

# Methods of Mathematical Physics

## Exam style questions

1. Cauchy's integral formula states that if  $f(z)$  is analytic within and on a closed contour  $C$

$$f(z_0) = \frac{1}{2\pi i} \oint_C \frac{f(z)}{z - z_0} dz$$

where  $z_0$  is within the contour  $C$ .

Consider a function  $f(z)$  analytic in the upper half plane and on the real axis, and such that  $|f(z)| \rightarrow 0$  as  $|z| \rightarrow \infty$  for  $0 \leq \arg z \leq \pi$ . Show by using Cauchy's integral formula and an appropriate closed contour  $C$  that

$$f(z_0) = \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{f(x)}{x - z_0} dx$$

where the integral is along the real axis . [6]

Now set  $z_0 = x \pm i\epsilon$  and let  $z_0$  approach the real axis. Show that

$$\int_{-\infty}^{\infty} \frac{f(x)}{x - x_0 \mp i\epsilon} dx = P \int_{-\infty}^{\infty} \frac{f(x)}{x - x_0} dx \pm i\pi f(x_0)$$
[6]

where P indicates the Cauchy principal value.

Hence deduce

$$\frac{1}{x - x_0 \mp i\epsilon} = P \frac{1}{x - x_0} \pm i\pi\delta(x - x_0)$$

and

$$f(x_0) = \frac{1}{\pi i} P \int_{-\infty}^{\infty} \frac{f(x)}{x - x_0} dx$$
[4]

Now consider a function  $f$  such that

$$f(-x) = f^*(x) .$$

Let  $f(x_0) = u(x_0) + iv(x_0)$ , where  $u(x_0)$  and  $v(x_0)$  are real functions, and show

$$\begin{aligned} u(x_0) &= \frac{2}{\pi} P \int_0^{\infty} \frac{xv(x)}{x^2 - x_0^2} dx \\ v(x_0) &= -\frac{2}{\pi} P \int_0^{\infty} \frac{x_0 u(x)}{x^2 - x_0^2} dx \end{aligned}$$
[9]

2.

(a) The Gamma function may be defined as

$$\Gamma(x) = \int_0^{\infty} u^{x-1} e^{-u} du \quad \text{where } \operatorname{Re}[x] > 0 .$$

Show that

$$\Gamma(1) = 1 \quad ; \quad \Gamma(1/2) = \sqrt{\pi} \quad ; \quad \Gamma(-1/2) = -2\sqrt{\pi} \quad [5]$$

Show that the Laplace transform of  $t^\lambda$  is given by

$$L[t^\lambda] = \frac{\Gamma(\lambda + 1)}{s^{\lambda+1}} \quad [3]$$

(b) The temperature distribution in a semi-infinite rod ( $x > 0$ ) obeys

$$\frac{\partial^2 T(x, t)}{\partial x^2} = \frac{1}{\kappa} \frac{\partial T}{\partial t}$$

with boundary conditions

$$\begin{aligned} T(x, 0) &= 0 \quad x > 0 \\ T(0, t) &= T_0 \end{aligned}$$

Show that the Laplace transform with respect to  $t$

$$F(x, s) = \int_0^{\infty} T(x, t) e^{-st} dt$$

obeys

$$\frac{\partial^2 F(x, s)}{\partial x^2} = s \frac{F(x, s)}{\kappa} . \quad [3]$$

Hence deduce

$$F(x, s) = \frac{T_0}{s} \exp\left(-\left(s/\kappa\right)^{1/2} x\right) . \quad [4]$$

(c) Identify the singularity of  $F(x, s)$  which controls the large  $t$  behaviour of  $T(x, t)$

Using the results of part a) compute the first three non-zero terms in the large  $t$  asymptotic expansion of  $F(x, t)$ . [10]

3. The generating function for the Bessel functions  $J_n(x)$  is given by

$$G(x, t) = \sum_{n=-\infty}^{\infty} t^n J_n(x) = \exp \left[ \frac{x}{2} \left( t - \frac{1}{t} \right) \right] .$$

