

SUSY Model Building for a 126 GeV Higgs

Moritz McGarrie

I .The Higgs mass 126 GeV

$$m_h^2 \simeq m_z^2 \cos^2(2\beta) + \frac{3}{(4\pi)^2} \frac{m_t^4}{v_{ew}^2} \left[\log \frac{m_{\tilde{t}}^2}{m_t^2} + \frac{X_t^2}{m_{\tilde{t}}^2} \left(1 - \frac{X_t^2}{12m_{\tilde{t}}^2} \right) \right]$$

$$X_t = A_t - \mu \cot \beta$$

$$126^2 = 91^2 + 81^2$$

- Radiative corrections are same order as tree level piece
- Fine tuning effects all models: R-S, composite Higgs, little Higgs, SUSY etc
- corrections run logarithmically in SUSY
- MSSM case implies either heavy stops or large $X_t = A_t + \dots$

Fine tuning: little hierarchy problem

$$m_z^2 = -2(m_{H_u}^2 + |\mu|^2) + \dots$$

For a natural cancellation these should be of the same order

Light Higgsino and Wino

Light stops

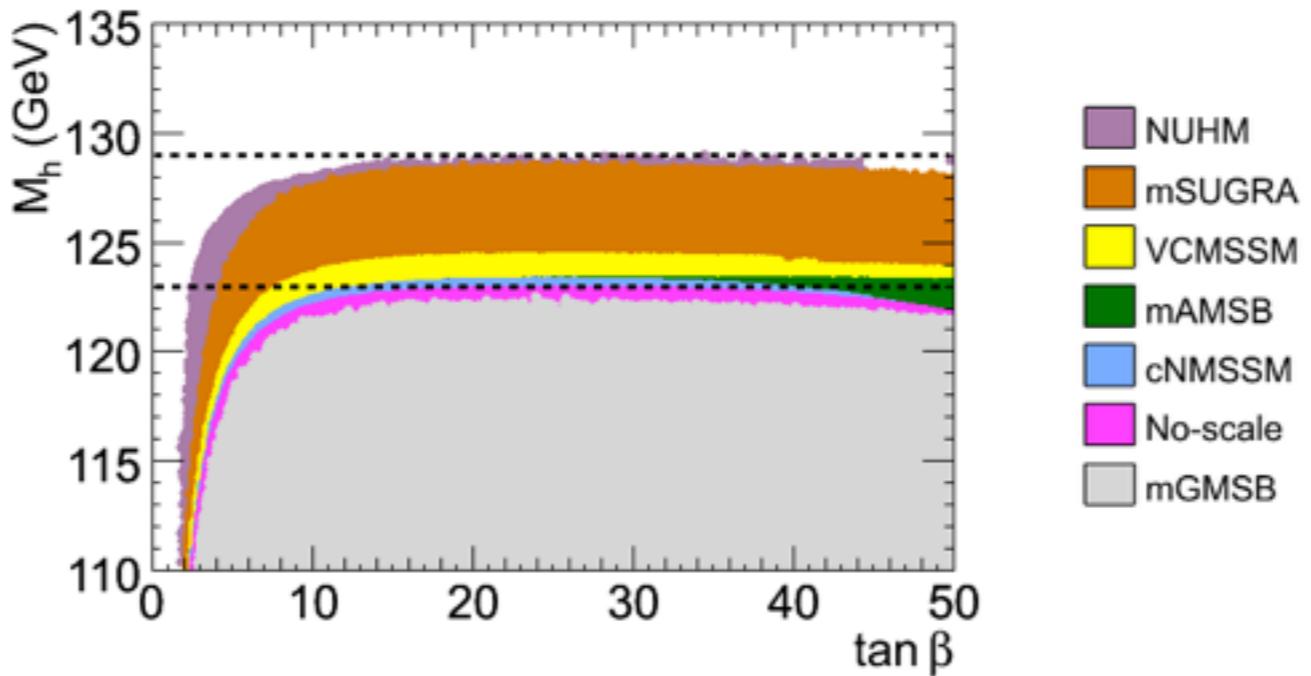


$$\delta m_{H_u}^2 \sim -\frac{3y_t^2 m_{\tilde{t}}^2}{4\pi^2} \text{Log}\left(\frac{\Lambda}{m_{\tilde{t}}}\right)$$

Light(ish) Gluino

$$\delta m_{\tilde{t}}^2 = \frac{8\alpha_s M_3^2}{3\pi} \text{Log}\left(\frac{\Lambda}{M_3}\right)$$

How do benchmark (minimal) models fair?



I207.1348
Arbey, Battaglia, Djouadi, Mahmoudi

stops less than 2 TeV

The point is that these models all work, but only if the spectrum is very heavy e.g. 5-10 TeV stops: *unnatural and unobservable at LHC14.*

if you give up on naturalness, is there any reason to see anything at the LHC?

Reasons to (still) believe in SUSY

- Solves big hierarchy problem M_{EW}/M_{Pl}
- Improves GUT unification (and so too, stability of Proton)
- necessary for string theory (unification of all forces)
- radiatively induced EWSB in MSSM-like models.
- Believable dark matter candidate(s)
- It is consistent with low energy EW observables
- Improves fit to electric dipole moment, muon magnetic moment etc.
- MSSM bounds Higgs < 135 GeV

what are our choices?

The trivial / *Immoral option*

Split SUSY, some landscape (multiverse) gives an explanation for Y_u, Λ_{cc}, m_h^2 (maybe more?)

The hard / ugly / *unbelievable option*

Generate large At (vanishing in mGMSB, uncalculable in SUGRA and has other problems.)

The innovative *option*

or get creative....

In this talk I am going to give 2 examples of models that improve naturalness

Non decoupled D-terms

(Aoife Bharucha, Andreas Goudelis & MM) [I310.4500](#)

ISS & Flavour Gauge Messengers

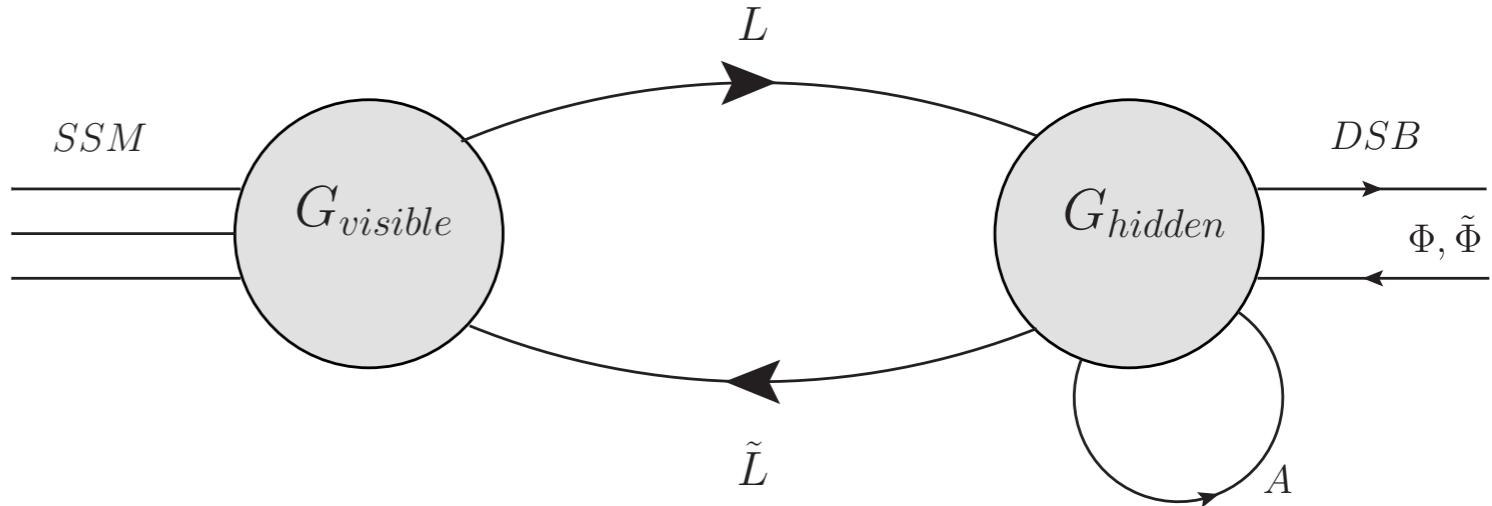
(F.Bruemmer, A.Weiler & MM) [I312.0935](#)

(S.Abel & MM) [I404.1318](#)

(Both presented at SUSY 2013)

A quiver model: motivation

- non decoupled D-terms: lifts the Higgs
(sometimes substantially)
Batra, Delgado, Kaplan, Tait 0309149
- extra adjoints of $SU(2), SU(3)$: lifts the Higgs
- More natural than NMSSM?
- embeds into magnetic SQCD
- Deconstructs an extra dimension
- “Split families” Batra, Kaplan, Tait, Delgado 0404251/0409073



$$m_h^2 \simeq m_z^2 \cos^2(2\beta) + \frac{3}{(4\pi)^2} \frac{m_t^4}{v_{ew}^2} \left[\log \frac{m_{\tilde{t}}^2}{m_t^2} + \frac{X_t^2}{m_{\tilde{t}}^2} \left(1 - \frac{X_t^2}{12m_{\tilde{t}}^2} \right) \right]$$

$$\Delta = \left(\frac{g_A^2}{g_B^2} \right) \frac{2m_L^2}{m_v^2 + 2m_L^2}$$

Related works:

Csaki, Erlich, Grojean, Kribs 0106044
 Medina, Shah, Wagner 0904.1625
 “GGM and Deconstruction”
 M.M. 1009.0012 and 1101.5158

Auzzi, Giveon, Gudnason, Shacham
 1009.1714
 1011.1664

easyDiracGauginos

S.Abel, M. Goodsell

Bharucha, Goudelis, M.M. 1310.4500

$$\begin{aligned} \delta \mathcal{L} = & -g_1^2 \Delta_1 (H_u^\dagger H_u - H_d^\dagger H_d)^2 \\ & - g_2^2 \Delta_2 \sum_a (H_u^\dagger \sigma^a H_u + H_d^\dagger \sigma^a H_d)^2 \end{aligned}$$

$$m_z^2 \rightarrow m_z^2 + \left(\frac{g_1^2 \Delta_1 + g_2^2 \Delta_2}{2} \right) v_{ew}^2$$

D-terms Vs NMSSM

$$m_h^2 \simeq m_z^2 \cos^2(2\beta) + \frac{3}{(4\pi)^2} \frac{m_t^4}{v_{ew}^2} \left[\log \frac{m_{\tilde{t}}^2}{m_t^2} + \frac{X_t^2}{m_t^2} \left(1 - \frac{X_t^2}{12m_{\tilde{t}}^2} \right) \right]$$

