

Quark Model



Outline

Hadrons

Isospin, Strangeness

Quark Model

3 Flavours u, d, s

Mesons

Pseudoscalar
and vector mesons

Baryons

Decuplet, octet

Hadron Masses

Spin-spin coupling

Heavy Quarks

Charm, bottom,
Heavy quark

Mesons

Top quark

Hadrons known in 1960

Mesons	$\langle \text{Mass} \rangle$	J^{PC}	I	S
π^-, π^0, π^+	138.0	0^{-+}	1	0
K^0, K^+	495.7	0^{-}	1/2	+1
K^-, \bar{K}^0				-1
η	547.3	0^{-+}	0	0
ρ^-, ρ^0, ρ^+	770.0	1^{--}	1	0
ω	781.9	1^{--}	0	0
K^{*0}, K^{*+}	893.7	1^{-}	1/2	+1
K^{*-}, \bar{K}^{*0}				-1
η'	957.8	0^{-+}	0	0
ϕ	1019.5	1^{--}	0	0

Baryons	$\langle \text{Mass} \rangle$	J^P	I	S
p, n	938.9	$1/2^{+}$	1/2	0
Λ	1116	$1/2^{+}$	0	-1
$\Sigma^-, \Sigma^0, \Sigma^+$	1193	$1/2^{+}$	1	-1
$\Delta^-, \Delta^0, \Delta^+, \Delta^{++}$	1232	$3/2^{+}$	3/2	0
Ξ^-, Ξ^0	1318	$1/2^{+}$	1/2	-2
$\Sigma^{*-}, \Sigma^{*0}, \Sigma^{*+}$	1385	$3/2^{+}$	1	-1
Ξ^{*-}, Ξ^{*0}	1533	$3/2^{+}$	1/2	-2

Motivation for Quark Model

Particle "Zoo" proliferates

" ... the finder of a new particle used to be rewarded by a Nobel prize, but such a discovery ought to be punished by a \$10000 fine"

Lamb, 1955

Isospin



Nucleons

Proton and neutron have almost equal mass
Strong nuclear force is charge independent

$$V_{pp} \approx V_{pn} \approx V_{nn}$$

Isospin

p and n form part of single entity with isospin $\frac{1}{2}$ analogous to \uparrow and \downarrow of spin $\frac{1}{2}$
Isospin I is conserved in strong interactions
Addition by rules of angular momentum

Isospin Multiplets

Useful for classification of hadrons, see slide 1
 $2I+1$ states in a isospin multiplet $|I, I_3\rangle$

$$\eta = |0, 0\rangle$$

$$p = \left| \frac{1}{2}, \frac{1}{2} \right\rangle$$

$$n = \left| \frac{1}{2}, -\frac{1}{2} \right\rangle$$

$$\pi^+ = |1, 1\rangle$$

$$\pi^0 = |1, 0\rangle$$

$$\pi^- = |1, -1\rangle$$

$$\Delta^{++} = \left| \frac{3}{2}, \frac{3}{2} \right\rangle$$

$$\Delta^+ = \left| \frac{3}{2}, \frac{1}{2} \right\rangle$$

$$\Delta^0 = \left| \frac{3}{2}, -\frac{1}{2} \right\rangle$$

$$\Delta^- = \left| \frac{3}{2}, -\frac{3}{2} \right\rangle$$

Quark Model

Gives natural explanation for Isospin

$$I_3 = \frac{1}{2}(n_u - n_d + n_{\bar{d}} - n_{\bar{u}}) \quad n_i \text{ number of } i \text{ quarks}$$

Isospin works well

Masses of u and d quark are almost equal

Isospin Conservation

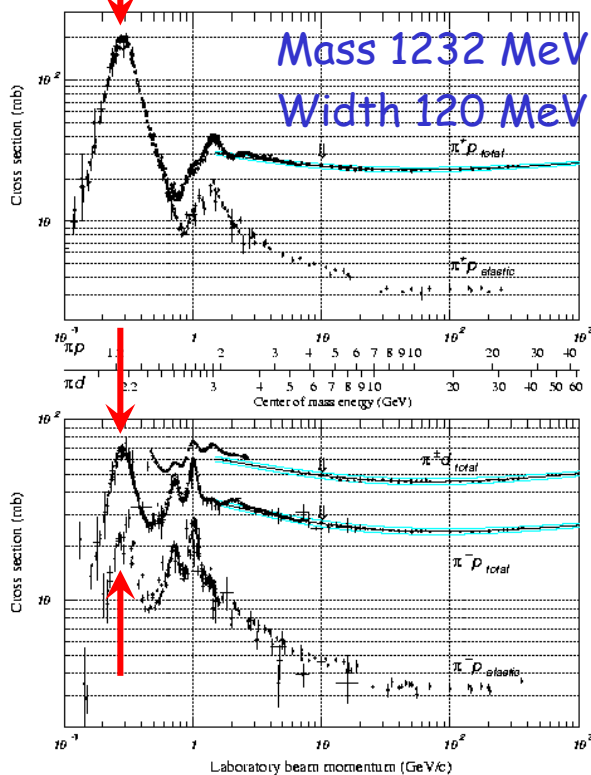


Conservation Law

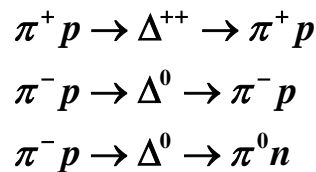
Isospin I is conserved in strong interactions
 Allows to calculate ratios of cross sections and branching fractions in strong interactions

