

Probing Lepton Number Violation on Three Frontiers



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Overview

- Neutrinos
 - Oscillations
 - Absolute Mass
- Neutrinoless Double Beta Decay
 - Light Neutrino Exchange
 - New Physics Mechanisms
 - In Combination with Cosmology
- Neutrino Mass Models
 - Effective Mass and Seesaw
 - Minimal Left-Right Symmetry
- LFV and LNV at the LHC
- Conclusion

Neutrino Oscillations

- **Neutrino interaction states different from mass eigenstates**

Neutrino flavour can change through propagation

$$\nu_i = \sum_{\alpha} U_{i\alpha} \nu_{\alpha}, \quad \nu_i(t) = e^{-i(E_i t - p_i x)} \nu_i$$

$$\Rightarrow P_{\alpha \rightarrow \beta} = \sin^2(2\theta) \sin^2\left(1.27 \frac{\Delta m^2}{\text{eV}^2} \frac{L/\text{km}}{E/\text{GeV}}\right)$$

- **Solar neutrino oscillations**

Large mixing

- **Atmospheric oscillations**

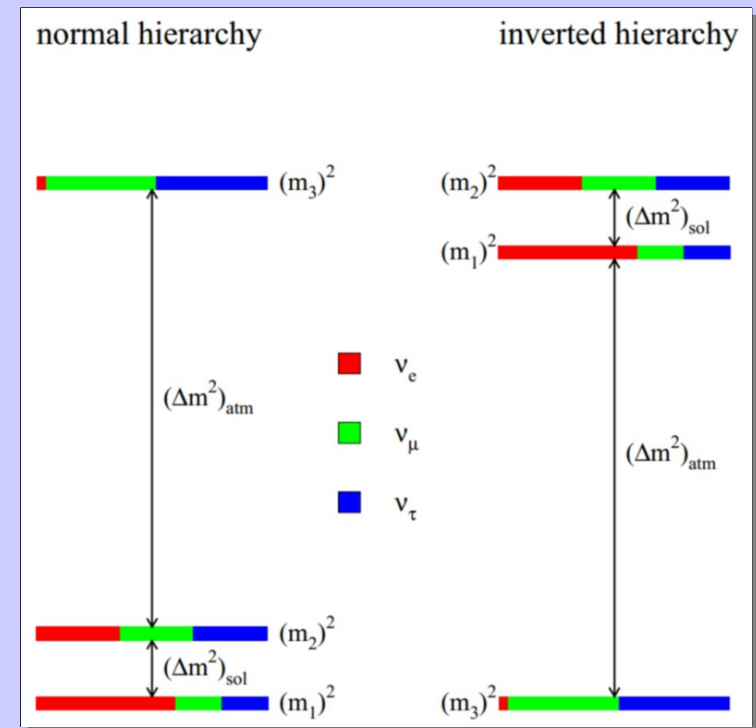
≈ Maximal mixing

- **Reactor and accelerator neutrinos**

$$\sin^2(2\theta_{13}) = 0.092 \pm 0.021$$

- **Experimental unknowns and anomalies**

CP violation? Sign of Δm_{23} ? Sterile Neutrinos?



Absolute Neutrino Mass

- Energy endpoint in Beta decay

$$m_{\beta}^2 = \sum_i |U_{ei}|^2 m_i^2 < (2.2 \text{ eV})^2$$

Katrin: $m_{\beta} \approx 0.2 \text{ eV}$

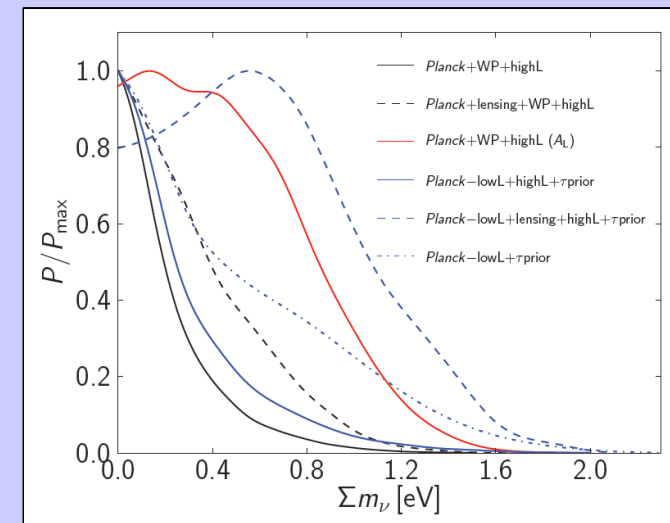
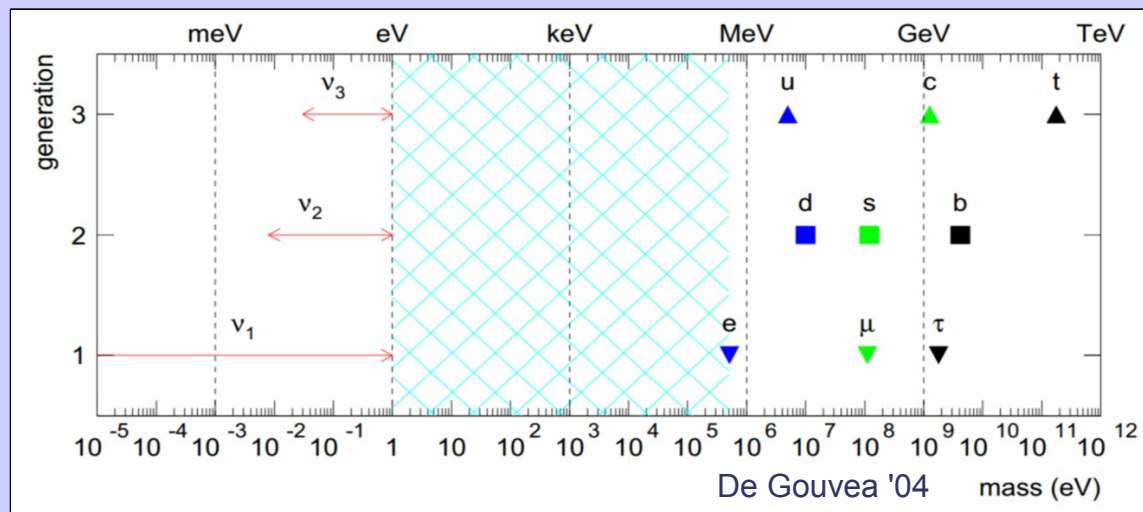
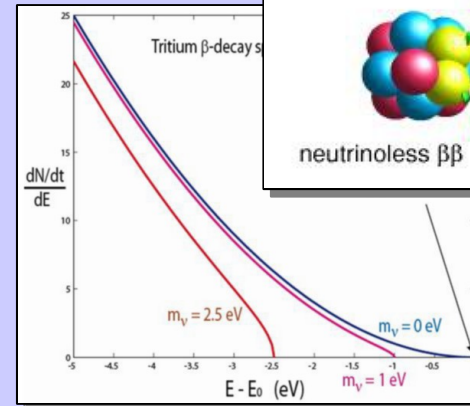
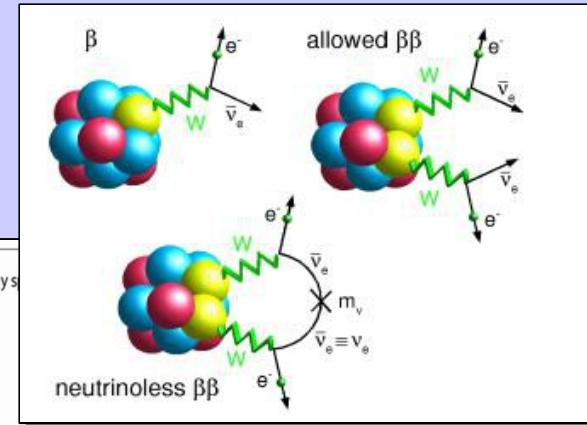
- Impact on Large Scale Structure

$$\Sigma = \sum_i m_i < 0.3 - 1 \text{ eV}$$

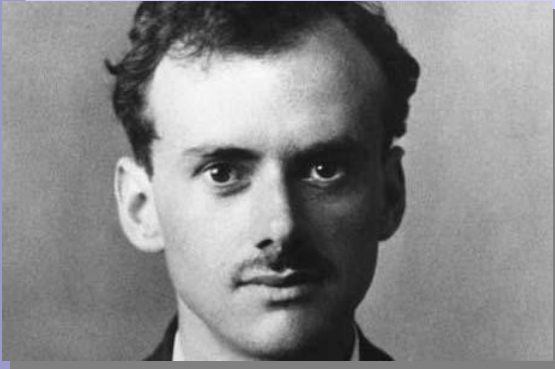
- Neutrinoless Double Beta Decay

$$m_{\beta\beta} = \left| \sum_i U_{ei}^2 m_{\nu_i} \right| < 0.2 - 2.0 \text{ eV}$$

Future: $m_{\beta\beta} \approx 0.01 \text{ eV}$



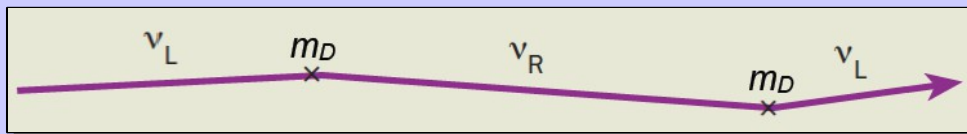
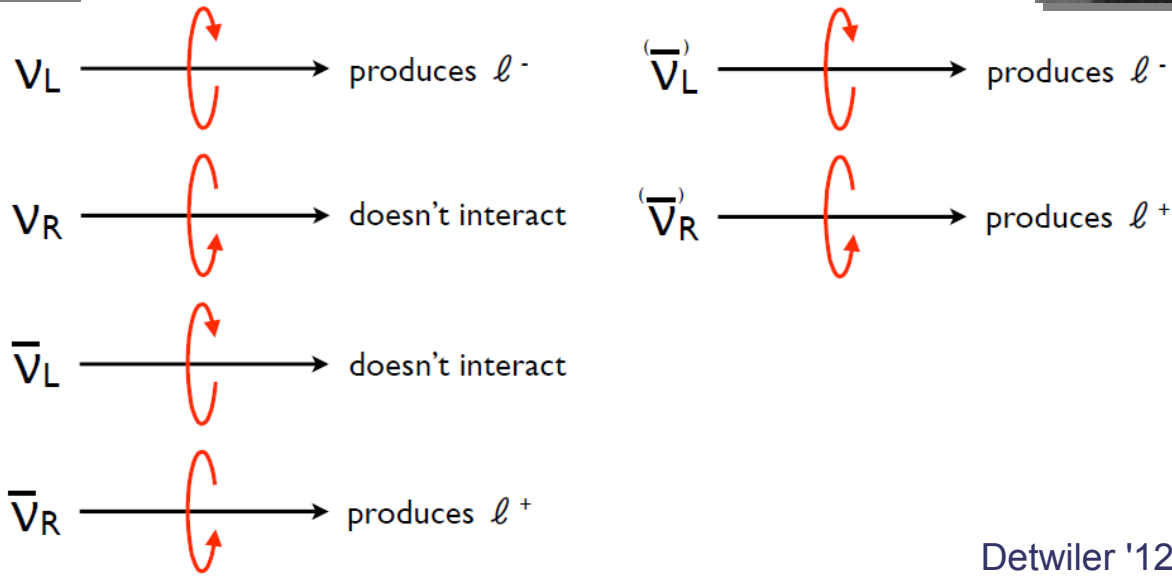
Dirac vs Majorana