$$\Delta = \left(\frac{g_A^2}{g_B^2} \right) \frac{2m_L^2}{m_v^2 + 2m_L^2}$$

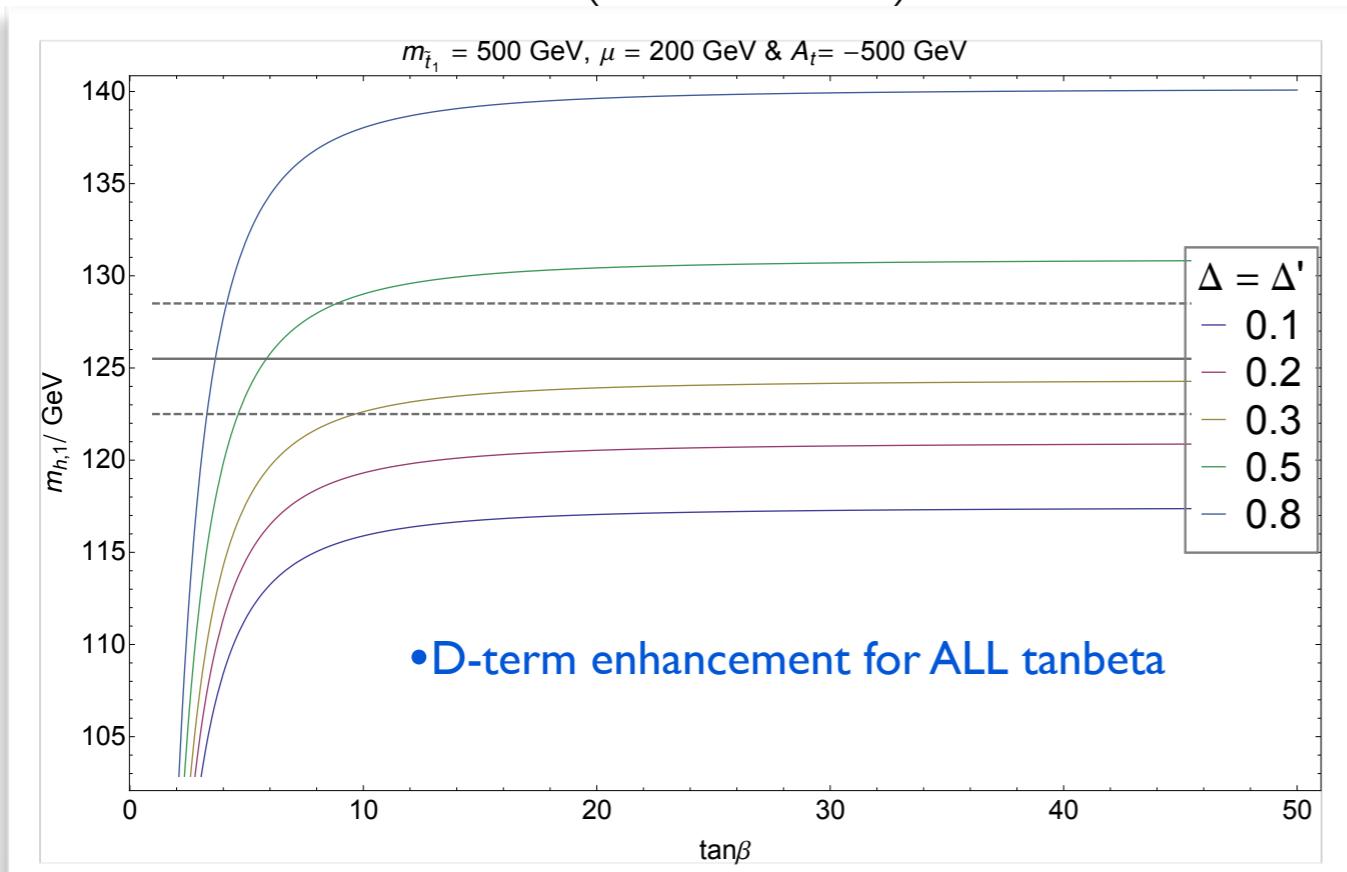
$$\begin{aligned} \delta\mathcal{L} = & -g_1^2 \Delta_1 (H_u^\dagger H_u - H_d^\dagger H_d)^2 \\ & - g_2^2 \Delta_2 \sum_a (H_u^\dagger \sigma^a H_u + H_d^\dagger \sigma^a H_d)^2 \end{aligned}$$

$$m_z^2 \rightarrow m_z^2 + \left(\frac{g_1^2 \Delta_1 + g_2^2 \Delta_2}{2} \right) v_{ew}^2$$

$$W_{NMSSM} \supset \lambda S H_u H_d$$

$$V(\phi' s) \supset \lambda^2 |H_u H_d|^2$$

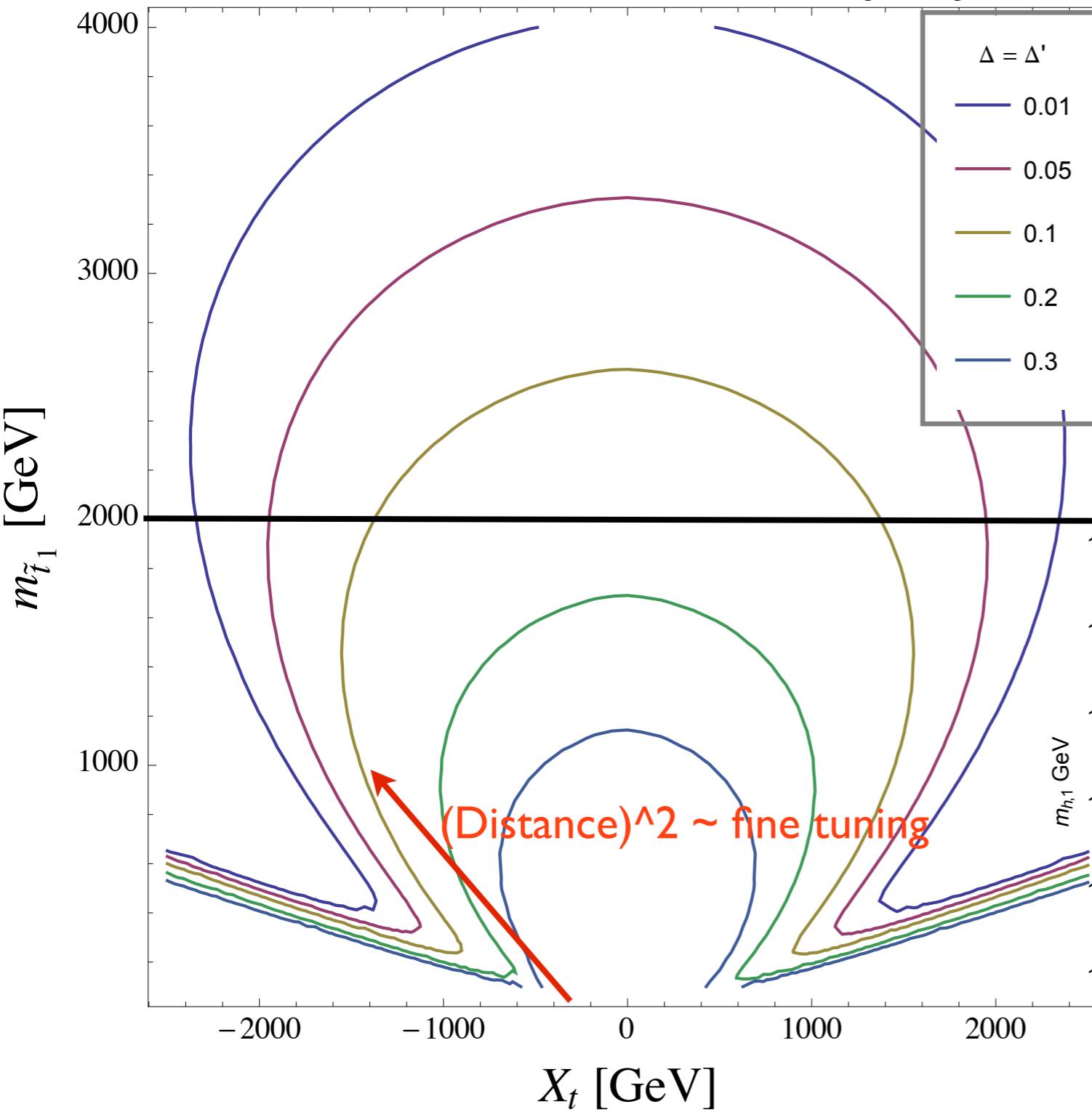
$$m_{h_0}^2 = m_z^2 \cos(2\beta) + \lambda^2 v_{ew}^2 \sin(2\beta)$$



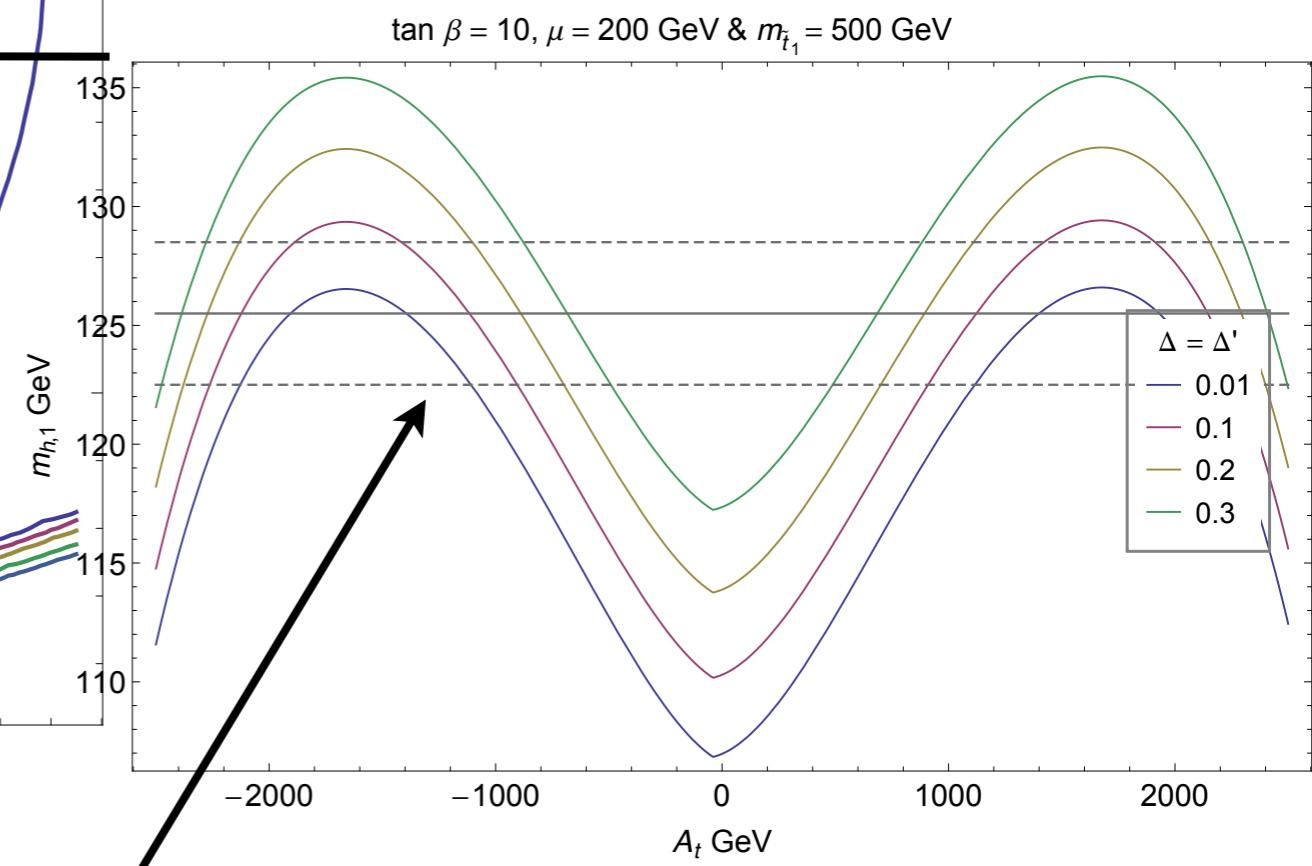
•F-term enhancement only for small tanbeta

$m_h = 125.5 \text{ GeV}$, $\tan \beta = 10$, $\mu = 200 \text{ GeV}$ & $m_{Q_3} = m_{U_3}$

M.M. Moortgat-Pick & S. Porto, G.



Sub 2 TeV stops
for $\Delta \geq 0.1$



Combination of medium size A_t and Deltas can do the trick! (GMSB back in the game?)