From the symmetries  $G(x, t) = G(x, -1/t)$ ,  $G(x, t) = G(-x, 1/t)$  deduce that

$$J_{-|n|}(x) = (-1)^n J_{|n|}(x) = J_{|n|}(-x) . \quad [4]$$

Use the product of generating functions

$$G(x + y, t) = G(x, t)G(y, t)$$

to derive

$$J_n(x + y) = \sum_{k=-\infty}^{\infty} J_k(x) J_{n-k}(y) . \quad [4]$$

By expanding the generating function show that

$$J_n(x) = \sum_{s=0}^{\infty} \frac{(-1)^s}{(n+s)! s!} \left( \frac{x}{2} \right)^{n+2s} \quad \text{for } n \geq 0 . \quad [6]$$

Show that

$$\begin{aligned} R_n(x) J_n(x) &= J_{n+1}(x) \quad \text{where} \quad R_n(x) = \left( \frac{n}{x} - \frac{d}{dx} \right) \\ L_n(x) J_n(x) &= J_{n-1}(x) \quad \text{where} \quad L_n(x) = \left( \frac{n}{x} + \frac{d}{dx} \right) \end{aligned} \quad [6]$$

Hence deduce the second order differential equation satisfied by  $J_n(x)$ . [5]

#### 4. June 2002 Q1

Consider the partial differential equation

$$\nabla^2 u(\mathbf{x}) - m^2 u(\mathbf{x}) = \rho(\mathbf{x}) \quad (1)$$

in three dimensions.

a. Is this equation parabolic, elliptic, or hyperbolic? [1]

b. If  $G(\mathbf{x}, \mathbf{x}')$  satisfies  $(\nabla^2 - m^2)G(\mathbf{x}, \mathbf{x}') = \delta(\mathbf{x} - \mathbf{x}')$  show that

$$G(\mathbf{x}, \mathbf{x}') = -(2\pi)^{-3} \int d^3\mathbf{k} \frac{e^{-i\mathbf{k}\cdot(\mathbf{x}-\mathbf{x}')}}{|\mathbf{k}|^2 + m^2}. \quad (4)$$

c. Show that

$$\int_0^\pi d\theta \sin \theta e^{-i\kappa r \cos \theta} = \frac{2 \sin \kappa r}{\kappa r}. \quad (5)$$

d. Show by contour integration that

$$\int_0^\infty d\kappa \frac{2\kappa \sin \kappa r}{(\kappa^2 + m^2)} = \pi e^{-mr} \quad (5)$$

e. Hence prove that

$$G(\mathbf{x}, \mathbf{x}') = -\frac{e^{-m|\mathbf{x}-\mathbf{x}'|}}{4\pi|\mathbf{x}-\mathbf{x}'|}. \quad (6)$$

f. Write down the solution to equation (1). [4]

#### 5. June 2002 Q2

a. Find the Fourier transform of  $f(x) = e^{-x^2/2V}$ . [5]

b. Consider the function

$$g(x) = \sum_{n=-\infty}^{\infty} f(x + 2\pi n) = \sum_{n=-\infty}^{\infty} e^{-(x+2\pi n)^2/2V}; \quad (2)$$

show it is a periodic function, and find its period. [5]

c. Deduce that  $g(x)$  has a Fourier series

$$g(x) = \sum_{k=-\infty}^{\infty} \tilde{g}_k e^{ikx}; \quad (3)$$

and show that  $\tilde{g}_k = \sqrt{2\pi V} e^{-\frac{1}{2}k^2 V}$ . [8]

d. Discuss how to use the expansions (2) and (3) to find approximate values for  $g(x)$  for an arbitrary value of  $x$  in the cases where

(a)  $V$  is large, [3]

(b)  $V$  is small. [4]

## 6. June 2002 Q3

The Bessel function  $J_0(x)$  may be defined by Schl\"afli's integral

$$J_0(x) = \frac{1}{2\pi i} \oint_{|t|=1} \frac{dt}{t} \exp \left[ \frac{x}{2} \left( t - \frac{1}{t} \right) \right].$$

For real  $x \gg 1$  its asymptotic expansion may be found using the method of steepest descents.