Two possibilities to define mass

Dirac ν

Majorana ν

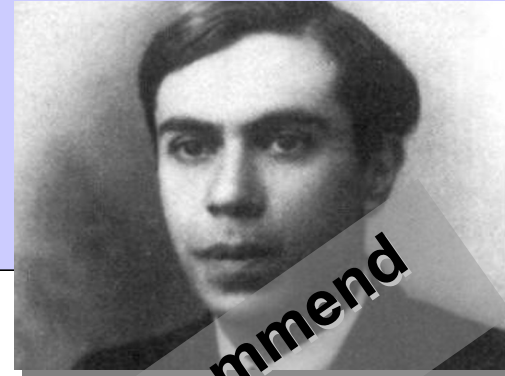
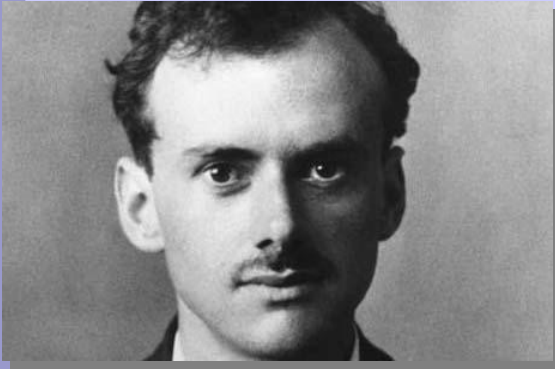


Dirac mass, analogous to other fermions (but with tiny coupling to Higgs)



Majorana mass, using only a left-handed neutrino → Lepton Number Violation

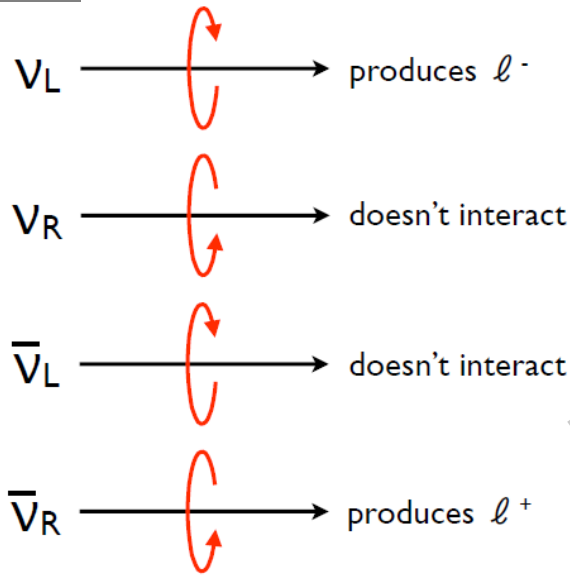
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Two possibilities to define mass

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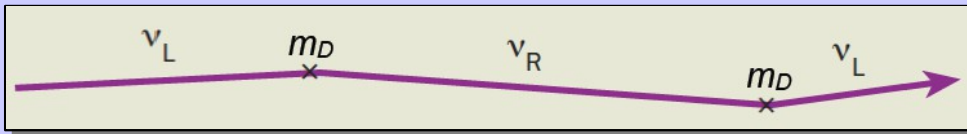
Majorana ν



9 out of 10 theorists recommend Majorana neutrinos



Detwiler



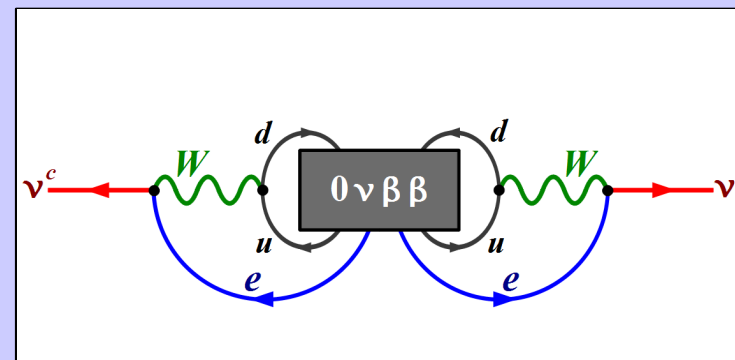
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Neutrinoless Double Beta Decay

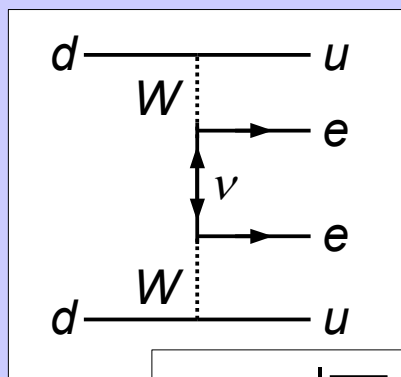
- **Process:** $(A, Z) \rightarrow (A, Z+2) + 2e^-$
- **Uncontroversial detection of $0\nu\beta\beta$ of utmost importance**
 - Prove lepton number to be broken
 - Prove neutrinos to be Majorana particles (Schechter, Valle '82)



$$\delta m_\nu \approx \frac{1}{(16\pi^2)^4} \frac{\text{MeV}^5}{M_W^4} \approx 10^{-23} \text{ eV}$$

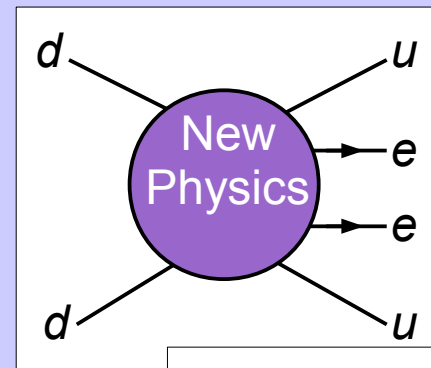
- **Which mechanism triggers the decay?**

Light Neutrino Exchange
(LH Current, Mass Mechanism)



$$T_{1/2}^{-1} \propto \left| \sum_i U_{ei}^2 m_{\nu_i} \right|$$

General Effective Operator

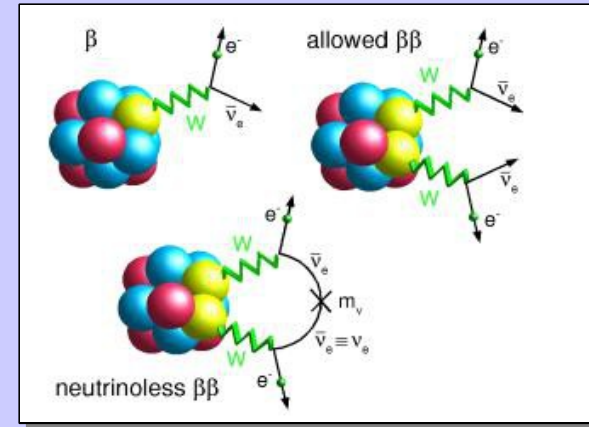
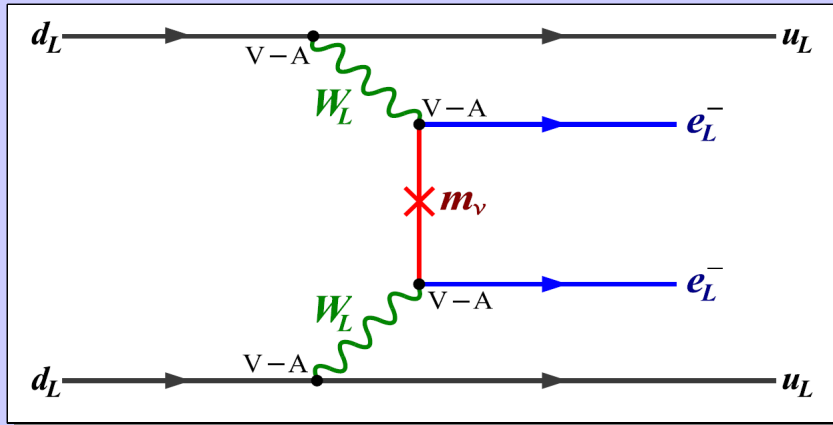