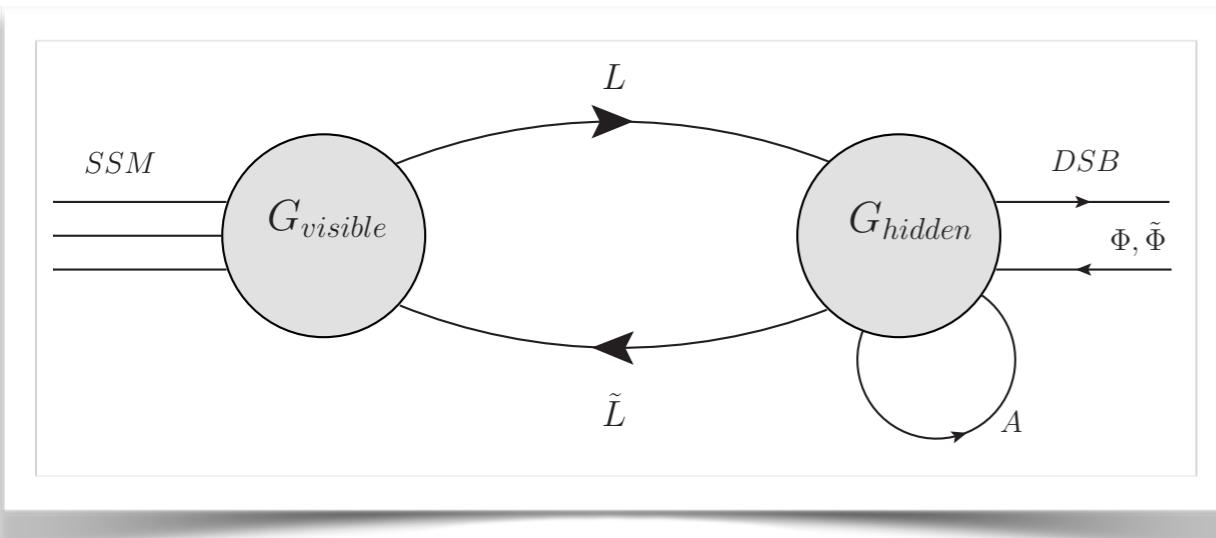
So far most studies of D-terms have been at tree-level and bottom-up

A meta model

Bharucha,Goudelis,M.M. | 310.4500

SF	Spin 0	Spin $\frac{1}{2}$	Generations	$(U(1)_A, SU(2)_B, SU(3)_c, U(1)_B, SU(2)_A)$	R-Parity
\hat{q}	\tilde{q}	q	3	$(\frac{1}{6}, 1, \mathbf{3}, 0, \mathbf{2})$	-1
\hat{l}	\tilde{l}	l	3	$(-\frac{1}{2}, 1, \mathbf{1}, 0, \mathbf{2})$	-1
\hat{H}_d	H_d	\tilde{H}_d	1	$(-\frac{1}{2}, 1, \mathbf{1}, 0, \mathbf{2})$	+1
\hat{H}_u	H_u	\tilde{H}_u	1	$(\frac{1}{2}, 1, \mathbf{1}, 0, \mathbf{2})$	+1
\hat{d}	d_R^*	d_R^*	3	$(\frac{1}{3}, 1, \overline{\mathbf{3}}, 0, 1)$	-1
\hat{u}	\tilde{u}_R^*	u_R^*	3	$(-\frac{2}{3}, 1, \overline{\mathbf{3}}, 0, 1)$	-1
\hat{e}	\tilde{e}_R^*	e_R^*	3	$(1, 1, 1, 0, 1)$	-1
\hat{L}	L	ψ_L	1	$(-\frac{1}{2}, \overline{\mathbf{2}}, 1, \frac{1}{2}, \mathbf{2})$	+1
$\hat{\tilde{L}}$	\tilde{L}	$\psi_{\tilde{L}}$	1	$(\frac{1}{2}, \mathbf{2}, 1, -\frac{1}{2}, \overline{\mathbf{2}})$	+1
\hat{K}	K	ψ_K	1	$(0, 1, 1, 0, 1)$	+1
\hat{A}	A	ψ_A	1	$(0, 3, 1, 0, 1)$	+1

Table 2. Matter fields of the model.

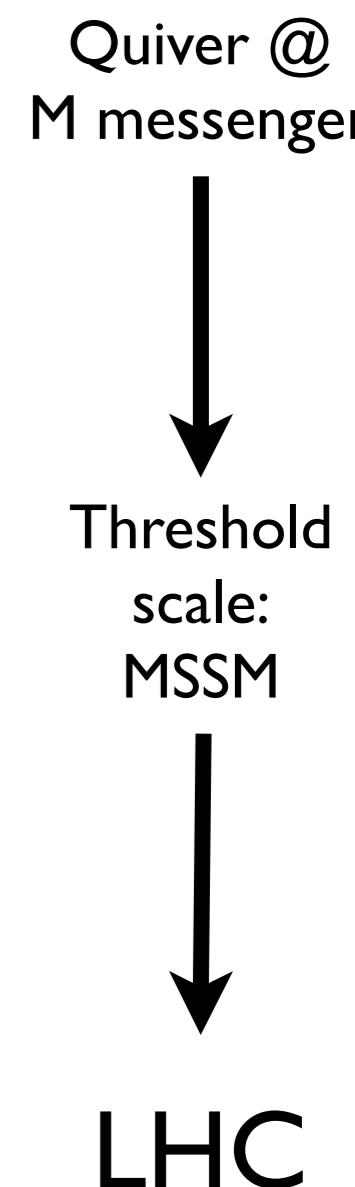


$$W_{\text{SSM}} = Y_u \hat{u} \hat{q} \hat{H}_u - Y_d \hat{d} \hat{q} \hat{H}_d - Y_e \hat{e} \hat{l} \hat{H}_d + \mu \hat{H}_u \hat{H}_d$$

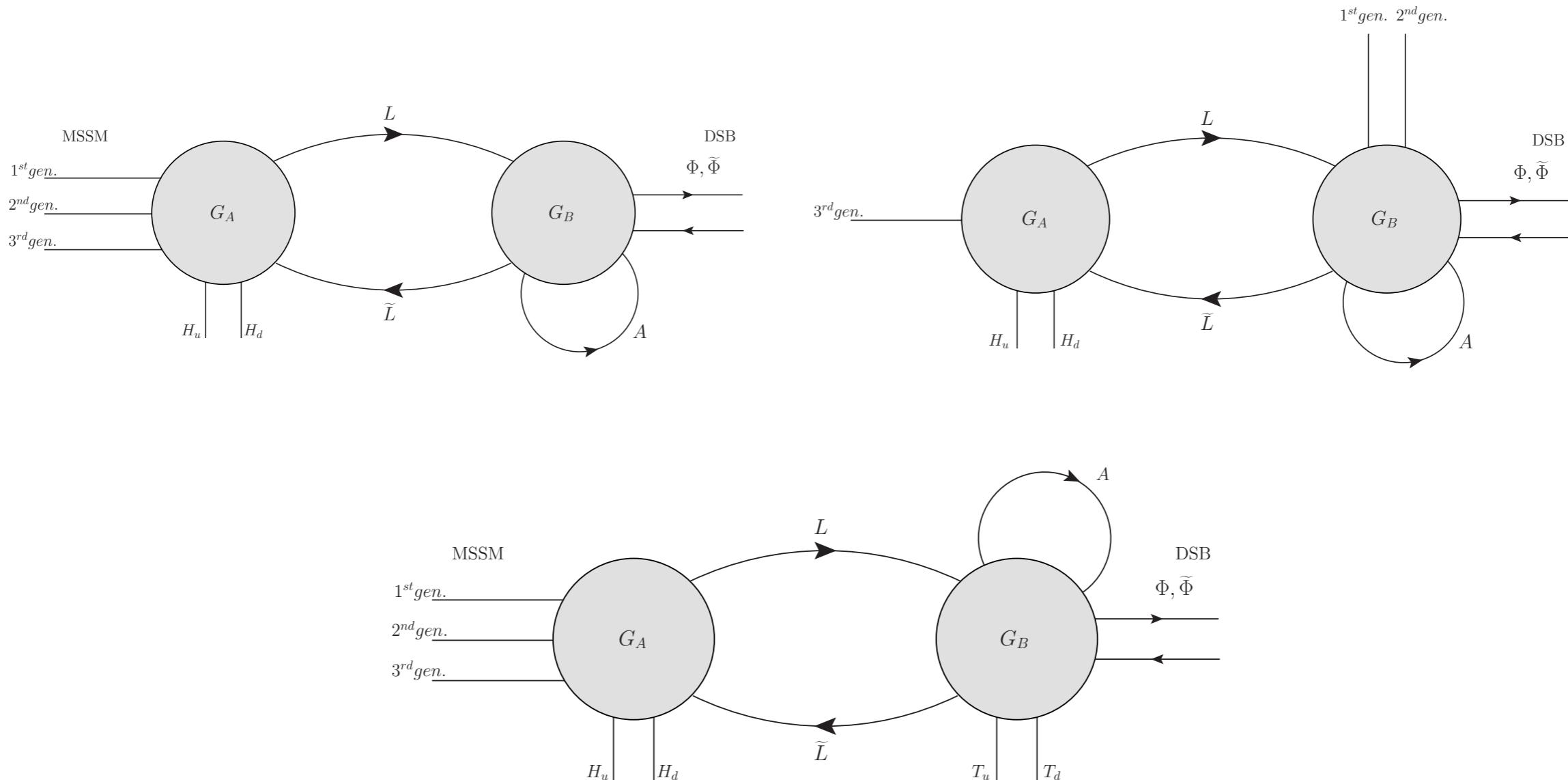
$$W_{\text{Quiver}} = \frac{Y_K}{2} \hat{K} (\hat{L} \hat{\tilde{L}} - V^2) + Y_A \hat{L} \hat{A} \hat{\tilde{L}}$$

The most sophisticated model so far implemented into a spectrum generator (SARAH/SPHENO)
 A meta-model i.e. *independent of the type of supersymmetry breaking:*
 AMSB, mSUGRA, GMSB, phenomenological, other?

Building a taylor made spectrum generator!

- We used SARAH mathematica package: “a spectrum generator generator” to write our own spectrum generator.
 - We implemented 5 gauge groups with full 2-loop RGE’s and one loop self energies (soon 6 and 9 gauge groups!).
 - Higgsing, and breaking to the diagonal 4 gauge groups, including all mixing matrices and assignment of Goldstones, Ghosts, RGEs of vevs, and Bmu at 2 loop.
 - All 3 and 4 vertices of all fields computed, and self energies.
 - All anomalous dimensions, tadpoles and running of all *additional soft terms and additional Yukawas, at 2 loop level.*
 - finite shifts and threshold corrections also accounted for.
 - Can talk to FeynArts, FormCalc, CalcHep, HiggsSignals, HiggsBounds, WHIZARD, micrOMEGAS, Vevacious and more.
- 

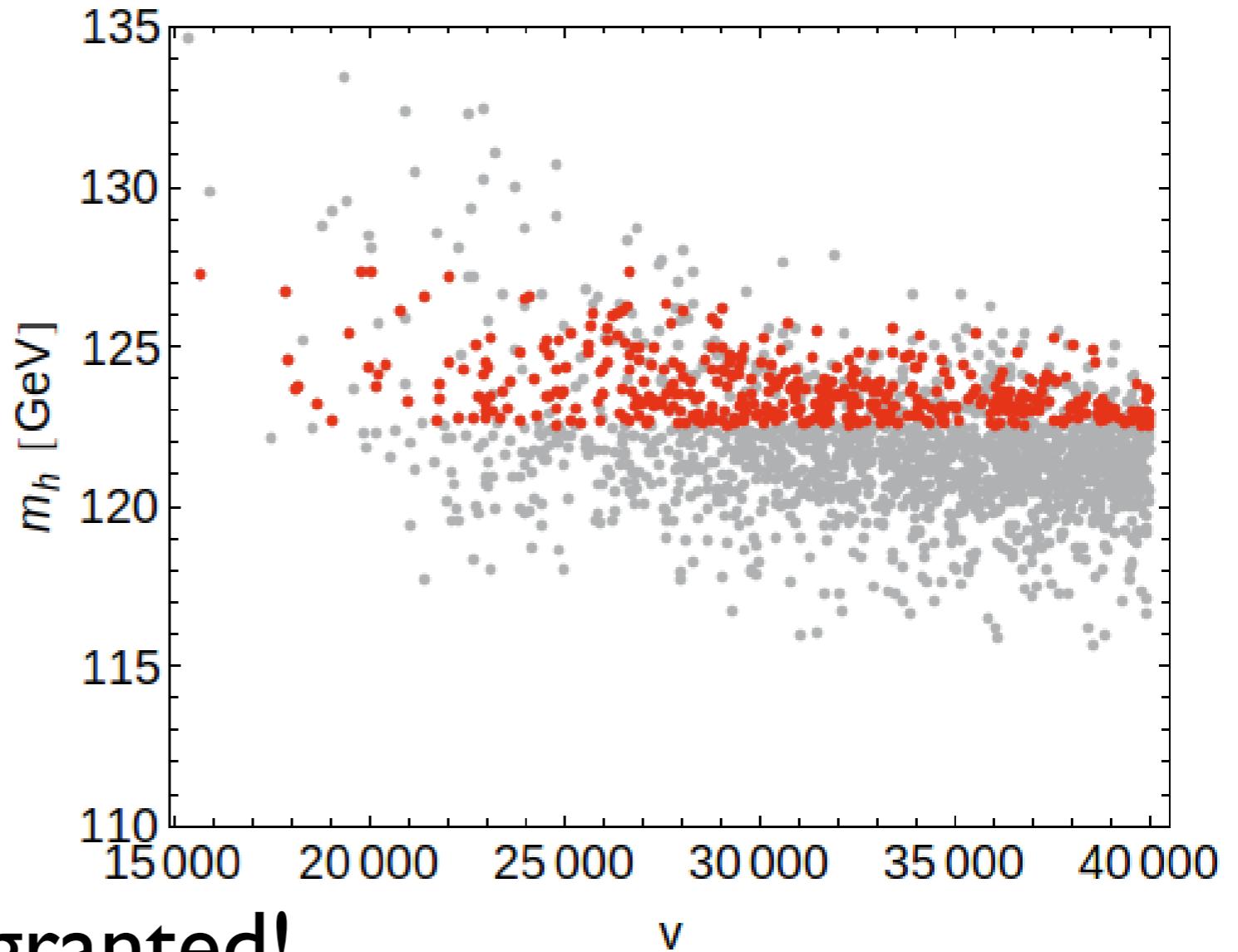
The Quiver Variations



Wish list

- $v < 10\text{TeV}$
- $m_L^2 > m_v^2$

We assumed GMSB
for soft terms!