- a. Show that in the complex  $t$  plane the dominant contributions to the integral come from regions around  $t_{\pm} = \pm i + 1/x + O(1/x^2)$ . [4]

- b. Show that the values of the integrand at the saddle points are

$$\exp \left\{ \pm i \left[ x - \frac{\pi}{2} + O\left(\frac{1}{x}\right) \right] \right\}. \quad [5]$$

- c. Show that the steepest descents contour passes through the saddle points  $t_{\pm}$  at angles  $\alpha_+ = 3\pi/4$  and  $\alpha_- = \pi/4$  to the positive real  $t$  axis. [8]

- d. Show that the leading asymptotic behaviour is

$$J_0(x) \sim \sqrt{\frac{2}{\pi x}} \cos \left( x - \frac{\pi}{4} \right). \quad [8]$$

[N.B. This question employs a slightly different approach to that of the lectures.]

## 7. Consider the integral

$$f_R(x) \equiv \frac{1}{2\pi i} \lim_{\varepsilon \rightarrow 0} \int_{-R}^R dz \frac{e^{ixz}}{z - i\varepsilon},$$

where  $x > 0$ .

- a. Show that

$$f_R(x) = 1 - \frac{1}{2\pi} \int_0^\pi d\theta e^{ixR\cos\theta}. \quad [4]$$

- b. Show that

$$|f_R(x) - 1| \leq \frac{1}{2\pi} \int_0^\pi d\theta e^{-xR\sin\theta}. \quad [4]$$

- c. Prove that  $\sin\theta \geq 2\theta/\pi$  for  $0 \leq \theta \leq \pi/2$ . [8]

d. Show that

$$|f_R(x) - 1| \leq \frac{1}{2xR}. \quad [6]$$

e. Prove that

$$\lim_{R \rightarrow \infty} f_R(x) = \theta(x) \equiv \begin{cases} 1 & \text{if } x > 0, \\ 0 & \text{if } x < 0. \end{cases} \quad [3]$$

### 8. June 2001 Q3

Calculate the leading asymptotic behaviour of the Airy integral

$$\text{Ai}(-x) = \frac{1}{\pi} \int_0^\infty dw \cos\left(\frac{1}{3}w^3 - xw\right) \quad x > 0, \quad [25]$$

as  $x \rightarrow \infty$ .

9. Let  $y_1$  and  $y_2$  be two solutions of Bessel's equation  $x^2 y'' + xy' + (x^2 - \nu^2)y = 0$ .

a. If  $W(y_1(x), y_2(x)) = W(x)$  is their Wronskian show that  $x^2 W' + xW = 0$ , and hence that  $W(x) = c/x$  where  $c$  is a constant. [4]

b. Using Frobenius' method one can show that the Bessel function

$$J_\nu(x) \equiv \sum_{r=0}^{\infty} \frac{(-1)^r}{r! \Gamma(\nu + r + 1)} \left(\frac{x}{2}\right)^{\nu+2r}$$

is a solution of of Bessel's equation. Manipulate this series to show that

$$J_{\nu-1}(x) \pm J_{\nu+1}(x) = \sum_{r=0}^{\infty} \frac{(-1)^r (\nu + r \mp r)}{r! \Gamma(\nu + r + 1)} \left(\frac{x}{2}\right)^{\nu+2r-1}.$$

and hence obtain the identities

$$\begin{aligned} J_{\nu-1}(x) + J_{\nu+1}(x) &= \frac{2\nu}{x} J_\nu(x) \\ J_{\nu-1}(x) - J_{\nu+1}(x) &= 2J'_\nu(x). \end{aligned} \quad [6]$$

