School of Physics & Astronomy





# Thermal Physics

# PHYS09061 (SCQF Level 9)

## Monday 7<sup>th</sup> May, 2018 14:30 - 17:30 (May Diet)

### Please read full instructions before commencing writing.

#### Examination Paper Information

Answer **ALL** questions from Section A and **THREE** questions from Section B & C, answering at least one question from each section.

#### **Special Instructions**

- Only authorised Electronic Calculators may be used during this examination.
- A sheet of physical constants is supplied for use in this examination.
- Attach supplied anonymous bar codes to *each* script book.

#### Special Items

- School supplied Constant Sheets
- School supplied barcodes

Chairman of Examiners: Prof A Trew External Examiner: Prof S J Clark

Anonymity of the candidate will be maintained during the marking of this examination.

### Section A: Answer ALL the questions from this Section

A.1 The low-temperature specific heat of diamond varies with temperature, T (in Kelvin) as

$$c_p = 124 \left(\frac{T}{1860}\right)^3 \,\mathrm{J/g/K} \;.$$

- (a) Calculate the entropy change when 1g of diamond is heated at constant pressure from 4K to 300K.
- (b) Does this equation for specific heat satisfy the third law of thermodynamics? Justify your answer.
- A.2 By considering entropy as a function of T and V, show that for an ideal gas

$$\left(\frac{\partial S}{\partial T}\right)_P = \frac{nR}{T} + \left(\frac{\partial S}{\partial T}\right)_V \,. \tag{5}$$

A.3 The specific Helmholtz free energy of a particular gas is

$$f = cT^2 - \frac{a}{v} - RT\ln(v - b)$$

where a, b and c are constants. Calculate the pressure of the gas in terms of T and v. [5]

- **A.4** In statistical mechanics, the entropy S(E) of a macrostate with energy E is defined as  $S(E) = k_B \ln \Omega(E)$ , where  $\Omega(E)$  is the weight function.
  - (a) Explain what is meant by "the weight function".
  - (b) By appealing to the principle of equal *a priori* probabilities, explain why entropy tends to increase over time.
- A.5 An electron is confined to a one-dimensional channel of length  $6 \times 10^{-12}$  m which is in contact with a heat reservoir at temperature T. The electron's energy is determined by its energy level n = 1, 2, 3, ... through the formula  $\epsilon_n = 0.01n^2$  eV. Determine the temperature T at which the electron is 20 times more likely to be in the ground state than the n = 3 state.

[5]

[2]

[3]

[3]

[2]

- A.6 (a) State which of the following quantities is fixed and which can fluctuate in the grandcanonical ensemble:
  - (i) Energy, E
  - (ii) Temperature, T
  - (iii) Particle number,  ${\cal N}$
  - (iv) Chemical potential,  $\mu$ .
  - (b) Write down the form of the grand-canonical distribution function in terms of the quantities (i)-(iv) above. [3]

[2]

# Section B: Answer AT LEAST ONE of the questions in this Section

B.1 This question concerns a three-stage engine.

A heat engine cycle using one mole of ideal gas consists of the following three stages:

1. an isobaric compression from volume  $V_a$  to  $V_b$  at pressure  $P_a$ 

2. an increase in pressure from  $P_a$  to  $P_b$  at a constant volume  $V_b$ 

3. an adiabatic expansion from  $(P_b, V_b)$  to the original state at  $(P_a, V_a)$ 

You may assume throughout that  $\mathrm{PV}^{\gamma}$  is constant in the adiabatic process.

(a) (b)	Write down the relationship between heat capacity $(C_v)$ and internal energy for an ideal gas.	[1]
	Sketch this cycle on a PV plot.	[4]
(c)	At what point is the temperature highest, and what is its value?	[2]
(d)	Calculate the heat and work at each stage of the cycle.	[6]
(e)	Define the efficiency of the engine in terms of the heat input and output.	[3]
(f)	Show that the efficiency of the cycle described above is	

$$\eta = 1 - \gamma \frac{r-1}{(r^{\gamma} - 1)}$$

where  $r = V_a/V_b$ .

[4]

- **B.2** This question examines  $CO_2$  going from solution to atmosphere.
  - (a) Write down the general expression relating chemical potential to Gibbs free energy. [2]
  - (b) Calculate the change in chemical potential of an ideal gas when pressurised isothermally from 1atm to 5atm.