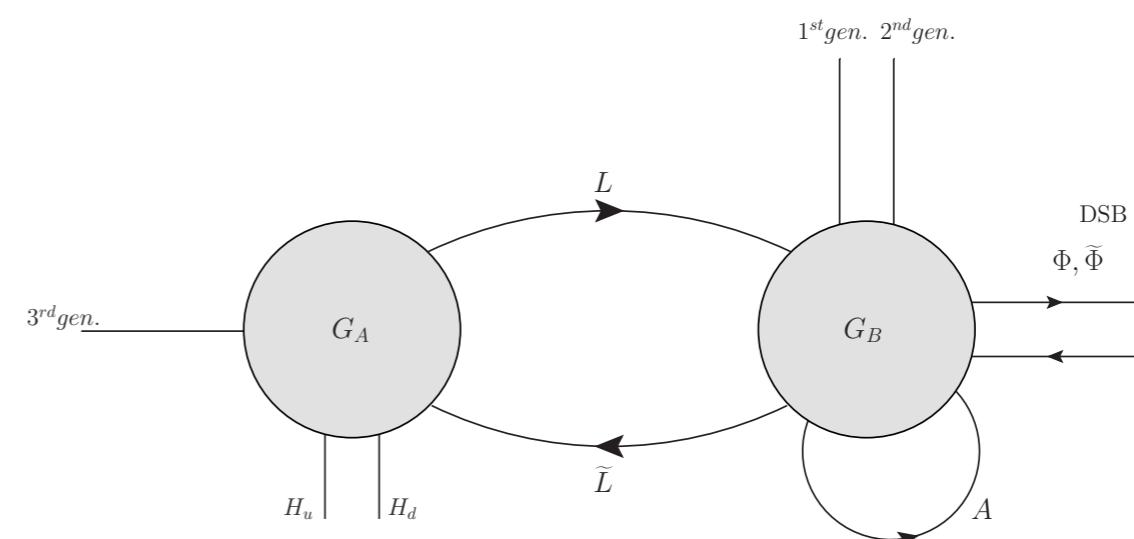
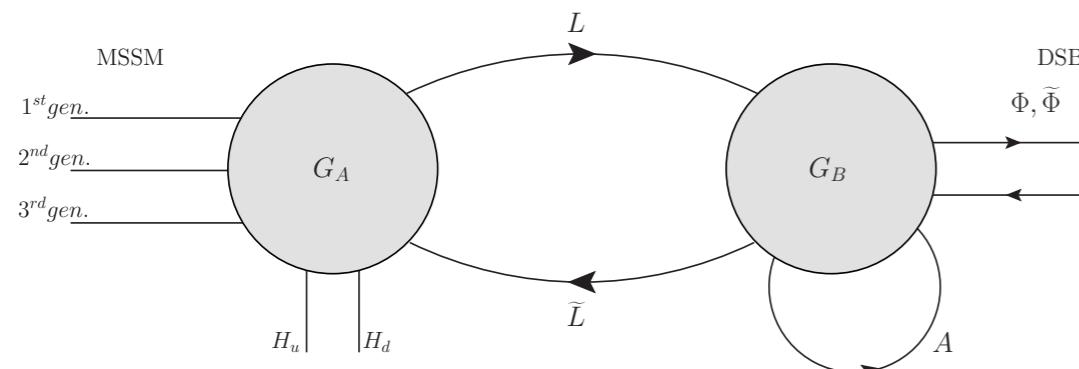
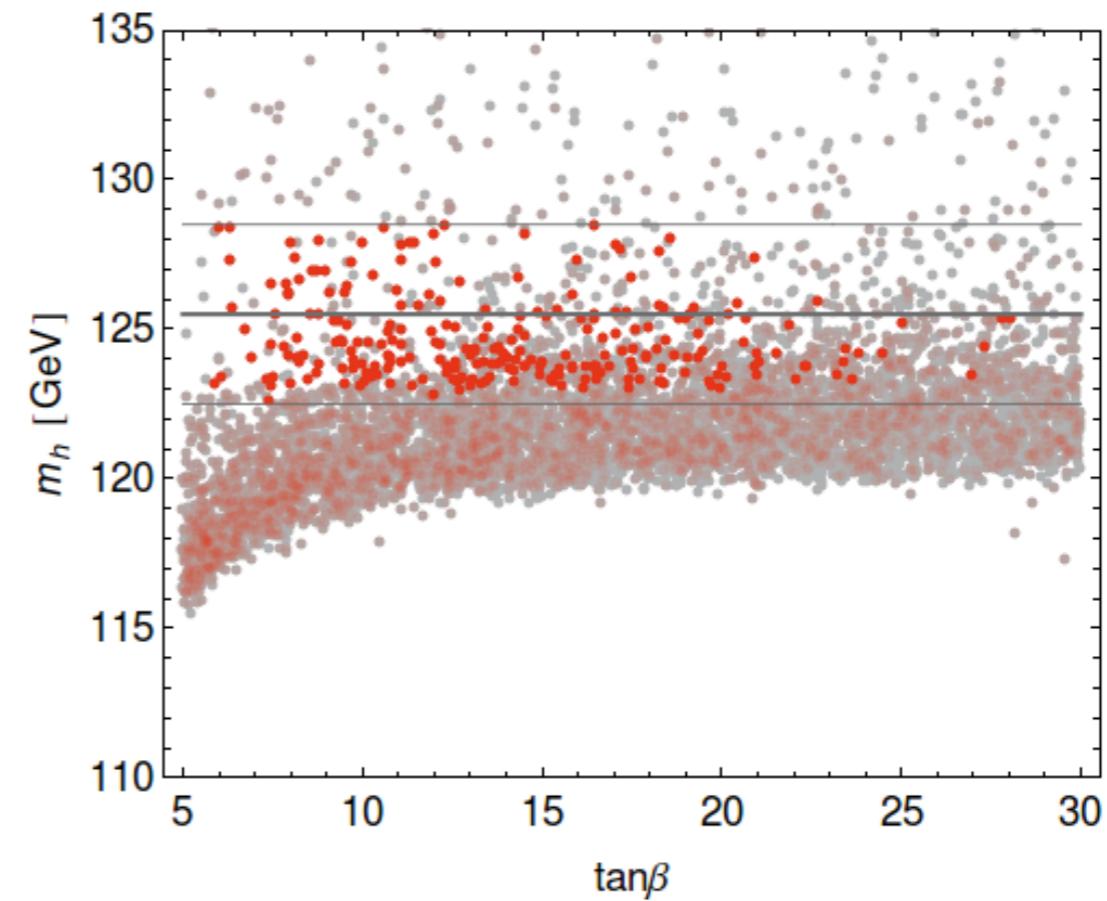
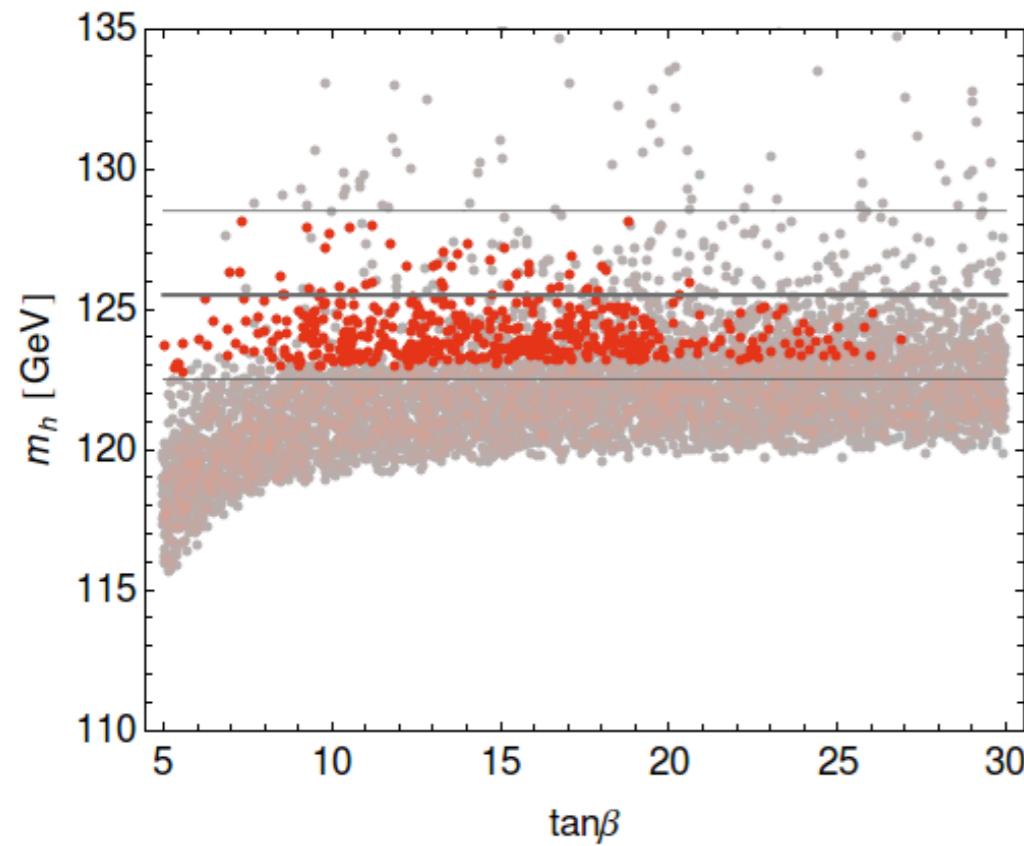