c. Prove that

$$W(J_\nu(x), J_{-\nu}(x)) = \frac{x}{2\nu} (J_{\nu-1}(x)J_{-\nu-1}(x) - J_{\nu+1}(x)J_{-\nu+1}(x)). \quad [6]$$

d. Show that

$$W(J_\nu(x), J_{-\nu}(x)) = -\frac{2 \sin \nu\pi}{\pi x}.$$

*Hint: Consider the behaviour of  $W(x)$  for small  $x$ , and recall that  $B(1 + \nu, -\nu) = -\pi / \sin \nu\pi$ .* [6]

e. Under what circumstances are  $J_\nu$  and  $J_{-\nu}$  linearly independent? [3]

**10. June 2003 Q1**

- (a) Find the position and nature of the singularities of the following differential equation.

$$x^2 y''(x) + xy'(x) - (x^2 + \nu^2)y(x) = 0. \quad (4) \quad [6]$$

- (b) Set  $y(x) = x^{-\frac{1}{2}}u(x)$  and show that  $u$  satisfies

$$u''(x) - \left[1 + \left(\frac{\nu^2}{x^2} - \frac{1}{4x^2}\right)\right]u(x) = 0. \quad [5]$$

- (c) Hence deduce that a solution which is bounded as  $x \rightarrow \infty$  has the asymptotic behaviour

$$y(x) \sim Ax^{-\frac{1}{2}}e^{-x} \quad \text{as } x \rightarrow \infty. \quad [4]$$

- (d) A solution of (4) is given by

$$y(x) = \frac{1}{2} \int_0^\infty dt t^{\nu-1} \exp\left[-\frac{x}{2}\left(t + \frac{1}{t}\right)\right],$$

[you are not required to show this].

Using Laplace's method determine the asymptotic behaviour of this solution as  $x \rightarrow \infty$ .

[10]

## 11. June 2003 Q2

A recurrence relation for Hermite polynomials is given by

$$H_{n+1}(x) = 2x H_n(x) - 2n H_{n-1}(x) \quad n \geq 0 \quad (1)$$

where  $H_0(x) = 1$  and  $H_{-1}(x) = 0$ .

Consider the generating function defined as

$$G(x, t) = \sum_{n=0}^{\infty} \frac{t^n}{n!} H_n(x). \quad (2)$$

(a) Show that

$$\begin{aligned} \frac{\partial G(x, t)}{\partial t} &= \sum_{n=0}^{\infty} \frac{t^n}{n!} H_{n+1}(x) \\ tG(x, t) &= \sum_{n=1}^{\infty} \frac{n t^n}{n!} H_{n-1}(x). \end{aligned} \quad [5]$$

(b) Hence deduce

$$G(x, t) = e^{-t^2+2tx}. \quad (3) \quad [4]$$

(c) From expression (3) for  $G(x, t)$  establish the relation

$$\frac{dH_n(x)}{dx} = 2n H_{n-1}(x). \quad (4) \quad [4]$$

(d) Using the definition (2) of  $G(x, t)$  and the expression (3) deduce the integral representation

$$H_n(x) = \frac{n!}{2\pi i} \oint dt \frac{e^{-t^2+2tx}}{t^{n+1}}, \quad (5)$$

where the closed contour encircles the origin.

Verify that this expression satisfies the relation (4). [5]

(e) Use the integral representation (5) to demonstrate the orthogonality of the Hermite polynomials

$$\int_{-\infty}^{\infty} dx e^{-x^2} H_m(x) H_n(x) = 2^n n! \sqrt{\pi} \delta_{m,n}. \quad [7]$$

Hint: Perform the  $x$  integral first then evaluate the subsequent two contour integrals.