$$\bar{u}\bar{u}\bar{e}\bar{e}dd/M^5$$

$$T_{1/2}^{-1} \approx 10^{25} \text{ y} \rightarrow M \approx 1 \text{ TeV}$$

Light Neutrino Exchange

- Standard Mass Mechanism



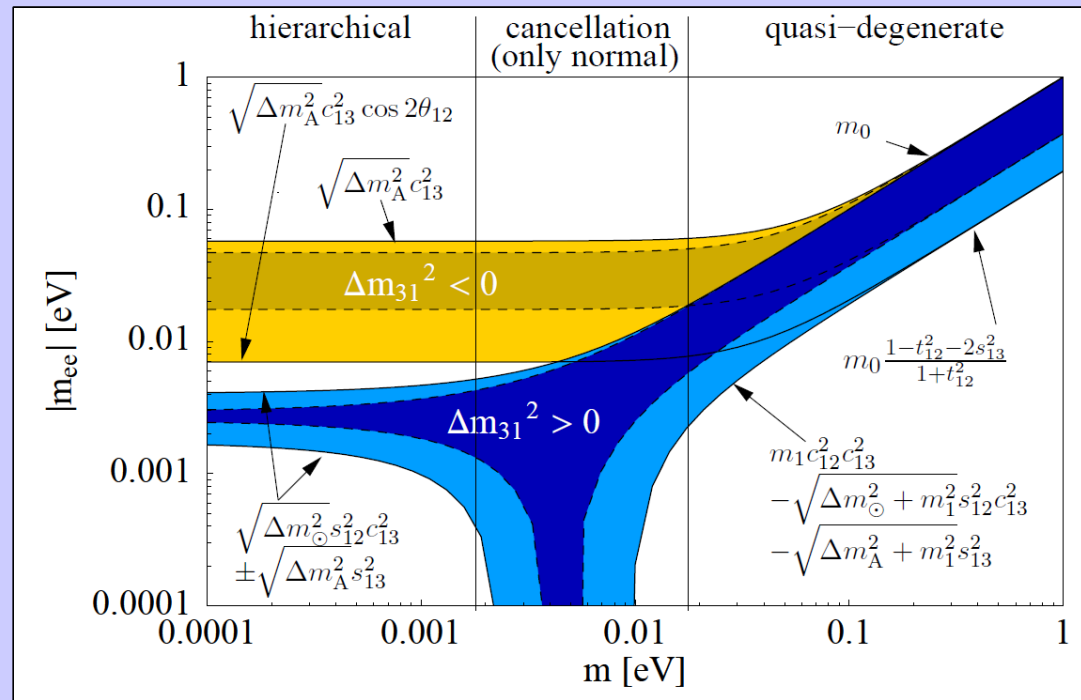
Lindner, Merle, Rodejohann (2005)

- Decay Rate

$$\Gamma = T_{1/2}^{-1} = \frac{m_{\beta\beta}^2}{m_e^2} G^{0\nu} |M^{0\nu}|^2$$

- Effective Mass

$$m_{\beta\beta} = \left| \sum_i U_{ei}^2 m_{\nu_i} \right| \equiv (m_{\nu})_{ee}$$



Nuclear Matrix Elements

- Decay Rate

$$\Gamma = T_{1/2}^{-1} = \frac{m_{\beta\beta}^2}{m_e^2} G^{0\nu} |M^{0\nu}|^2$$

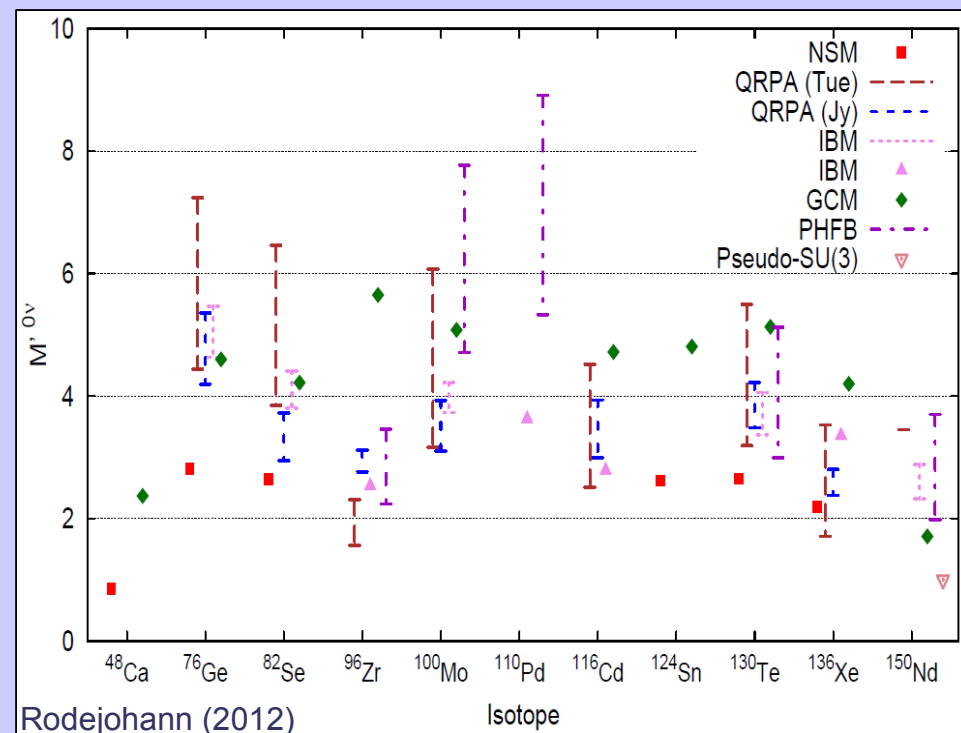
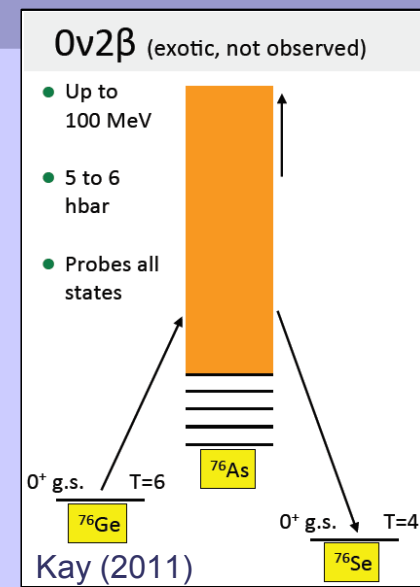
- Requires calculation of matrix element of the nuclear transition via intermediate state

- Many-body problem not solvable from first principle

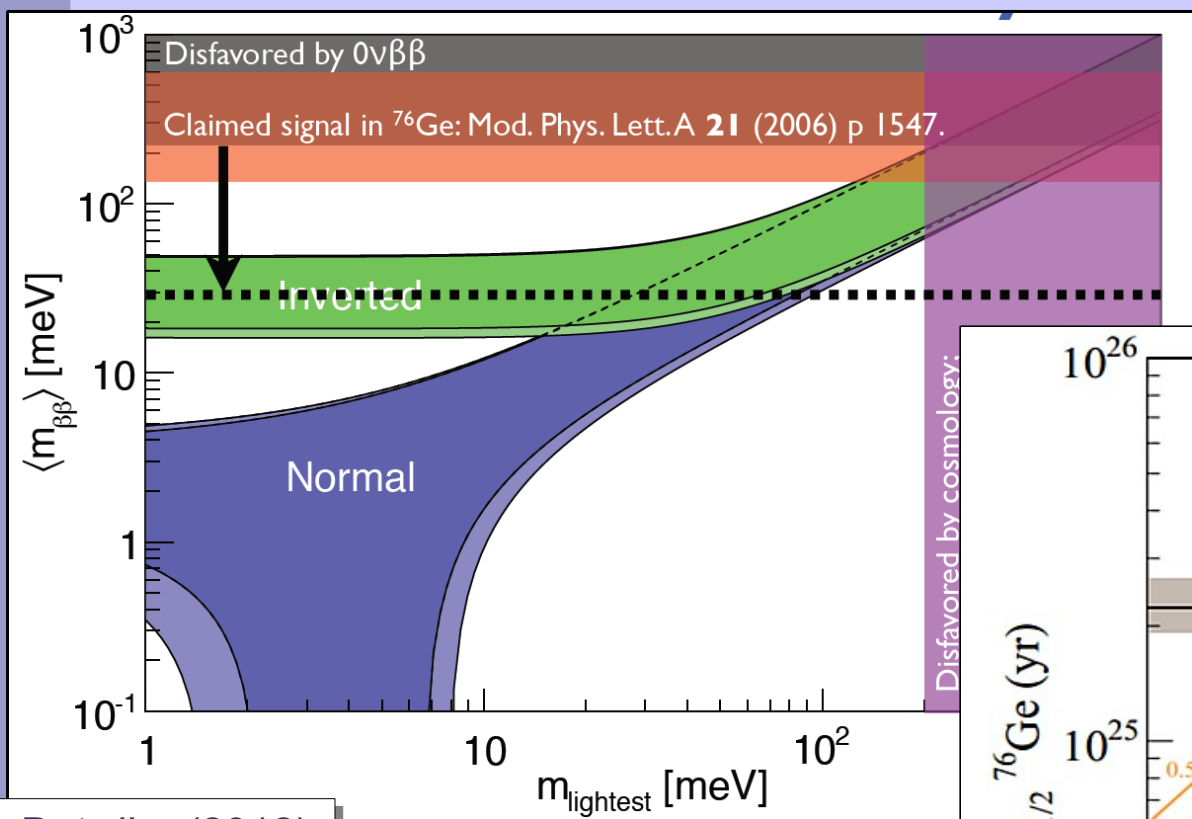
- Factor 2 – 3 uncertainty between different nuclear models