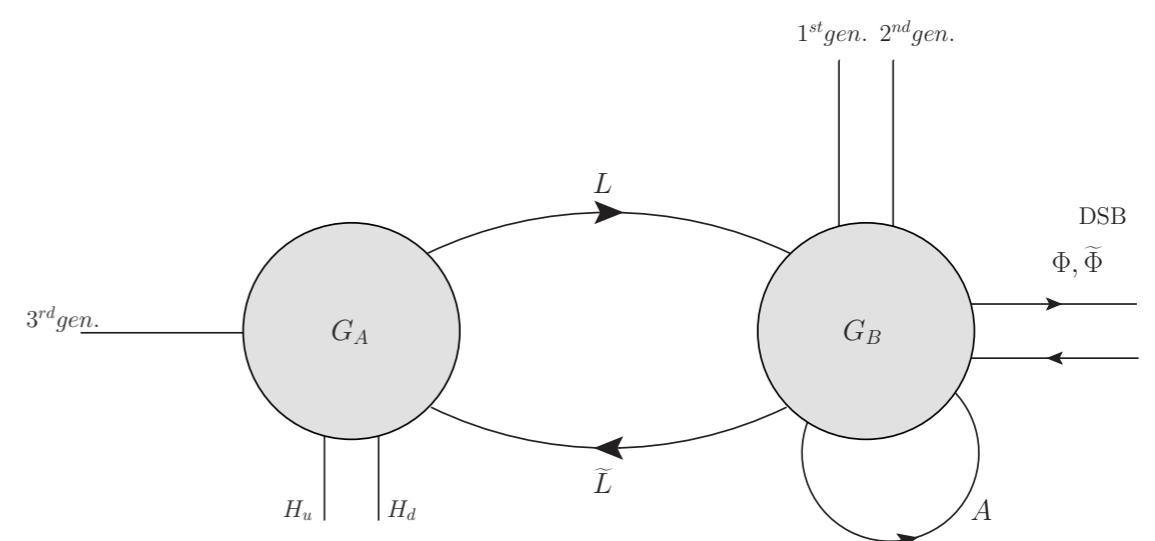
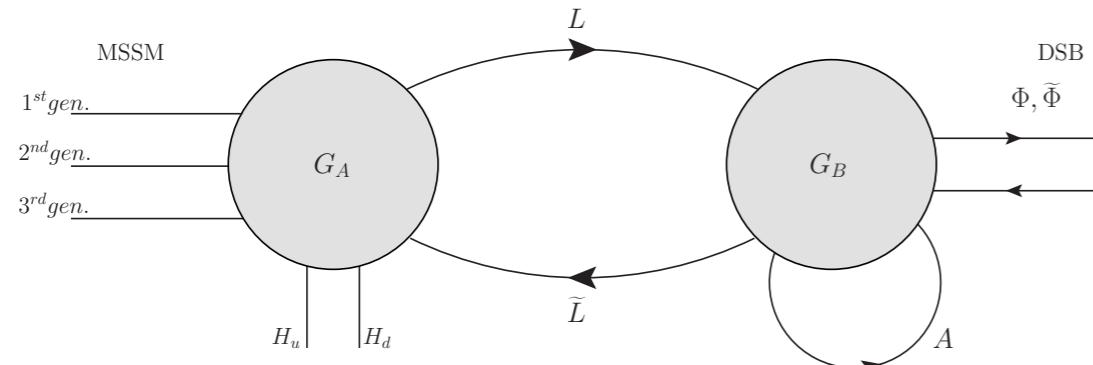
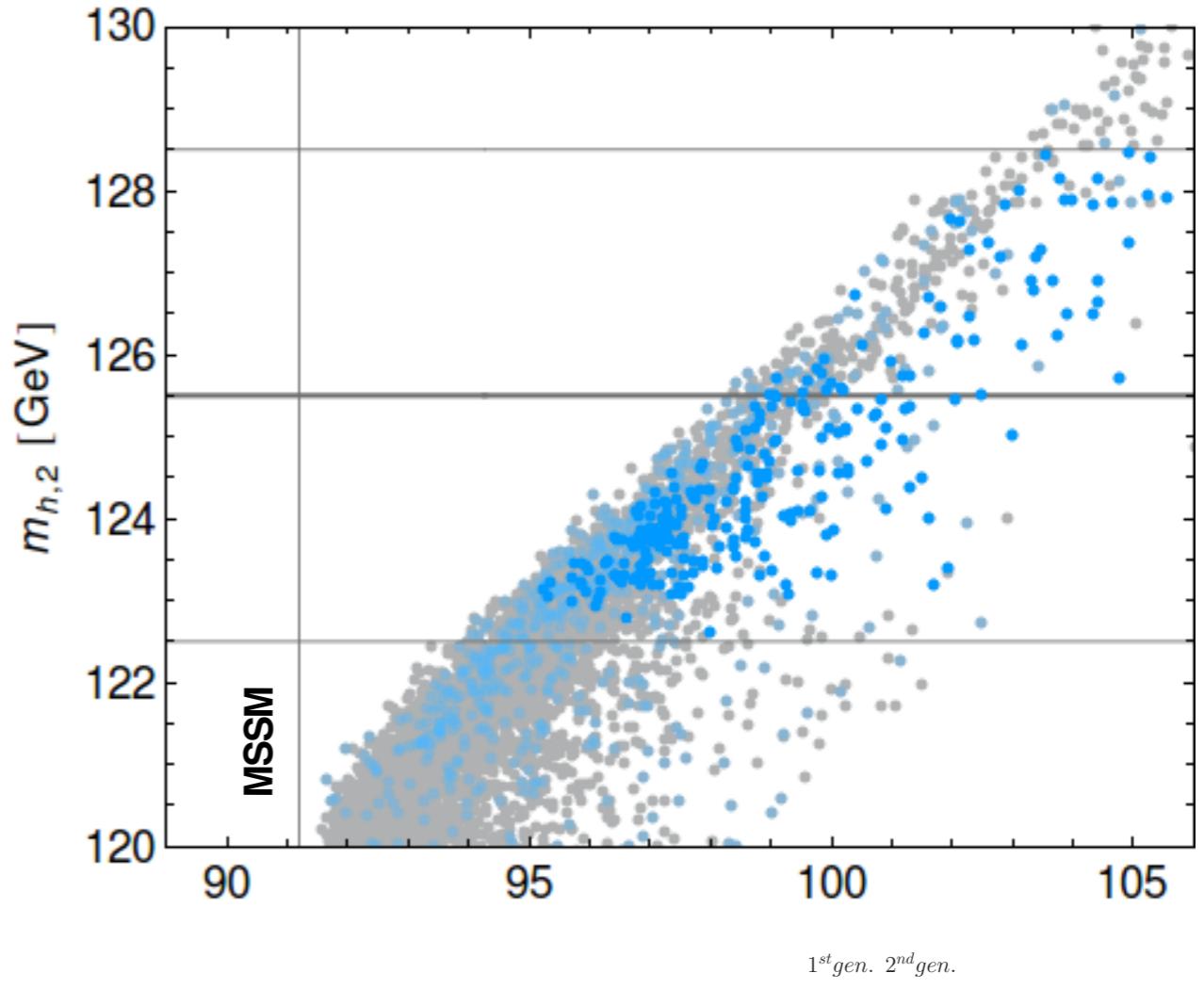
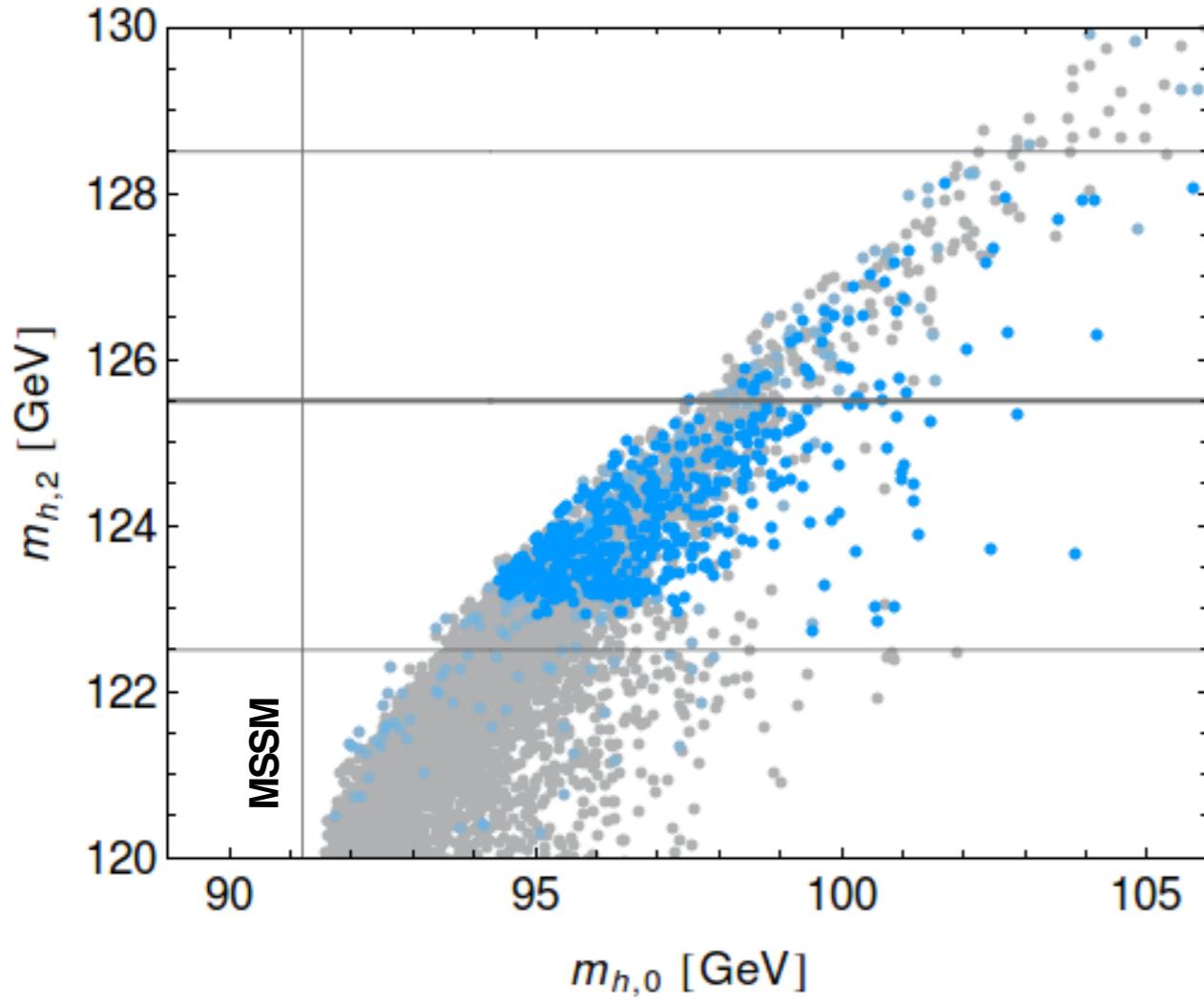
Our wishes were not granted!

we had to enlarge m_L^2 to
make it work

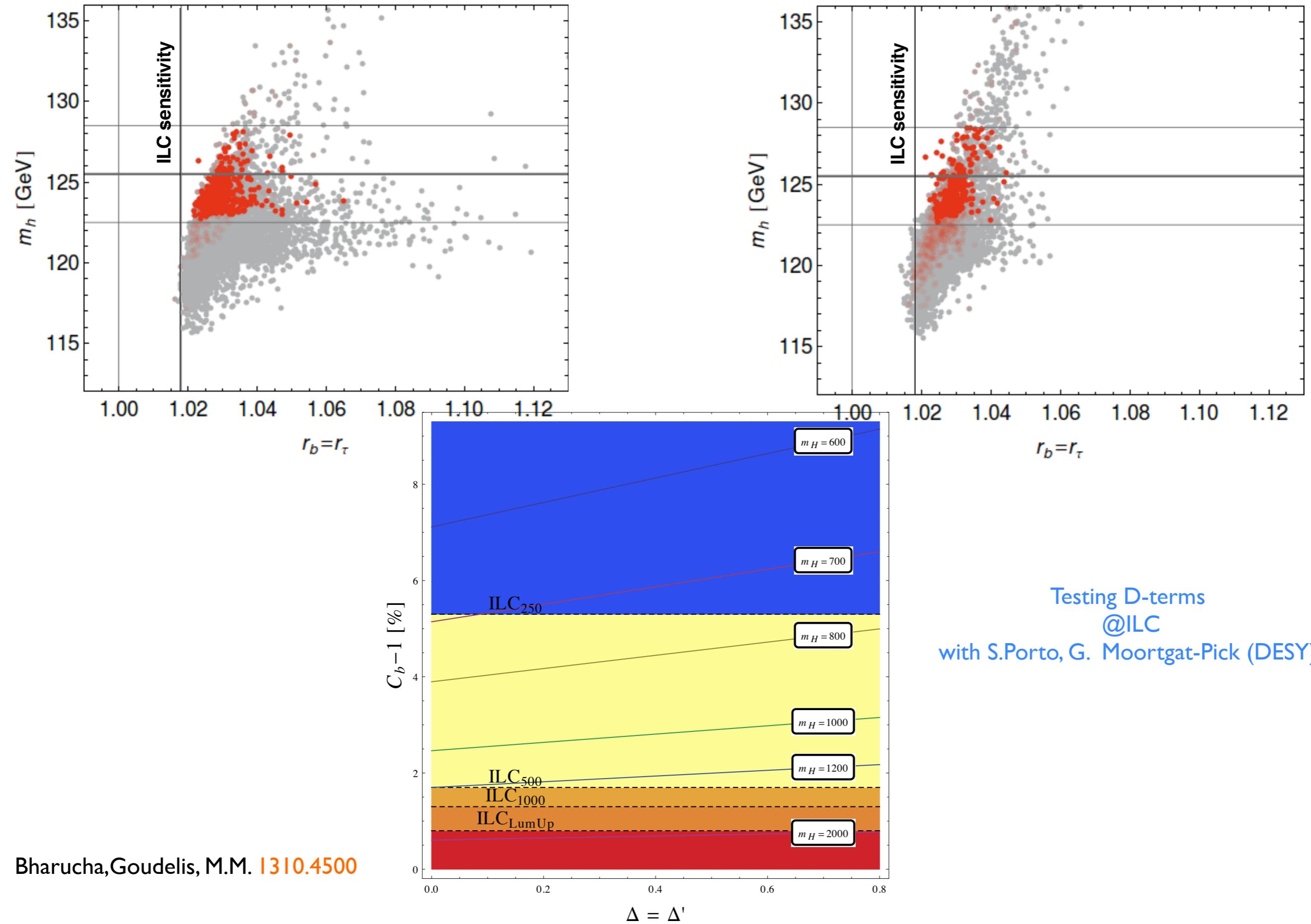
But perhaps this can be improved by
considering
different SUSY breaking scenarios