A 1.05 litre bottle contains 1.00 litres of fizzy lemonade pressurized to  $P_g = 5$  atm with CO<sub>2</sub> gas. The CO<sub>2</sub> concentration in the lemonade is initially x = 0.2 moles/litre. Treat the gaseous CO<sub>2</sub> as an ideal gas at 300K, and take the chemical potential of the dissolved CO<sub>2</sub> proportional to the concentration,  $\mu = Ax$ .

(c) What three thermodynamic quantities have equal value in the lemonade and in the gas above it?

The bottle cap is briefly loosened, so that the gas comes fully into equilibrium with the atmosphere, but all the dissolved  $CO_2$  remains in solution. The bottle is then sealed.

(d) What are the chemical potentials of the CO<sub>2</sub> in the lemonade, and in the gas above it, at the moment when the bottle is sealed? [4]

Now the bottle is shaken, which allows the gas and solution to come into equilibrium.

- (e) Describe what happens and estimate how many moles of CO<sub>2</sub> come out of solution. [5]
   (*Hint: Because the volume of liquid is larger than the gas, assume that it acts as a reservoir of constant chemical potential.*)
- (f) Approximately how many times can the process of loosening and tightening the cap be repeated before the drink goes flat? [3]

[3]

[3]

- **B.3** This question concerns compressing an elastic solid.
  - (a) Show that for a piston of area A, moving some distance  $d\mathbf{x}$ , subject to force  $\mathbf{F}$ , the two expressions for work,  $\mathbf{F}.d\mathbf{x}$  and PdV are equivalent. [2]
  - (b) Write down the central equation of thermodynamics.
  - (c) Use this equation to show that

$$\left(\frac{\partial U}{\partial V}\right)_T = T \left(\frac{\partial P}{\partial T}\right)_V - P .$$
[3]

[2]

[2]

[5]

A block of solid material has an equation of state that involves two constants a and  $V_0$ :

$$P = aT\left(\frac{V_0^2}{V^2} - \frac{V}{V_0}\right)$$

- (d) Calculate the volume of the block at zero pressure.
- (e) Define and evaluate the isothermal bulk modulus of the material at  $V = V_0$ .
- (f) In a high pressure experiment, the block is compressed isothermally from  $V_0$  to  $V_0/2$ . Calculate work done, change in internal energy, and heat supplied in the experiment. [6]

# Section C: Answer AT LEAST ONE of the questions in this Section

C.1 This question concerns an assembly of N weakly-interacting dipoles, each of which can be oriented in one of four directions,  $\underline{m}_1, \underline{m}_2, \underline{m}_3$  and  $\underline{m}_4$ , as shown in the figure. The magnitude of the dipole moment, m, is the same in all four configurations.

 $\begin{array}{c|c} \underline{H} \\ & \underline{m}_1 \\ & & \underline{m}_2 \\ & & & \underline{m}_3 \end{array} \qquad \underbrace{\underline{m}_4} \\ & & & & \underbrace{\underline{m}_4} \\ & & & & & & \underbrace{\underline{m}_4} \\ & & & & & & & \underbrace{\underline{m}_4} \\ & & & & & & & \underbrace{\underline{m}_4} \\ & & & & & & & & \underbrace{\underline{m}_4} \\ & & & & & & & & \underbrace{\underline{m}_4} \\ & & & & & & & & & & \\ \end{array}$ 

The energy of a dipole is given by  $\epsilon(\underline{m}) = -\underline{m} \cdot \underline{H}$  where  $\underline{H}$  is the applied field. The assembly is in equilibrium with a heat reservoir at temperature T.

- (a) Determine the distinct energies that a single dipole (N = 1) can have, and the number of microstates corresponding to each one. [3]
- (b) Show that the partition function for a single dipole, Z(1), is

$$Z(1) = 2[1 + \cosh\left(\beta mH\right)]$$

where m and H are the magnitudes of  $\underline{m}$  and  $\underline{H}$ , respectively.