- Shell Models
- QRPA
- IBM

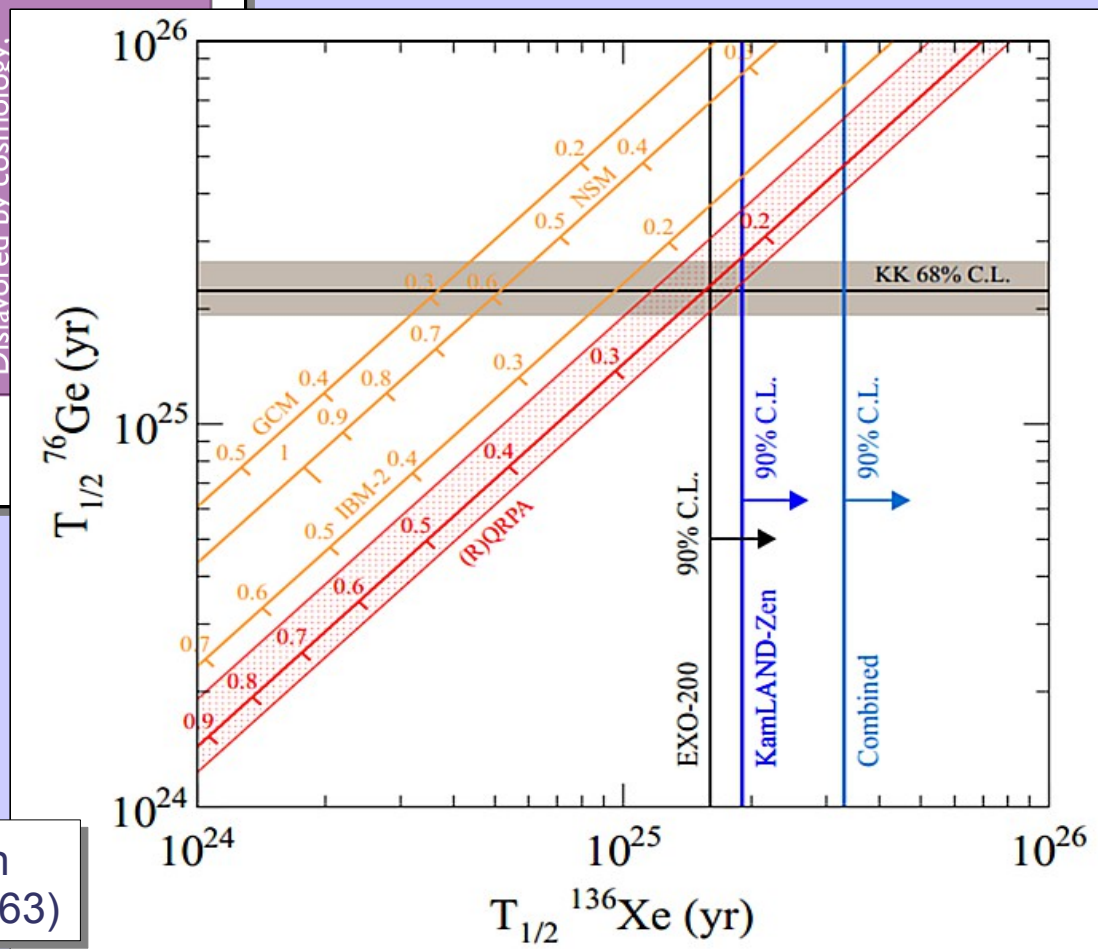
- Important to search for $0\nu\beta\beta$ in several isotopes



Experimental Situation



Detwiler (2012)



KamLAND-Zen
(arXiv:1211.3863)

Combining $0\nu\beta\beta$ and Cosmology

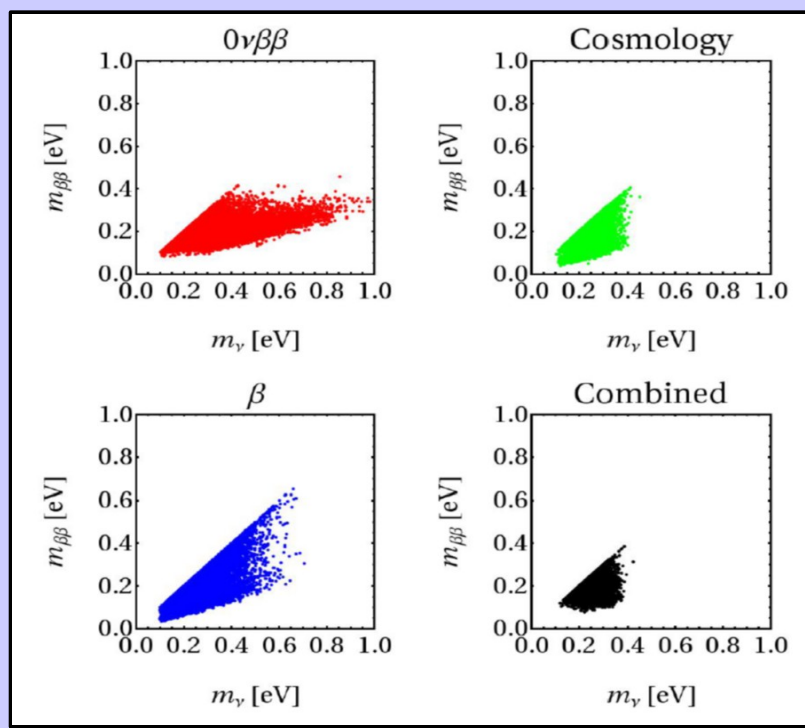
J. Auger, FFD, O. Lahav, I. Sadeh, D. Waters, Work in Progress

• Effective $0\nu\beta\beta$ Observable

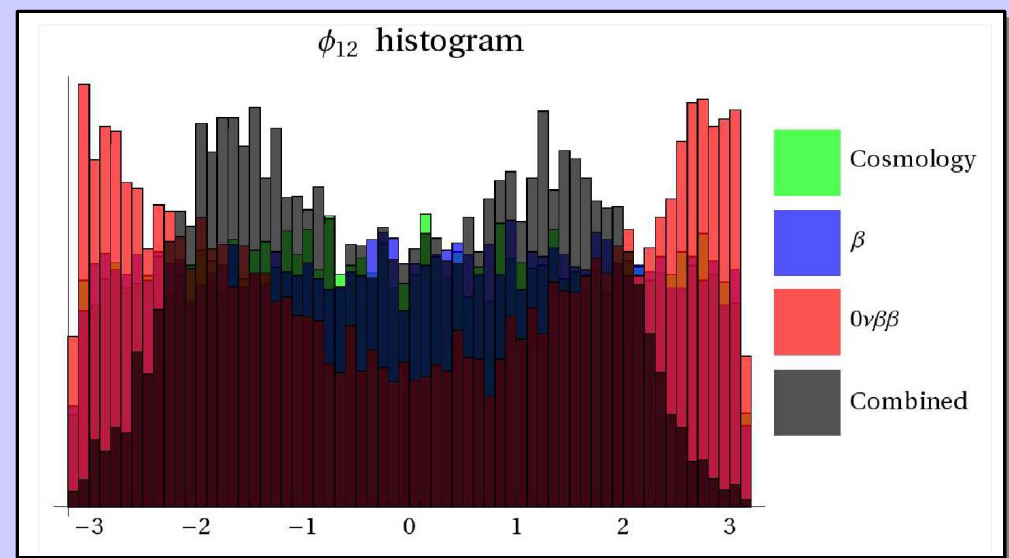
$$m_{\beta\beta} \equiv (m_\nu)_{ee} = \left| \sum_i U_{ei}^2 m_{\nu_i} \right| = \left| m_{\nu_1} |V_{e1}|^2 + m_{\nu_2} |V_{e2}|^2 e^{2i\phi_{12}} + m_{\nu_3} |V_{e3}|^2 e^{2i\phi_{23}} \right|$$

$$\approx m_1 \sqrt{1 - \sin^2(2\theta_{12}) \sin^2(\phi_{12})} \quad \text{quasi-deg. neutrinos}$$

• Interplay of mass probes

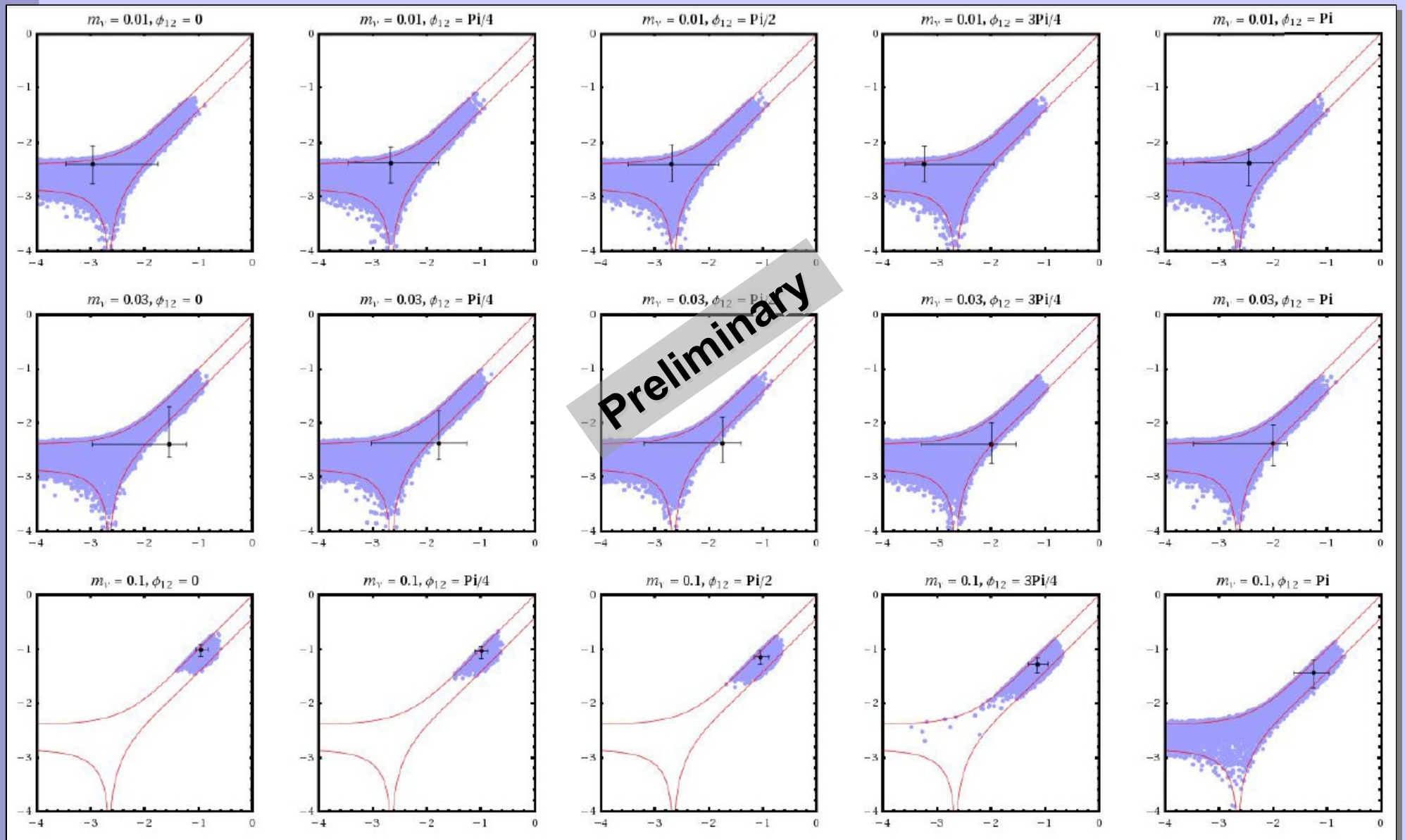


Simulation of observables incl. NME and expected exp. uncertainties with $m_{\nu_1} = 0.25 \text{ eV}, \phi_{12} = \pi/2$