Ex. Applying mSUGRA here may get v down

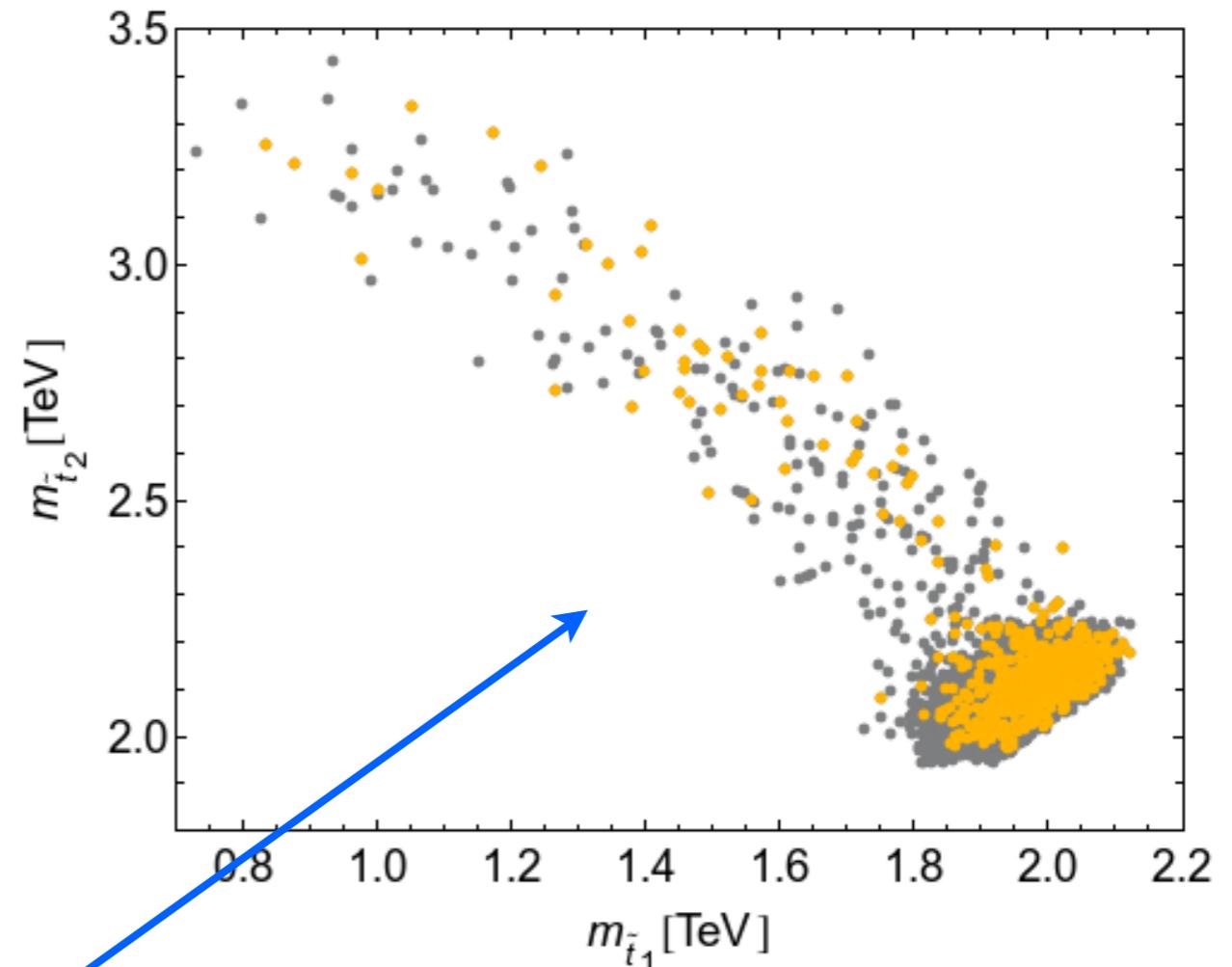
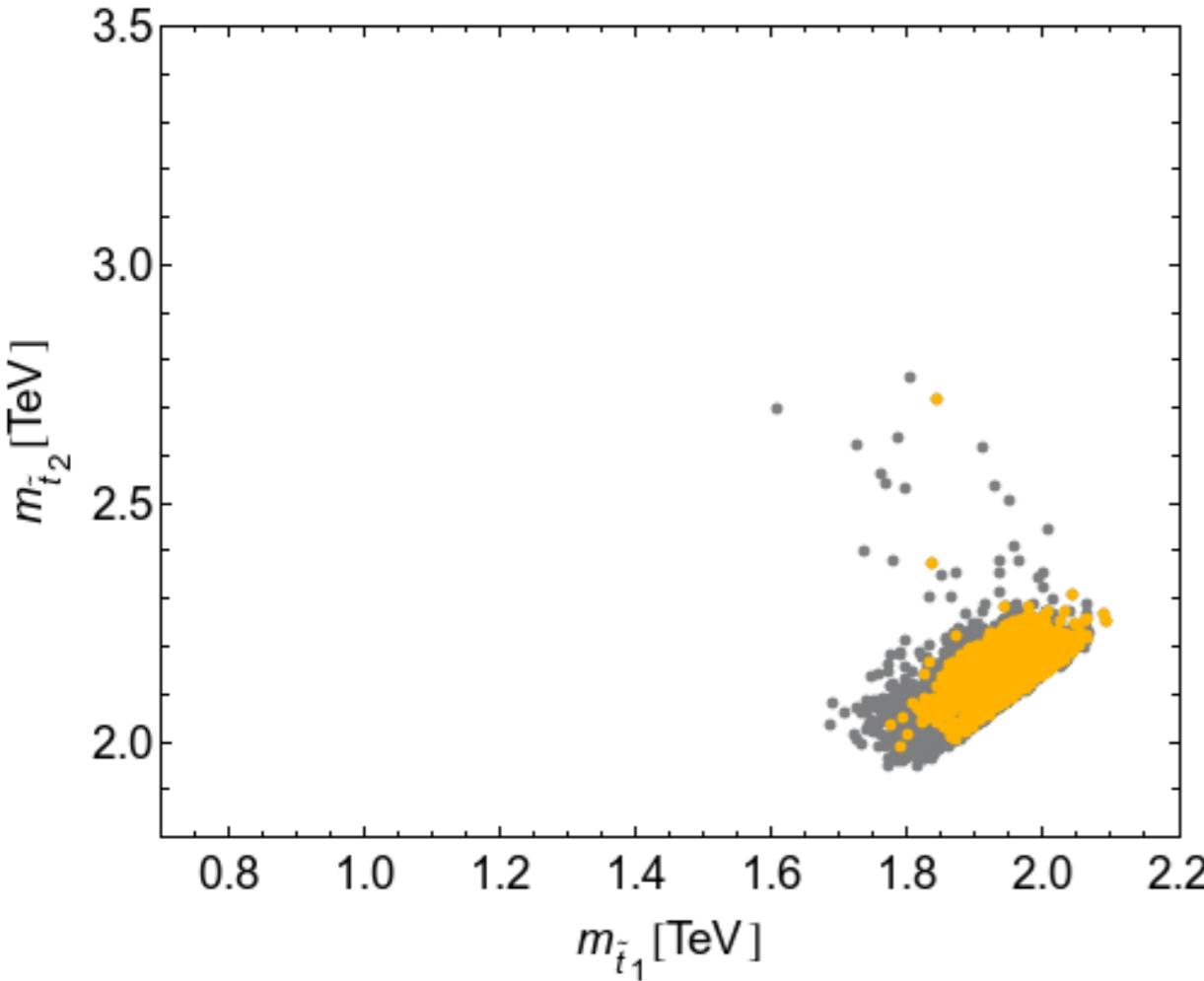




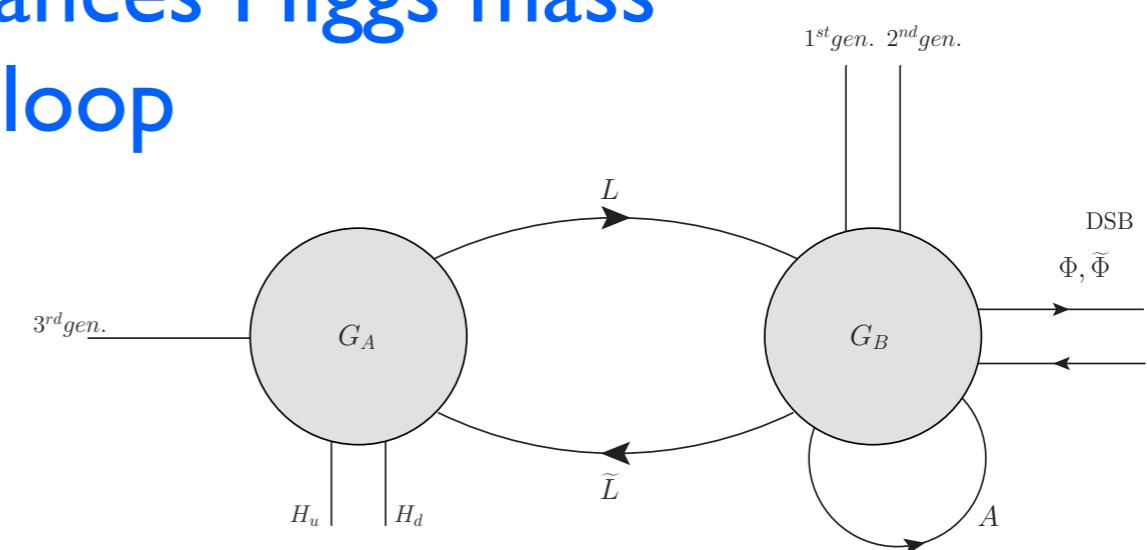
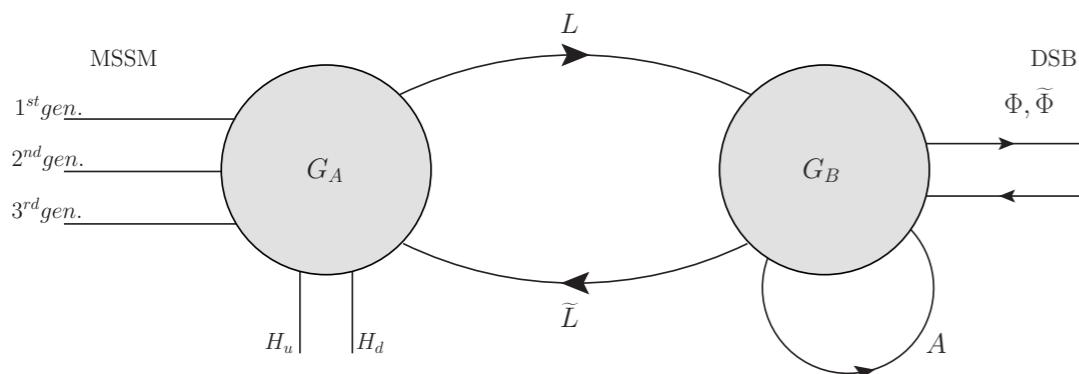
Testable at ILC!