## 12. June 2003 Q3

Consider the wave equation in three space dimensions and one time dimension for a system with periodic forcing

$$\nabla^2 u(\underline{x}, t) - \frac{1}{c^2} \frac{\partial^2 u(\underline{x}, t)}{\partial t^2} = f(\underline{x}) e^{-i\omega_0 t}. \quad (1)$$

(a) Show that  $u(\underline{x}, t)$  may be written as

$$u(\underline{x}, t) = u(\underline{x}) e^{-i\omega_0 t},$$

where  $u(\underline{x})$  satisfies

$$(\nabla^2 + k_0^2) u(\underline{x}) = f(\underline{x}) \quad \text{with} \quad k_0 = \frac{\omega_0}{c}. \quad [2]$$

(b) If  $G(\underline{x}, \underline{x}')$  satisfies

$$(\nabla^2 + k_0^2) G(\underline{x}, \underline{x}') = \delta(\underline{x} - \underline{x}')$$

use  $G(\underline{x}, \underline{x}')$  to construct a solution to (1) .

[4]

(c) Show by using Fourier transformation that

$$G(\underline{x}, \underline{x}') = -\frac{1}{(2\pi)^3} \int d^3k \frac{e^{i\mathbf{k} \cdot (\underline{x} - \underline{x}')}}{k^2 - k_0^2} \quad \text{where} \quad k = |\underline{k}|. \quad [5]$$

(d) Perform the angular integrals in  $\underline{k}$ -space to obtain

$$G(\underline{x}, \underline{x}') = \frac{1}{(2\pi)^2} \int_{-\infty}^{\infty} dk \frac{k}{ir} \frac{e^{-ikr}}{k^2 - k_0^2} \quad \text{where} \quad r = |\underline{x} - \underline{x}'|. \quad [5]$$

(e) Perform this integral by contour integration. You should choose a contour that corresponds to outgoing waves  $u \sim e^{i(k_0 r - \omega_0 t)}$ , i.e., you should obtain the retarded Green function.

[9]

## 13. June 2004 Q1

The Legendre Polynomials may be defined through the generating function

$$G(x, t) = \frac{1}{(1 - 2xt + t^2)^{1/2}}$$

where  $G(x, t)$  is defined as

$$G(x, t) = \sum_{n=0}^{\infty} t^n P_n(x).$$

- (a) By taking derivatives of  $G(x, t)$  with respect to  $t$  and  $x$  derive the following recursion relations

$$\begin{aligned}(n+1)P_{n+1}(x) + nP_{n-1}(x) &= (2n+1)xP_n(x) \\ P'_{n+1}(x) - 2xP'_n(x) + P'_{n-1}(x) &= P_n(x) .\end{aligned}\tag{8}$$

- (b) Show that

$$\int_{-1}^1 G^2(x, t) dx = \frac{1}{t} [\ln(1+t) - \ln(1-t)] .$$

By expanding the r.h.s. in powers of  $t$  and assuming

$$\int_{-1}^1 dx P_n(x) P_m(x) = 0 \quad \text{for } n \neq m$$

deduce

$$\int_{-1}^1 P_n(x) P_{n'}(x) dx = \frac{2\delta_{n,n'}}{2n+1} .\tag{7}$$

- (c) A function  $f(x)$  may be expressed as

$$f(x) = \sum_{k=0}^{\infty} a_k P_k(x) .$$

Use the result of part (b) to show that

$$a_n = \left(n + \frac{1}{2}\right) \int_{-1}^1 dx f(x) P_n(x)$$

and

$$\delta(x-x') = \frac{1}{2} \sum_{k=0}^{\infty} (2k+1) P_k(x) P_k(x') .\tag{5}$$

- (d) Show that  $G(x, t)$  implies the following integral representation of  $P_n(x)$

$$P_n(x) = \oint \frac{dt}{2\pi i} \frac{1}{t^{n+1} (1-2xt+t^2)^{1/2}} ,$$

where the closed contour encircles the origin. Specify the singularities of the integrand for the cases  $|x| < 1$  and  $|x| \geq 1$  and discuss the restrictions they impose on the radius of a circular contour.