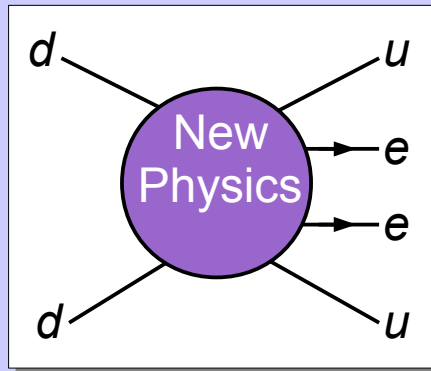
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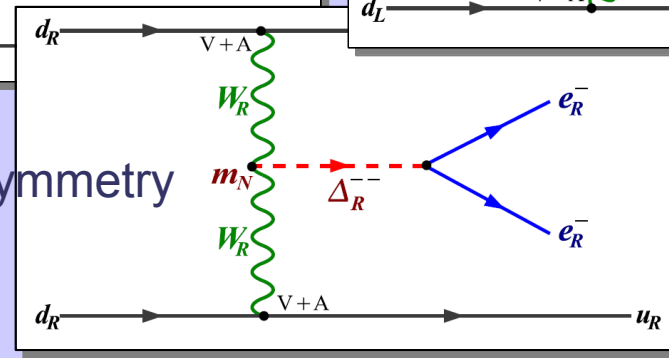
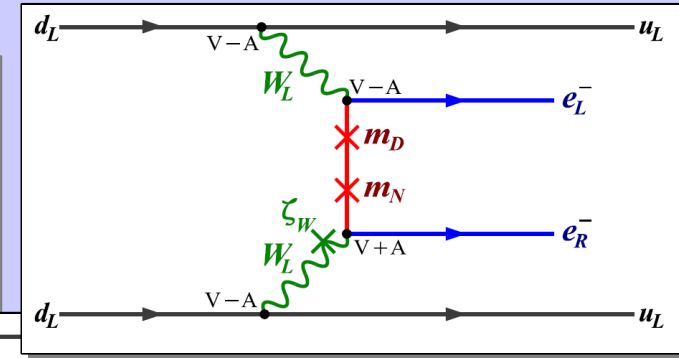
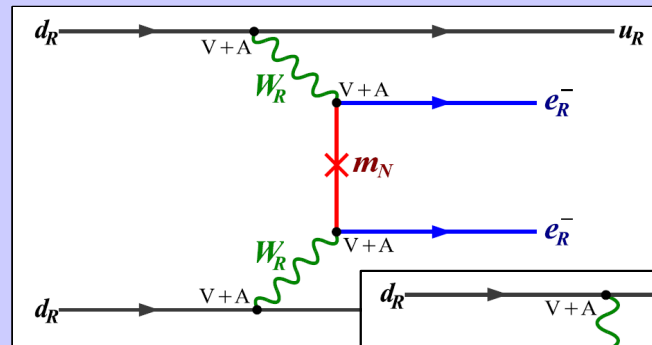


New Physics Contributions to $0\nu\beta\beta$

• Plethora of New Physics Scenarios



=



Left-Right Symmetry

$$\Gamma = T_{1/2}^{-1} = \epsilon_{\text{NP}}^2 G_{\text{NP}}^{0\nu} |M_{\text{NP}}^{0\nu}|^2$$

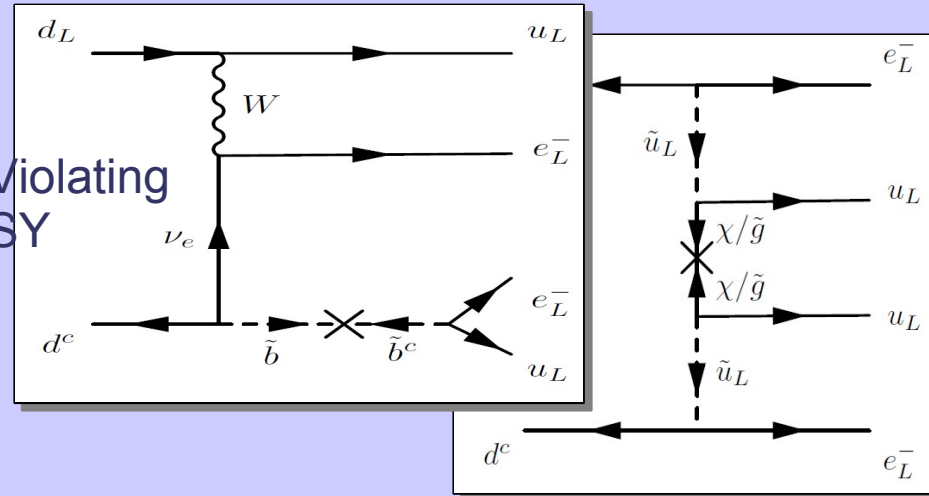
Extra Dimensions

Majorons

R-Parity Violating SUSY

Leptoquarks

...



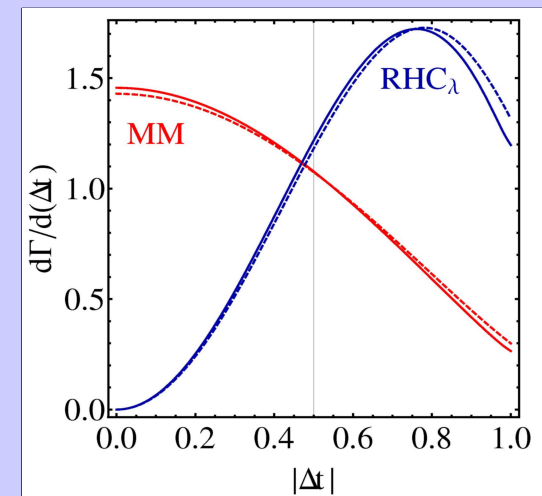
Disentangling New Physics Scenarios

- **Angular and Energy distribution of emitted electrons**

(Doi et al. '83; Ali et al. '06; Arnold et al. '10; FFD, Jackson, Nasteva, Söldner-Rembold '10)

$$\frac{d\Gamma}{dE_1 dE_2 d\cos\theta} = \frac{\Gamma}{2} (1 - k(E_1, E_2) \cos\theta) \quad -1 < k < 1$$

- Linear in $\cos\theta$
- $k(E_1, E_2)$ depends on $0\nu\beta\beta$ mechanism



- **Comparison of $0\nu\beta\beta$ in multiple isotopes**

(FFD, Päs PRL 2007)

$$\frac{T_{1/2}(^A X)}{T_{1/2}(^B Y)} = \frac{G(^B Y) |M(^B Y)|^2}{G(^A X) |M(^A X)|^2}$$

- Depends on $0\nu\beta\beta$ mechanism
- Independent of details of new physics (if one mechanism dominates)

Effective Mass and Seesaw Mechanism

- Effective operator for Majorana neutrino mass

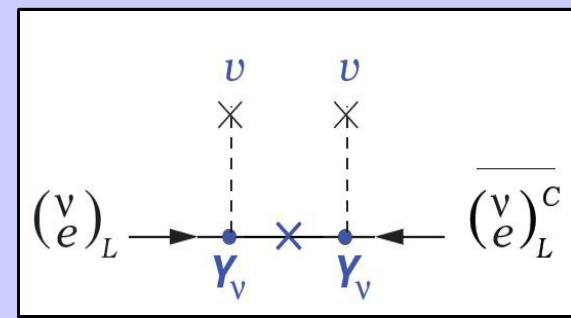
$$L = \frac{1}{2} \frac{h_{ij}}{\Lambda_{\text{LNV}}} (\bar{L}_i^c \cdot \tilde{H}) (\tilde{H}^T \cdot L_j) \rightarrow \frac{1}{2} (m_\nu)_{ij} \bar{\nu}_i^c \nu_j$$

Unique dim-5
Operator

- Seesaw Mechanism

Add right-handed neutrinos to the Standard Model particle content, $M \approx 10^{14}$ GeV

$$L = L_{\text{SM}} - \frac{1}{2} \bar{\nu}_R M \nu_R^c + \bar{\nu}_R Y_\nu L \cdot H_u$$



- Light neutrino mass matrix at low energies

$$m_\nu = m_D^T M^{-1} m_D \text{ for } m_D \ll M_R \quad m_\nu \approx 0.1 \text{ eV} \left(\frac{m_D}{100 \text{ GeV}} \right)^2 \left(\frac{M}{10^{14} \text{ GeV}} \right)^{-1}$$

Effective Mass and Seesaw Mechanism

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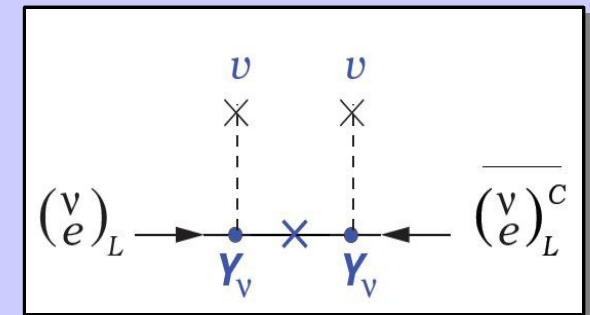
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Unique dim-5
Operator

- Seesaw Mechanism

- Sterile Neutrino Mass Scale Unknown

- $\approx 10^{14}$ GeV: Naive Seesaw, GUTs
 - $> 10^9$ GeV: Thermal Leptogenesis
 - $\approx 10^2$ GeV: Resonant Leptogenesis, Production at LHC
 - ≈ 1 keV: Dark Matter Candidate
 - ≈ 1 eV: Oscillation, Cosmology, (Double) Beta Decay



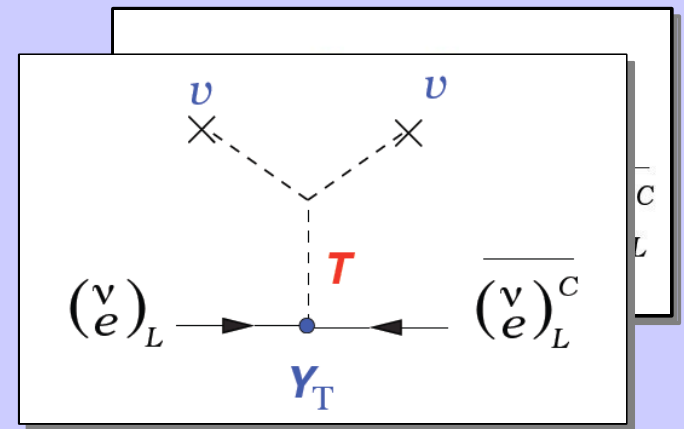
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Unique dim-5 Operator

