 \pm 0.7 \cdot 10^{-6}$	$q^2$ -spectra
$a_{CP}(B \rightarrow X_s \bar{\ell}\ell)$	$-0.22 \pm 0.26$	$-0.2 \pm 0.2 \%$ <a href="#">hep-ph/9812267</a>	CPX
$A_{FB}^{CP}(B \rightarrow K^* \bar{\ell}\ell)$	—	$\lesssim 10^{-3}$ <a href="#">hep-ph/0006136</a>	CPX in $bsZ$
$R_K \mu\mu$ vs. $ee$	$1.06 \pm 0.48$ <a href="#">BaBar'05</a>	$1 + \mathcal{O}(m_\mu^2/m_b^2)$ <a href="#">hep-ph/0310219</a>	non-SM Higgs
$\mathcal{B}(B \rightarrow K\nu\bar{\nu})$	$< 3.6 \cdot 10^{-5}$ <a href="#">Belle'05</a>	$3.8_{-0.6}^{+1.2} \cdot 10^{-6}$	$O(10)$ from SM
$\mathcal{B}(B_s \rightarrow \mu^+\mu^-)$	$< 1.0 \cdot 10^{-7}$ <a href="#">CDF'06</a>	$3.2 \pm 1.5 \cdot 10^{-9}$	$O(30)$ from SM
$\mathcal{B}(B_s \rightarrow \tau^+\tau^-)$	$< \mathcal{O}(5\%)$	$7.2 \pm 1.1 \cdot 10^{-7}$	$O(10^5)$ from SM
$\mathcal{B}(B_d \rightarrow \tau^+\tau^-)$	$< 3.4 \cdot 10^{-3}$ <a href="#">BaBar'05</a>	$2.1 \pm 0.3 \cdot 10^{-8}$	$O(10^5)$ from SM

$$\Delta m_s^{\text{SM}} \sim (V_{tb}^* V_{ts})^2 \times \frac{g^4}{16\pi^2} \times \frac{(\bar{b}\Gamma_s)(\bar{b}\Gamma'_s)}{m_W^2}$$

$$\Delta m_s^{\text{NP}} \sim \text{flavor} \times \text{loop/tree} \times \frac{(\bar{b}\Gamma_s)(\bar{b}\Gamma'_s)}{\Lambda^2}$$

NP in  $B_s$ -mixing not bigger than SM implies:

	NP from loops	tree level NP
MFV	$\Lambda \gtrsim m_W \sim O(100) \text{ GeV}$	$\Lambda \gtrsim 4\pi m_W \sim O(1) \text{ TeV}$
non MFV	$\Lambda \gtrsim m_W/ V_{ts}  \sim O(2-3) \text{ TeV}$	$\Lambda \gtrsim 4\pi m_W/ V_{ts}  \sim O(30) \text{ TeV}$

origin of flavor is still a mystery

whether it is connected to the NP scale  $\Lambda$  we are probing now