(c) Show that the mean energy of a single dipole is

$$\bar{\epsilon} = -mH \frac{\sinh(\beta mH)}{1 + \cosh(\beta mH)} \,. \tag{3}$$

[3]

[4]

The magnetisation of a system of N dipoles is defined as  $M = m(N_{\downarrow} - N_{\uparrow})$ , where  $N_{\downarrow}$ and  $N_{\uparrow}$  are the total number of dipoles that lie parallel and antiparallel to the external field, respectively. The magnetic susceptibility is defined as  $\chi(H,T) = \frac{\partial \overline{M}}{\partial H}$ .

- (d) Show that the mean magnetisation and energy are related by  $\overline{M} = -N\overline{\epsilon}/H$ . [4]
- (e) Consider the regime where  $mH \ll k_B T$  and show that the quantity

$$\chi(0,T) \equiv \lim_{H \to 0} \chi(H,T)$$

satisfies the Curie law  $\chi(0,T) \propto 1/(k_B T)$ .

(f) Explain physically why we expect the susceptibility to diverge in the  $T \to 0$  limit. [3]

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C.2 This question concerns the three-dimensional ideal gas within the semi-classical approximation. The density of states of energy  $\epsilon$  for a single particle of mass M confined to a volume V is

$$g(\epsilon) = \pi V \left(\frac{M}{\hbar^2 \pi^2}\right)^{\frac{3}{2}} \sqrt{\frac{\epsilon}{2}} .$$

The gas is at equilibrium with a reservoir at temperature T.

(a) Explain why, within the semi-classical approximation, the partition function for N weakly-interacting particles is written as

$$Z(N) \approx \frac{[Z(1)]^N}{N!}$$

where Z(1) is the single-particle partition function.

(b) By approximating the partition function  $Z(1) = \sum_{\epsilon} \Omega(\epsilon) e^{-\epsilon/k_B T}$  as an integral involving the density of states, show that

$$Z(1) \approx C\pi V \left(\frac{Mk_BT}{\hbar^2 \pi^2}\right)^{\frac{3}{2}}$$

where C is a dimensionless constant that you do not need to calculate.

- (c) Determine the equilibrium Helmholtz free energy of the N-particle gas under the assumption that  $N \gg 1$ .
- (d) Given that the mean energy of the ideal gas is  $\bar{E} = \frac{3}{2}Nk_BT$ , show that the entropy is

$$S = Nk \left[ \frac{3}{2} \ln \left( \frac{Mk_B T}{\hbar^2 \pi^2} \right) - \ln \left( \frac{N}{C\pi V} \right) + \frac{5}{2} \right] .$$
 [3]

(e) Show that the entropy becomes unphysical at a temperature proportional to  $\rho^{2/3}$ , where  $\rho = N/V$ . What is the origin of this problem? [5]

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[4]

[5]

[3]

C.3 The energy levels of a spin- $\frac{1}{2}$  fermion in one dimension are

$$\epsilon_n = \frac{\hbar^2 \pi^2 n^2}{2ML^2}$$

where n is a positive integer, M is the mass of the fermion and L is the length of a line that the fermion is confined to.

(a) Show that the density of states for a spin- $\frac{1}{2}$  fermion in one dimension is

$$g(\epsilon) = \sqrt{\frac{2ML^2}{\hbar^2 \pi^2 \epsilon}} .$$
[5]

(b) Sketch the  $T \to 0$  limiting form of the Fermi-Dirac distribution,

$$f_+(\epsilon) = \frac{1}{\mathrm{e}^{\beta(\epsilon-\mu)}+1} \; ,$$

noting the position of the Fermi energy  $\epsilon_{\rm f}$ .

(c) By considering the mean number of particles  $\bar{N}$  at T = 0, show that the Fermi energy for a one-dimensional electron gas is

$$\epsilon_{\rm f} = \frac{\hbar^2 \pi^2 \bar{N}^2}{8ML^2} \,. \tag{4}$$

[3]

- (d) Obtain an expression for the mean energy of this gas at zero temperature that contains only  $\epsilon_{\rm f}$ ,  $\bar{N}$  and a numerical factor. [5]
- (e) Explain why the energy of the Fermi gas at low temperatures differs from that of the classical ideal gas in one dimension, which satisfies  $\bar{E} = \frac{1}{2}\bar{N}k_BT$ . [3]