- Seesaw Mechanism
- Three possible mediators at tree level



Seesaw II

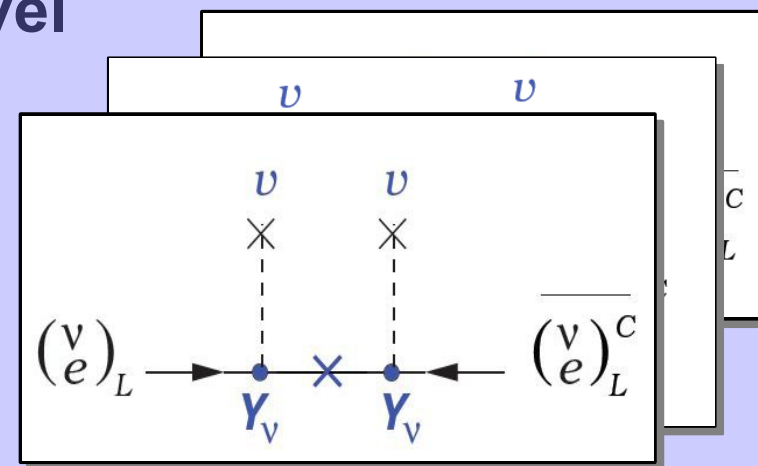
Effective Mass and Seesaw Mechanism

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Unique dim-5 Operator

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Seesaw III

Effective Mass and Loops

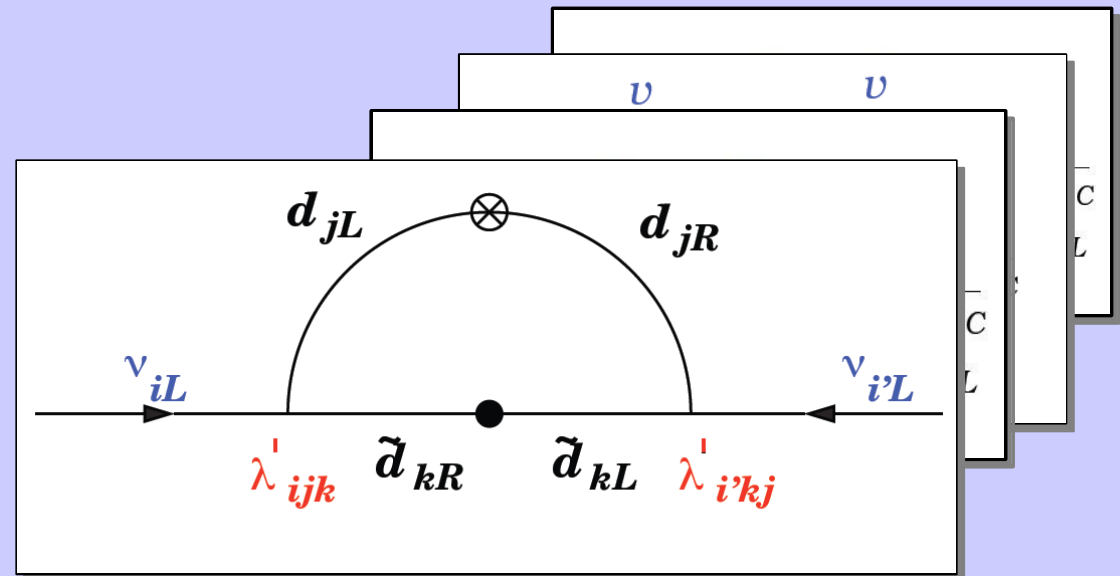
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Unique dim-5
Operator

- Radiative Generation via Loops

Alternative to Seesaw Mechanism



R-Parity Violating SUSY

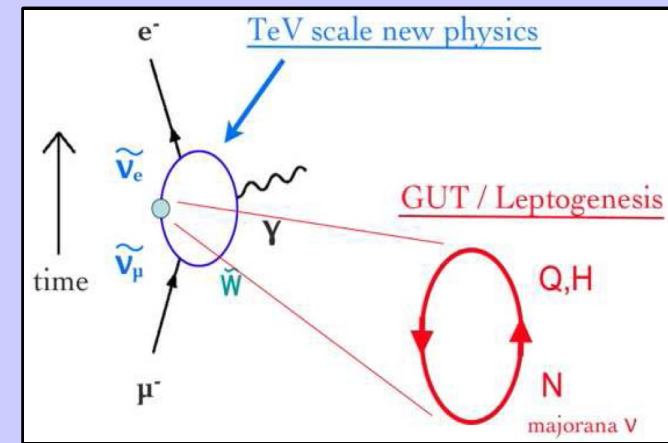
Problems of Seesaw Mechanism

- **Introduces high energy scale**
- **Right-handed neutrinos are singlets**
Couple only via small mixture with active neutrinos
- **Mechanism not testable with low energy observables**



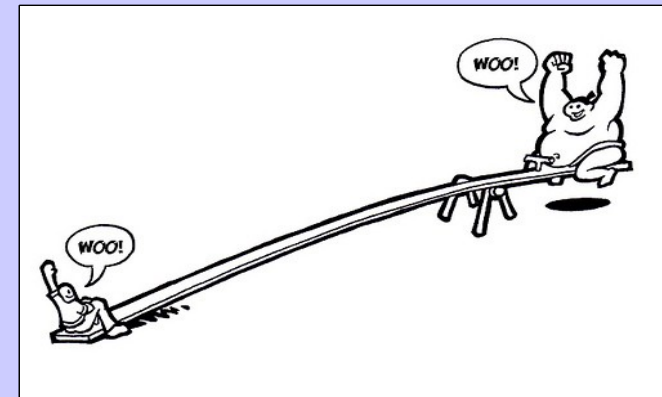
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- **Possible Solutions**
 - **SUSY Seesaw**
Testable LFV effects from sleptons



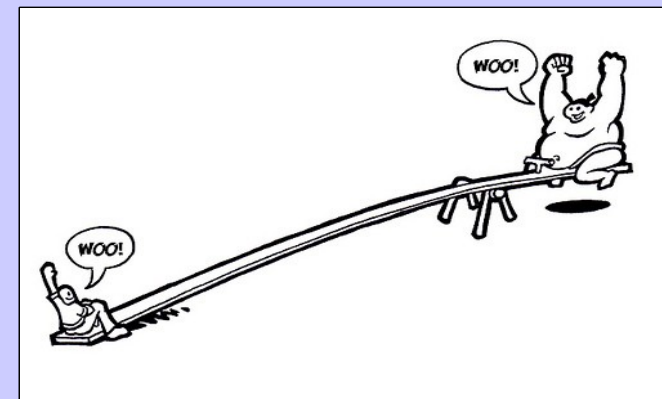
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LNV at low scale allows low mass of right-handed neutrinos



Problems of Seesaw Mechanism

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Testable LFV effects from sleptons
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LNV at low scale allows low mass of right-handed neutrinos
 - **Left-Right symmetric models**
Right-handed neutrinos couple with gauge strength to charged leptons



Minimal Left-Right Symmetrical Model

- Based on

$$SU(3) \times SU(2)_L \times SU(2)_R \times U(1)_{B-L}$$

Pati & Salam '74
Mohapatra & Senjanovic '75

- Higgs Sector:**

Bidoublet (EW Breaking) + Left-handed Triplet + Right-handed Triplet (Breaking Lepton Number + Parity + $SU(2)_R$)

- Generate $N_i + W_R + Z_R$ masses

$$M_{N_i} \approx M_{W_R} \approx M_{Z_R} \approx \langle \Delta_R \rangle \approx 0.5 - 5 \text{ TeV}$$

- General Seesaw II Mechanism

$$M_\nu = \begin{pmatrix} M_L & M_D \\ M_D^T & M_R \end{pmatrix},$$

- Charged current weak interactions

Neglect any
Left-Right mixing

$$J_{W_1}^{\mu-} = \frac{g_L}{\sqrt{2}} (\bar{\nu}_i U_{li}^{LL} + \bar{N}_i^c U_{li}^{LR}) \gamma^\mu \ell_L + \frac{g_R}{\sqrt{2}} \sin \zeta_W (\bar{\nu}_i U_{li}^{RL} + \bar{N}_i U_{li}^{RR}) \gamma^\mu \ell_R,$$

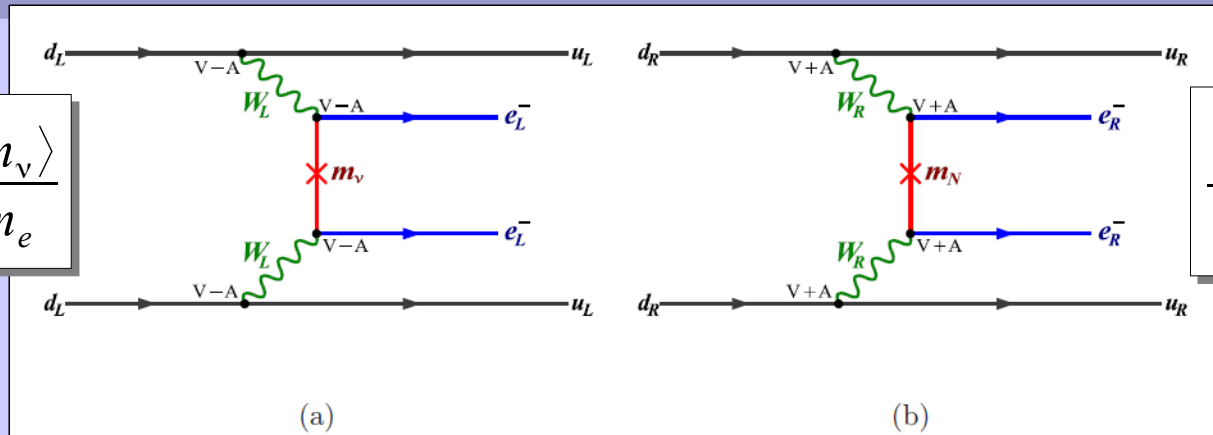
$$J_{W_2}^{\mu-} = -\frac{g_L}{\sqrt{2}} \sin \zeta_W (\bar{\nu}_i U_{li}^{LL} + \bar{N}_i U_{li}^{LR}) \gamma^\mu \ell_L + \frac{g_R}{\sqrt{2}} (\bar{N}_i U_{li}^{RR} + \bar{\nu}_i^c U_{li}^{RL}) \gamma^\mu \ell_R,$$

$$J_{W_L}^{\mu-} \approx \frac{g_L}{\sqrt{2}} U_{li} \bar{\nu}_i \gamma^\mu \ell_L,$$

$$J_{W_R}^{\mu-} \approx \frac{g_R}{\sqrt{2}} V_{li} \bar{N}_i \gamma^\mu \ell_R,$$