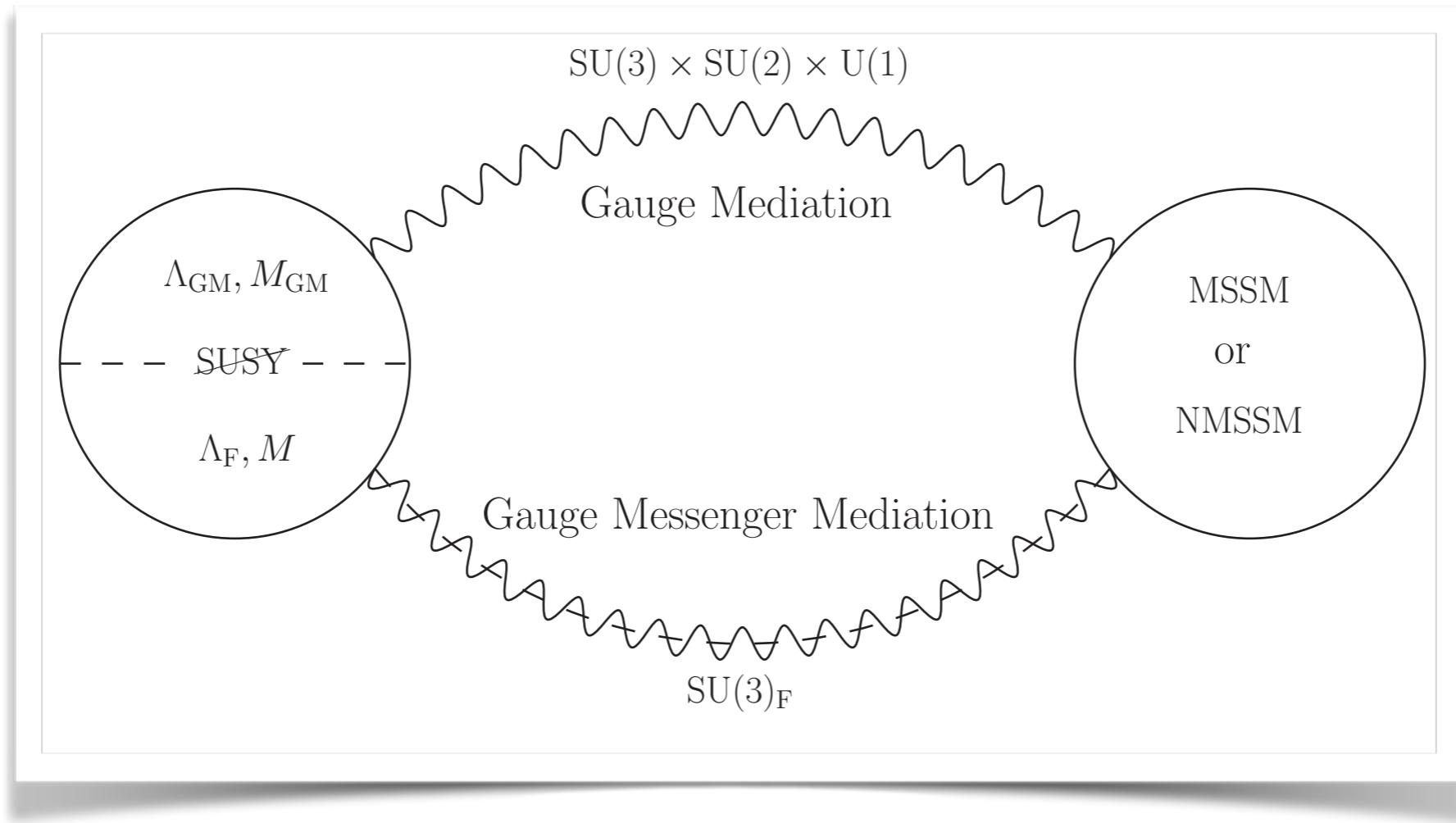
$$m_h^2 \simeq m_z^2 \cos^2(2\beta) + \frac{3}{(4\pi)^2} \frac{m_t^4}{v_{ew}^2} \left[\log \frac{m_{\tilde{t}}^2}{m_t^2} + \frac{X_t^2}{m_{\tilde{t}}^2} \left(1 - \frac{X_t^2}{12m_{\tilde{t}}^2}\right) \right] + \Delta m_{h,1}^2 = \frac{3m_Z^2}{16\pi^2 v_{ew}^2} \left(1 - \frac{8}{3} \sin \theta_W^2\right) \cos 2\beta m_t^2 \ln \left(\frac{m_{\tilde{q}_L}^2}{m_{\tilde{u}_R}^3} \right)$$



**Stop splitting enhances Higgs mass
@ 1-loop**



2. Flavour Gauge Messengers



- Extend gauge mediation to include a gauged flavour group
- Explain Yukawas and SUSY breaking
- Fields break $SU(3)_F$ and SUSY at the same time
- Fully dynamical origin in terms of Meta-stable SUSY breaking

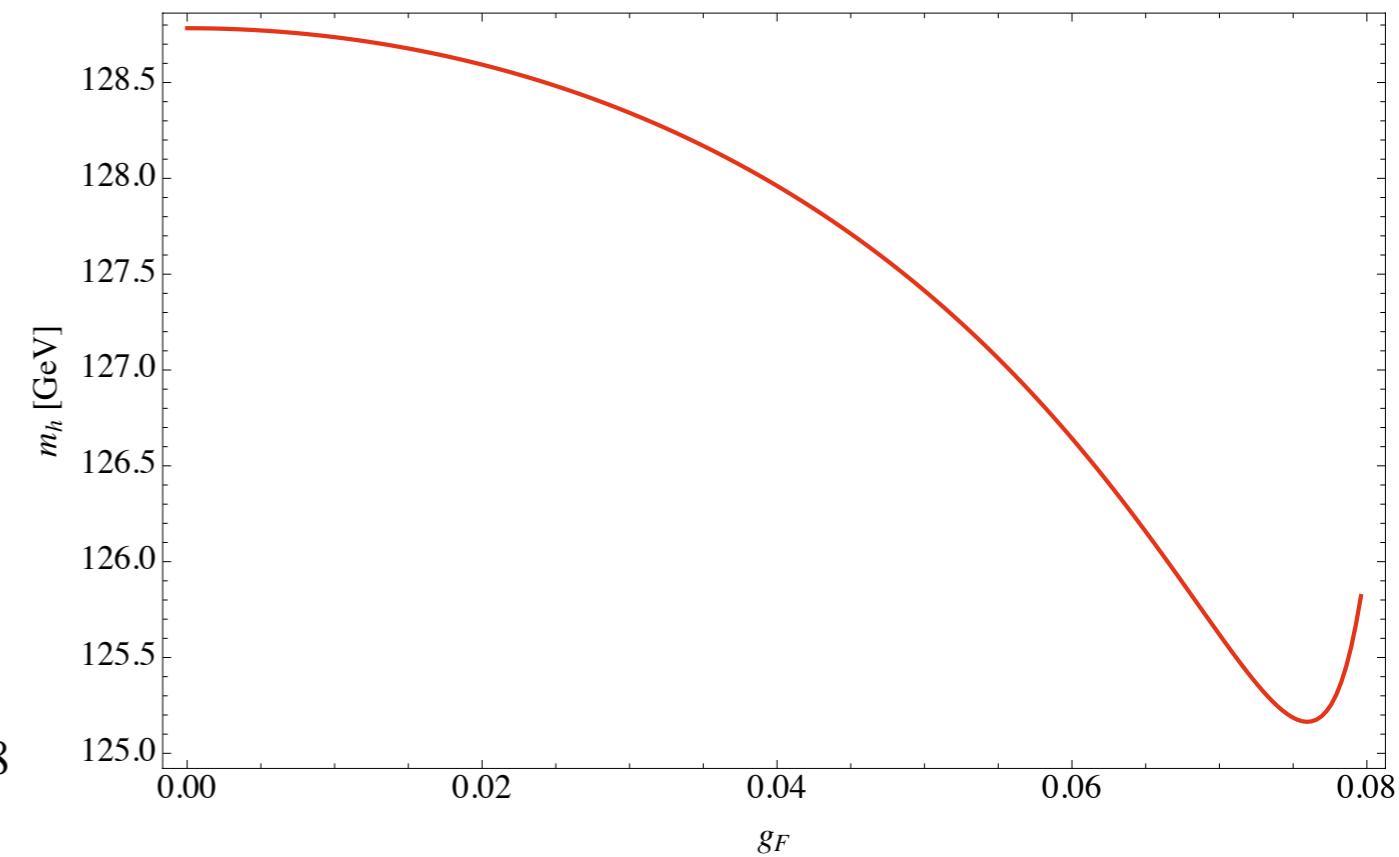
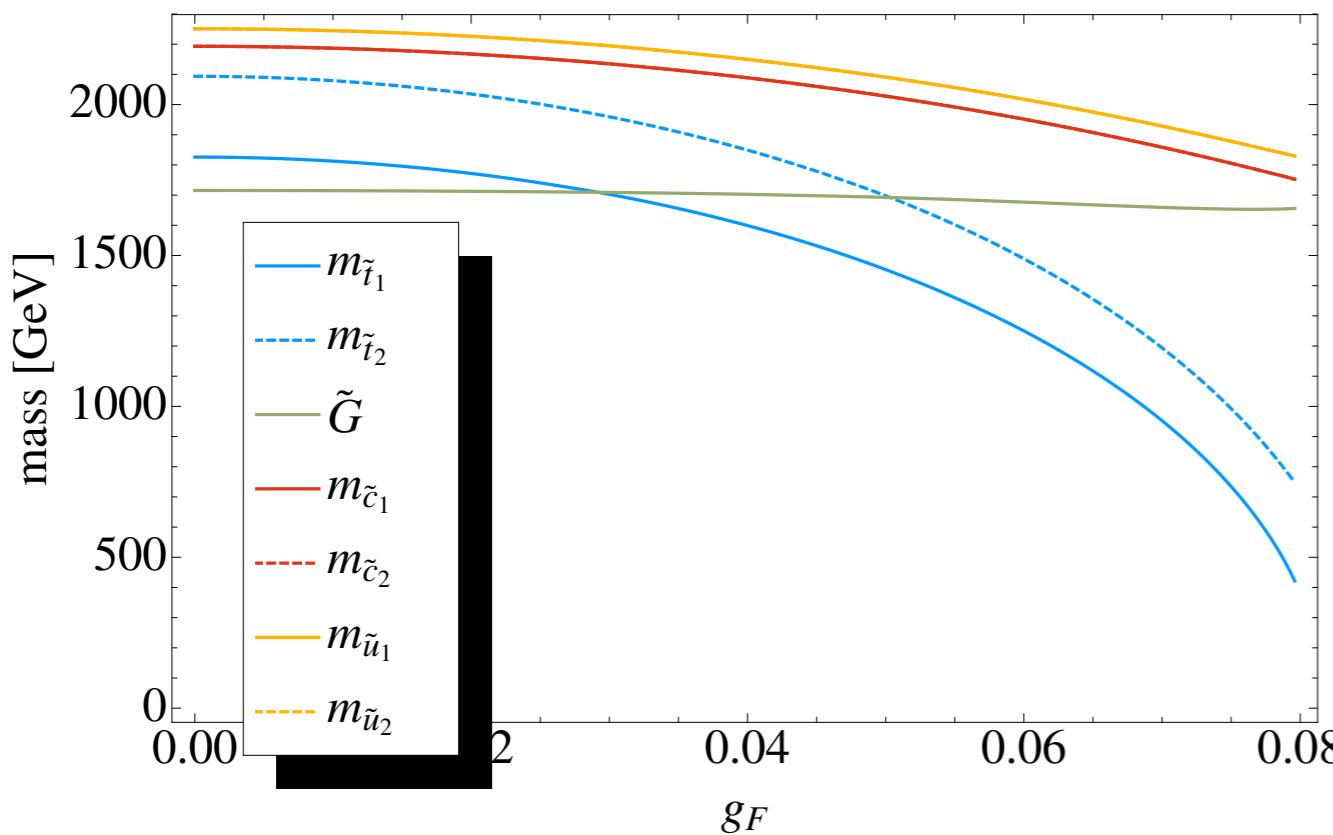
From GMSB

$$m_{Q,U,D,\text{GMSB}}^2 \sim + \sum_i \frac{g_{SM,i}^4}{(16\pi^2)^2} \left(\frac{F}{M}\right)^2 \begin{pmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

From flavour gauge mess.