[5]

**14. June 2004 Q2**

(a) The Laplace transform of a function  $f(t)$  is defined as

$$L[f] = \int_0^{\infty} f(t)e^{-st} dt, \quad \text{where } s > 0.$$

Show that

$$L[tf] = -\frac{\partial}{\partial s}L[f],$$

$$L\left[\frac{df}{dt}\right] = -f(0) + sL[f] \tag{3}$$

and

$$L[t^\lambda] = \frac{\Gamma(\lambda + 1)}{s^{\lambda+1}} \quad \text{where } \Gamma(x) = \int_0^{\infty} u^{x-1}e^{-u} du. \tag{2}$$

(b) Consider the differential equation

$$t^2 \frac{d^2 y(t)}{dt^2} + t \frac{dy(t)}{dt} - (1 + t^2)y(t) = 0.$$

Let  $g(s)$  denote the Laplace transform of a solution of this equation which is bounded as  $t \rightarrow \infty$ .

Show that  $g(s)$  obeys

$$(s^2 - 1)g''(s) + 3sg'(s) = 0. \tag{5}$$

Integrate this equation to obtain

$$g'(s) = \frac{A}{(s^2 - 1)^{3/2}}, \tag{3}$$

where  $A$  is a constant.

Assuming the inversion integral for Laplace transforms, deduce that

$$y(t) = -\frac{A}{t} \int_C \frac{ds}{2\pi i} \frac{e^{st}}{(s^2 - 1)^{3/2}},$$

where you should specify the contour  $C$ . \tag{2}

(c) Sketch the analytic structure of the integrand and identify the singularity that controls the large  $t$  behaviour of  $y(t)$ . \tag{3}

Compute the first two non-zero terms in the large  $t$  asymptotic expansion of  $y(t)$ . \tag{7}

**15. June 2004 Q3**

Consider the Green function  $G(\underline{x}, t)$  in three space dimensions and one time dimension defined by

$$\left[ \nabla^2 - \frac{\partial^2}{\partial t^2} - m^2 \right] G(\underline{x}, t) = \delta(\underline{x}) \delta(t) .$$

Show by Fourier transformation that

$$G(\underline{x}, t) = \frac{1}{(2\pi)^4} \int d^3k \int d\omega \frac{e^{i\mathbf{k}\cdot\underline{x} - i\omega t}}{(\omega^2 - k^2 - m^2)} \quad \text{where } k = |\underline{k}| . \quad [5]$$

Perform the angular integrations in  $\underline{k}$ -space to obtain

$$G(\underline{x}, t) = \frac{1}{(2\pi)^3 r} \int_{-\infty}^{\infty} dk \frac{k}{i} e^{ikr} \int_{-\infty}^{\infty} d\omega \frac{e^{-i\omega t}}{\omega^2 - \Omega^2} \quad \text{where } \Omega = \sqrt{k^2 + m^2}$$

and  $r = |\underline{x}|$ . [5]

Perform the  $\omega$  integral by contour integration, choosing the contour to give

$$\int_{-\infty}^{\infty} d\omega \frac{e^{-i\omega t}}{\omega^2 - \Omega^2} = -2\pi \frac{\sin(\Omega t)}{\Omega} \theta(t) ,$$

where  $\theta(t)$  is the usual Heaviside step function. [7]

Use the following results (which you are not required to show) for Bessel functions  $J_0, J_1$

$$\frac{1}{2\pi} \int_{-\infty}^{\infty} dk e^{ikr} \frac{\sin(\Omega t)}{\Omega} = \frac{1}{2} J_0 \left( m\sqrt{t^2 - r^2} \right) \theta(t^2 - r^2)$$

$$\begin{aligned} J_0'(y) &= -J_1(y) \\ J_0(0) &= 1 \end{aligned}$$

to deduce that

$$G(\underline{x}, t) = \left[ \frac{m}{4\pi\sqrt{t^2 - r^2}} J_1 \left( m\sqrt{t^2 - r^2} \right) \theta(t - r) - \frac{1}{4\pi r} \delta(r - t) \right] \theta(t) . \quad [5]$$

Interpret the two terms in  $G(\underline{x}, t)$  identifying the case  $m = 0$  . [3]