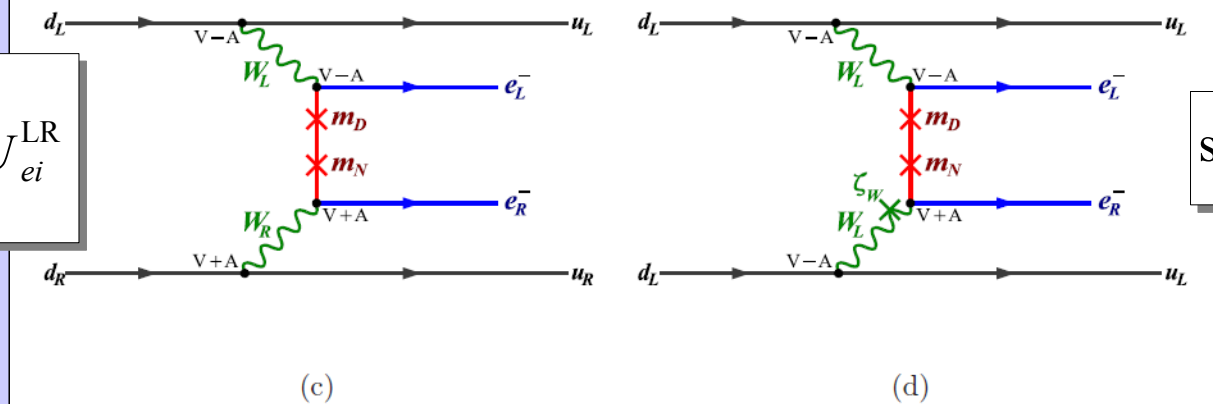
Neutrinoless Double Beta Decay in the LRSM

$$\sum_i (U_{ei}^{LL})^2 \frac{m_{\nu_i}}{m_e} = \frac{\langle m_\nu \rangle}{m_e}$$



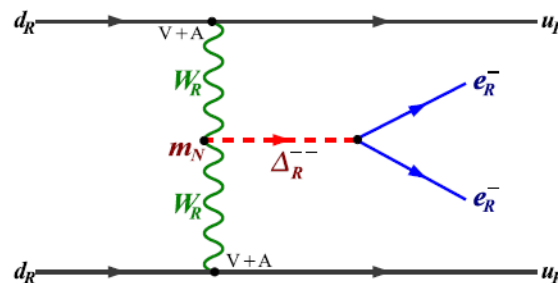
$$\frac{M_{W_L}^4}{M_{W_R}^4} \sum_i \frac{(U_{ei}^{RR})^2}{M_{N_i}}$$

$$\left(\frac{M_{W_L}}{M_{W_R}} \right)^2 \sum_i U_{ei}^{LL} U_{ei}^{LR}$$

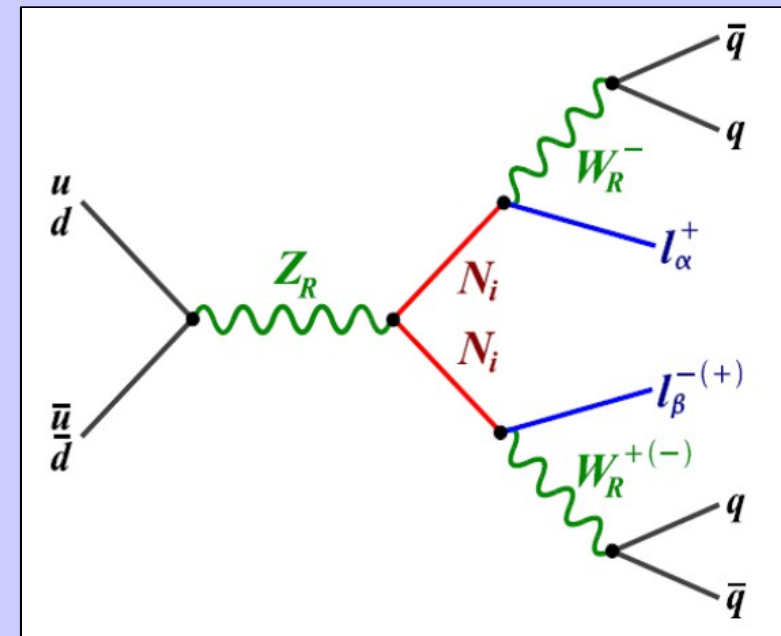
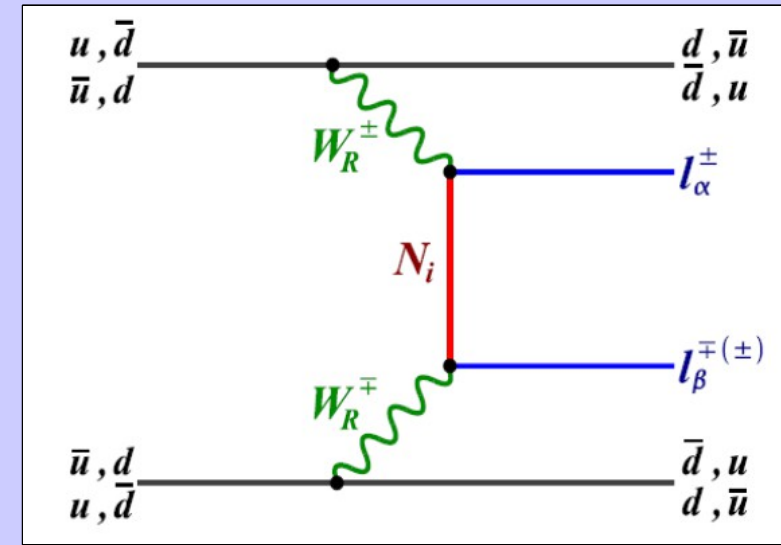
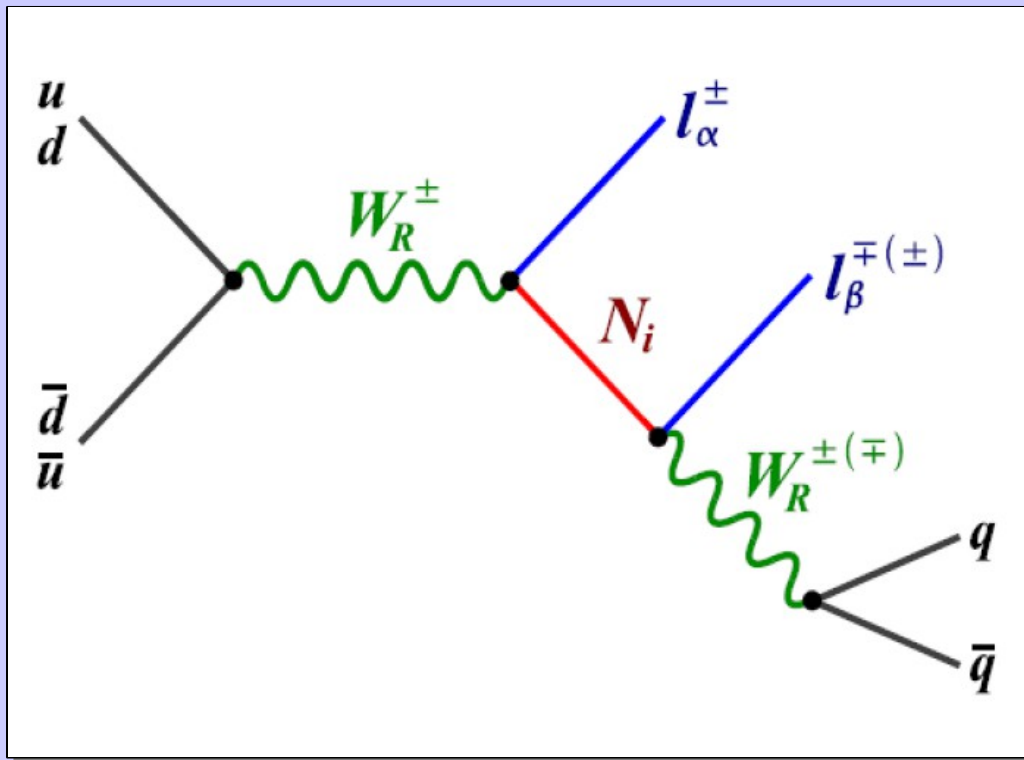


$$\sin^2 \zeta \sum_i U_{ei}^{LL} U_{ei}^{LR}$$

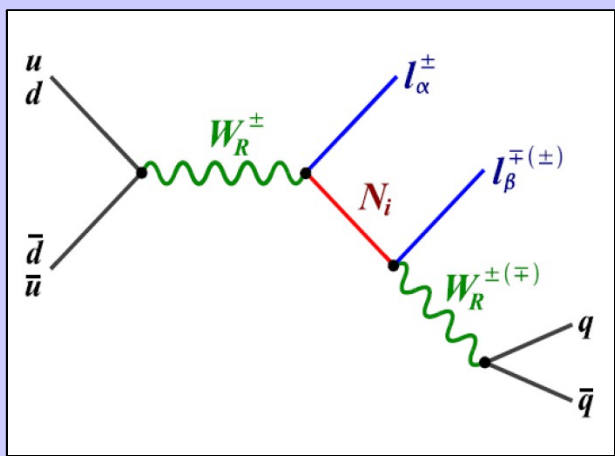
$$\frac{M_{W_L}^4}{M_{W_R}^4} \frac{m_p}{M_{\Delta_R^{--}}^2} \sum_i (U_{ei}^{RR})^2 M_{N_i}$$