$$\delta m_{Q,U,D}^2 = - \frac{g_F^2}{16\pi^2} \left(\frac{F}{M}\right)^2 \begin{pmatrix} \frac{7}{6} & 0 & 0 \\ 0 & \frac{7}{6} & 0 \\ 0 & 0 & \frac{8}{3} \end{pmatrix}$$

Stick the model into an NMSSM spectrum generator
(SPheno)



Squarks and Gluino

Higgs

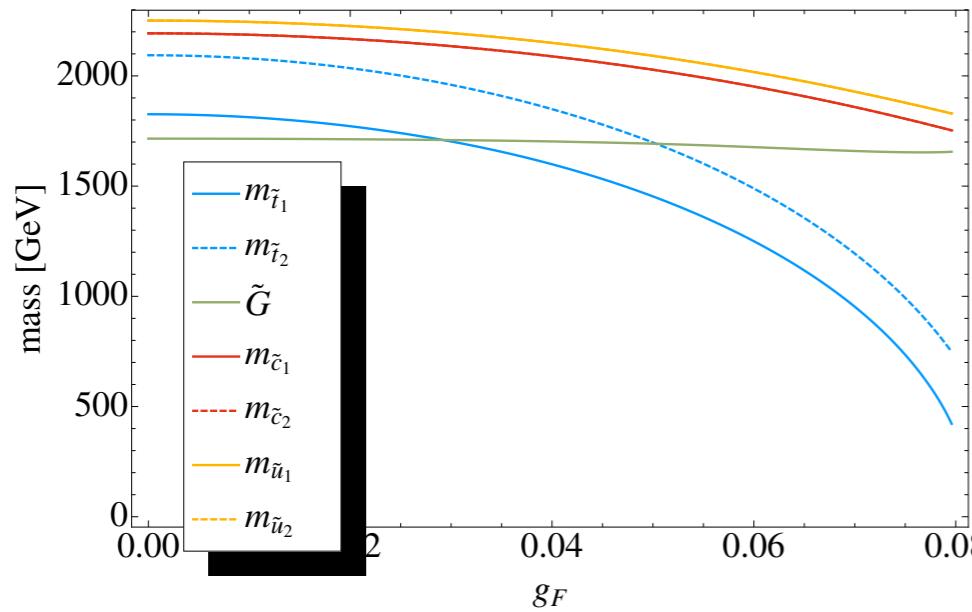
Figure 2. A plot [Left] of the squark and gluino masses for model 1 with the NMSSM. [Right] a plot of Higgs mass versus g_F for the same range. $\lambda = 0.8$, $\kappa = 0.8$, $v_s = 1000$, $m_{H_d}^2 = m_{H_u}^2 = 10^5$, $\Lambda = \Lambda_F = 2.3 \times 10^5$, $M = 10^7$, $\tan \beta = 1.5$.

Flavour changing neutral currents

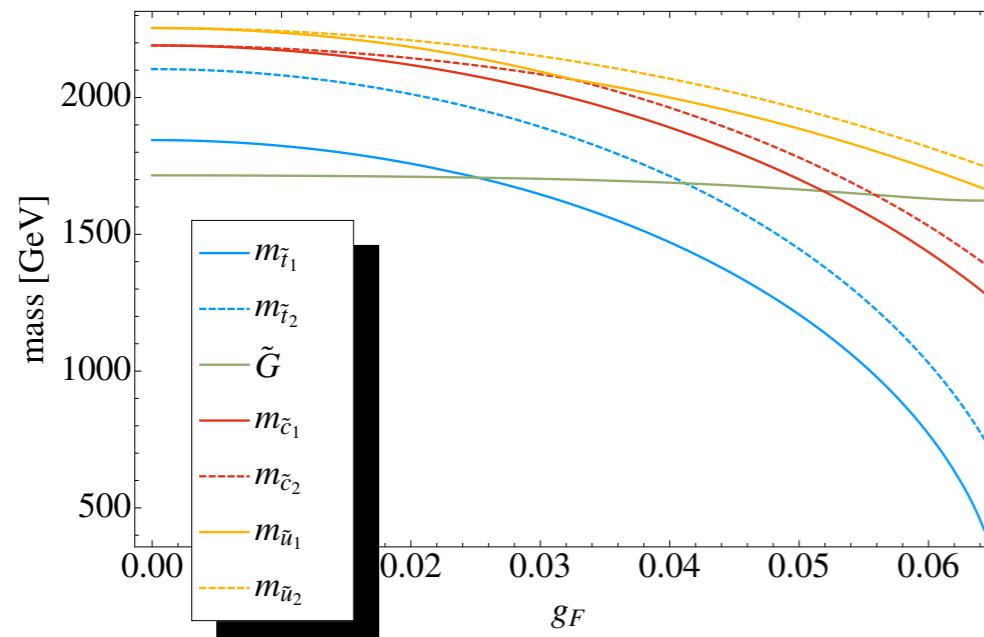
(S.Abel & MM) I404.I318

Model I: degenerate 1st & 2nd

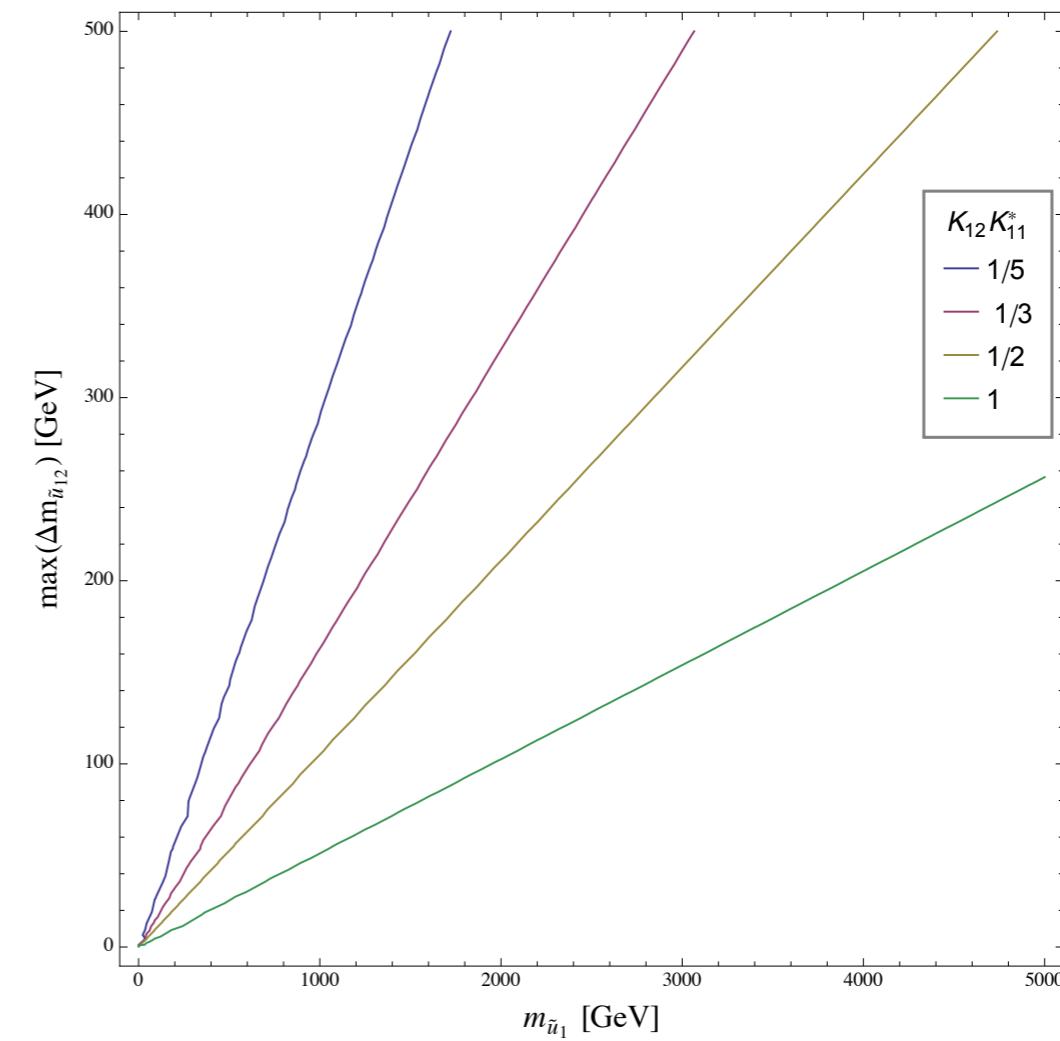
$$\delta_{u,12} < 0.1$$



Model 2: split 1st & 2nd



$$\delta_{ij} = \frac{m_{q_2}^2 - m_{q_1}^2}{\frac{1}{2}(m_{q_2}^2 + m_{q_1}^2)} K_{ij} K_{ii}^*$$



Tachyons are natural?!

For a natural cancellation these should be of the same order

$$m_z^2 = -2(m_{H_u}^2 + |\mu|^2) + \dots$$

Massless stops at Mplanck, turn tachyonic at messenger scale, are turned positive by gluino

stops run positive

$$\delta m_{\tilde{t}}^2 = -\frac{8\alpha_s M_3^2}{3\pi} \text{Log}\left(\frac{\Lambda}{M_3}\right)$$

$$\delta m_{H_u}^2 \sim -\frac{3y_t^2 m_{\tilde{t}}^2}{4\pi^2} \text{Log}\left(\frac{\Lambda}{m_{\tilde{t}}}\right) \quad (+) + (-) \sim 0$$

Reduces fine tuning on the Higgs.

It turns out that this model can embed into magnetic SQCD too!

Field	G_{SM}	$SU(3)_L \times SU(3)_R$
\hat{Q}^f	$(2, \frac{1}{6}, \mathbf{3})$	$(\bar{\mathbf{3}}, 1)$
\hat{L}^f	$(2, -\frac{1}{2}, \mathbf{1})$	$(\bar{\mathbf{3}}, 1)$
\hat{H}_d	$(2, -\frac{1}{2}, \mathbf{1})$	$(1, 1)$
\hat{H}_u	$(2, \frac{1}{2}, \mathbf{1})$	$(1, 1)$
\hat{D}^f	$(1, \frac{1}{3}, \bar{\mathbf{3}})$	$(1, \mathbf{3})$
\hat{U}^f	$(1, -\frac{2}{3}, \bar{\mathbf{3}})$	$(1, \mathbf{3})$
\hat{E}^f	$(1, 1, \mathbf{1})$	$(1, \mathbf{3})$
$\hat{\nu}^f$	$(0, 1, \mathbf{1})$	$(1, \mathbf{3})$

“Dynamical metastable flavour gauge mediation”

Field	$SU(\tilde{N})_{\text{mag}}$	$SU(3)_L \times SU(3)_R$
Φ	1	$(\mathbf{3}, \bar{\mathbf{3}})$
φ	\square	$(\bar{\mathbf{3}}, 1)$
$\tilde{\varphi}$	$\bar{\square}$	$(1, \mathbf{3})$

$$W_{\text{mag}} = h \text{Tr} \varphi \Phi \tilde{\varphi} - \mu^2 \text{Tr} \Phi.$$

$$\mu_{ij} = \begin{pmatrix} \mu & 0 & 0 \\ 0 & \mu & 0 \\ 0 & 0 & \mu_3 \end{pmatrix} \quad \text{and} \quad \varphi^T = \tilde{\varphi} = \begin{pmatrix} \mu \\ \mu \\ 0 \end{pmatrix}$$

$$F_\Phi = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & h\mu_3^2 \end{pmatrix} \quad \text{such that} \quad V_{\min} = |h^2 \mu_3^4|,$$

$$\delta m_{Q,U,D}^2 = -\frac{g_F^2}{16\pi^2} \frac{|h^2 \mu_3^4|}{\mu^2} \begin{pmatrix} \frac{8}{9} & 0 & 0 \\ 0 & \frac{8}{9} & 0 \\ 0 & 0 & \frac{20}{9} \end{pmatrix}$$

Perhaps we can explain Yukawas too!

$$W = \frac{\lambda_u}{\Lambda} H_u Q \Phi U + \frac{\lambda_d}{\Lambda} H_d Q \Phi D$$

(S.Abel & MM) 1404.1318

Available for hire :)

Thanks for listening!

Conclusions

- Most models are in bad shape
- Perhaps it is a time to panic?
- Natural SUSY is well motivated from bottom up
- Some of these have top-down motivation too
- It does mean sacrificing minimality!