Right-handed Neutrino Production at the LHC

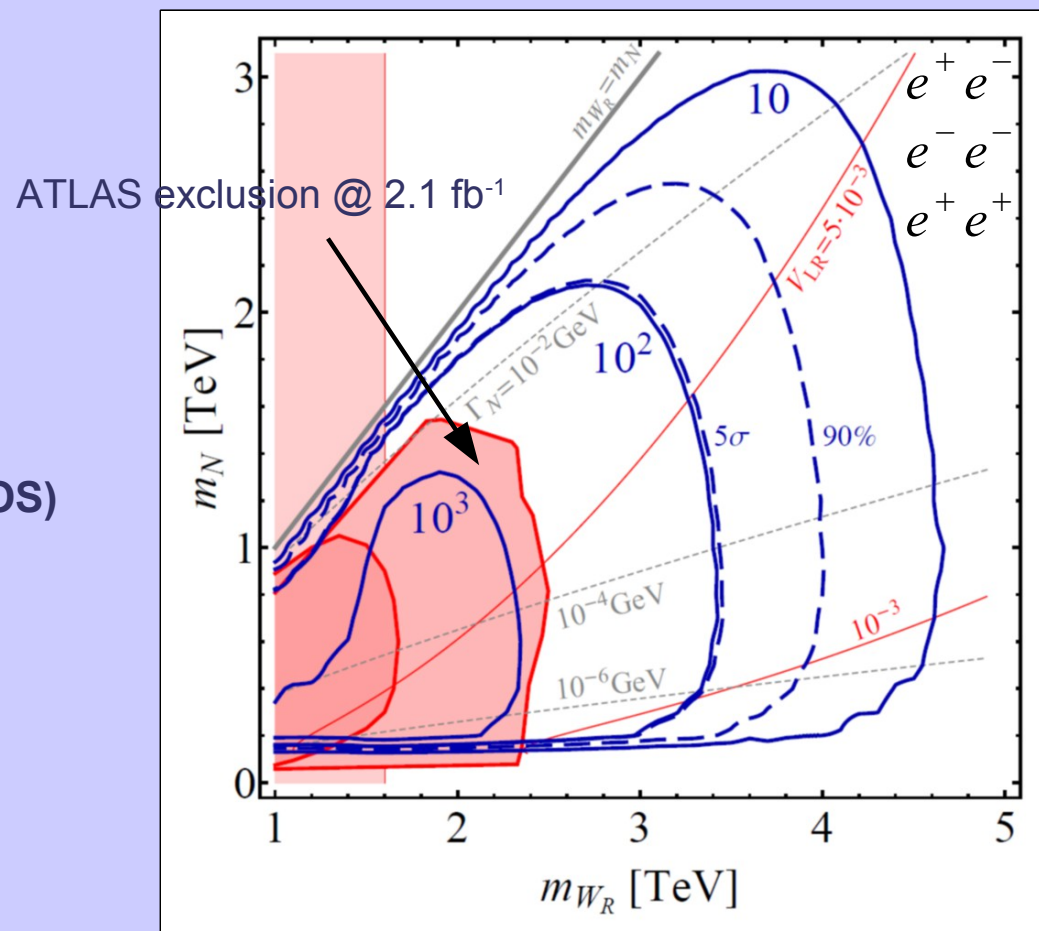


Single Neutrino Production



- Monte Carlo Simulation (PROTOS)
- Background $t\bar{t}$, Z + jets (Pythia, Alpgen)
- Fast Detector Simulation (AcerDET)
- Selection Criteria

number of jets	$N_j \geq 2$
number of isolated leptons	$N_\ell = 2$
invariant dilepton mass	$m_{\ell\ell} > 300 \text{ GeV}$
total invariant mass	$m_{\ell\ell jj} > 1.5 \text{ TeV}$



Opposite Sign + Same Sign Leptons
LHC reach @ 14 TeV, 30 fb⁻¹

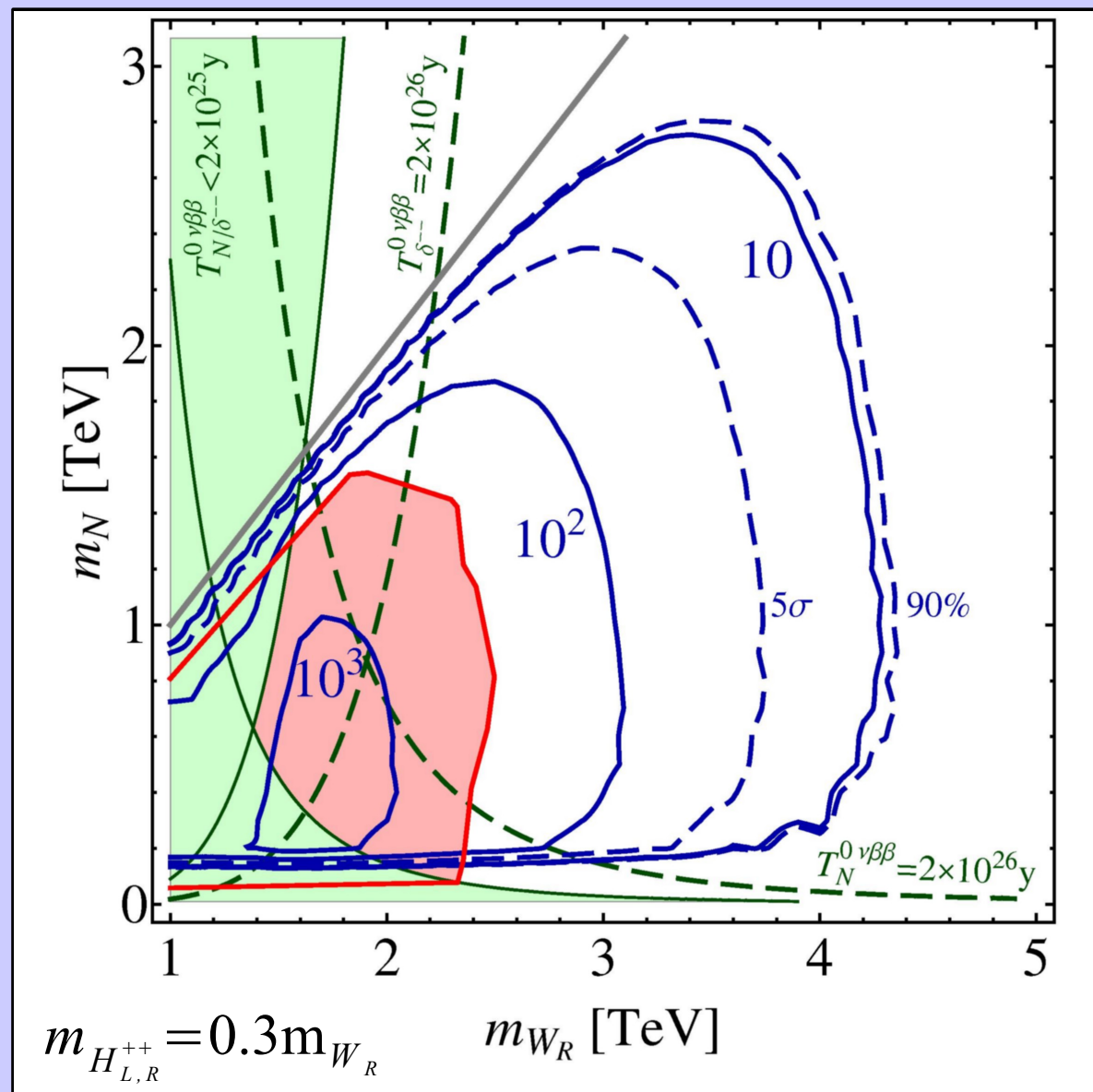
Comparison to $0\nu\beta\beta$

- Consider contributions to $0\nu\beta\beta$ from triplet Higgs

$$\frac{M_{W_L}^4}{M_{W_R}^4} \frac{m_p}{M_{\Delta_R^{--}}} \sum_i (U_{ei}^{RR})^2 M_{N_i}$$

- and heavy neutrinos

$$\frac{M_{W_L}^4}{M_{W_R}^4} \sum_i \frac{(U_{ei}^{RR})^2}{M_{N_i}}$$



LHC reach @ 14 TeV, 30 fb⁻¹

Conclusion

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Strong experimental program to probe absolute mass

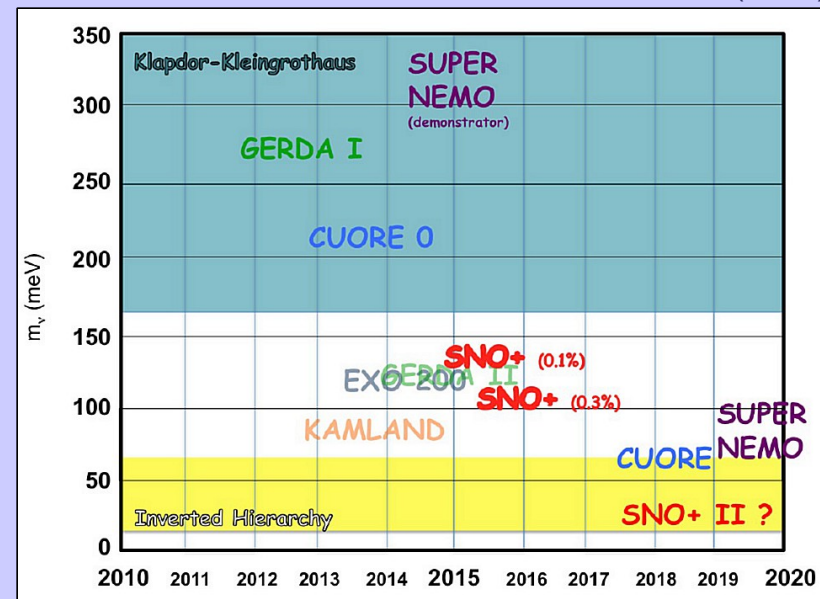
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New LNV physics at the EW scale
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- **Rich phenomenology in models of neutrino mass generation**

- Cosmological Observations
- Charged lepton flavour violation
- LFV and LNV processes at the LHC
- Connection to Leptogenesis?

Lefeuvre (2011)