Delta(1232) Resonance

36. Plots of cross sections and related quantities 17



Production



Isospin addition

$$\begin{aligned} \pi^+ p &: |1,1\rangle \left| \frac{1}{2}, \frac{1}{2} \right\rangle = \left| \frac{3}{2}, \frac{3}{2} \right\rangle \\ \pi^- p &: |1,-1\rangle \left| \frac{1}{2}, \frac{1}{2} \right\rangle = \sqrt{\frac{1}{3}} \left| \frac{3}{2}, -\frac{1}{2} \right\rangle - \sqrt{\frac{2}{3}} \left| \frac{1}{2}, -\frac{1}{2} \right\rangle \\ \pi^0 n &: |1,0\rangle \left| \frac{1}{2}, -\frac{1}{2} \right\rangle = \sqrt{\frac{2}{3}} \left| \frac{3}{2}, -\frac{1}{2} \right\rangle + \sqrt{\frac{1}{3}} \left| \frac{1}{2}, -\frac{1}{2} \right\rangle \end{aligned}$$

Matrix element

$$M_3 = \left\langle \frac{3}{2} \left| H_3 \right| \frac{3}{2} \right\rangle$$

depends on I, not I_3 $M_1 = \left\langle \frac{1}{2} \left| H_1 \right| \frac{1}{2} \right\rangle$

$$M(\pi^+ p \rightarrow \Delta^{++} \rightarrow \pi^+ p) = M_3$$

$$M(\pi^- p \rightarrow \Delta^0 \rightarrow \pi^- p) = \frac{1}{3} M_3 + \frac{2}{3} M_1$$

$$M(\pi^- p \rightarrow \Delta^0 \rightarrow \pi^0 n) = \frac{\sqrt{2}}{3} M_3 - \frac{\sqrt{2}}{3} M_1$$

$$\sigma(\pi^+ p \rightarrow \Delta^{++} \rightarrow \pi^+ p) \approx 200 \text{ mb} \approx 9x$$

$$\sigma(\pi^- p \rightarrow \Delta^0 \rightarrow \text{all}) \approx 70 \text{ mb} \approx 3x$$

$$\sigma(\pi^- p \rightarrow \Delta^0 \rightarrow \pi^- p) \approx 23 \text{ mb} \approx 1x$$

Cross sections

$$\sigma \propto |M|^2$$

In agreement with

I=3/2 Isospin prediction

Strangeness

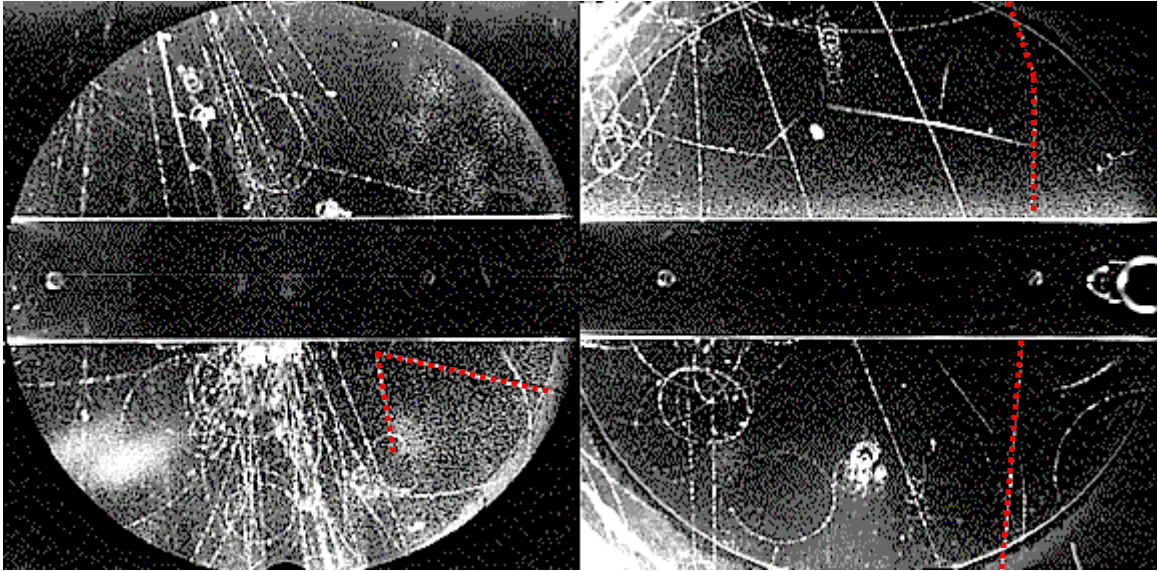


Strange Particles

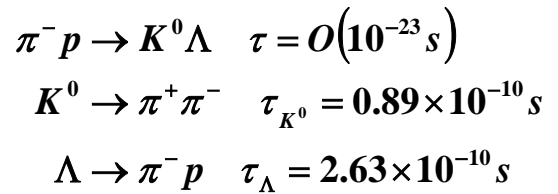
Discovered in 1947

Rochester and Butler

V, "fork", and K, "kink"



Production of $V(K^0, \Lambda)$ and K^\pm
via strong interaction,
weak decay



Associated Production

Strange particles produced in pairs

Pais

Strangeness S

Additive quantum number

Gell-Mann Nishijima

Conserved in strong and electromagnetic interactions

Violated in weak decays

Non-zero for Kaons and hyperons

$S = 0:$	π, p, n, Δ, \dots	$S = 1:$	K^+, K^0
$S = -1:$	$K^-, \bar{K}^0, \Lambda, \Sigma, \dots$	$S = -2:$	Ξ

Naturally explained in quark model $S = n_{\bar{s}} - n_s$

Quark Model



3 Quark Flavours u, d, s

1964 - introduced by Gell-Mann & Zweig

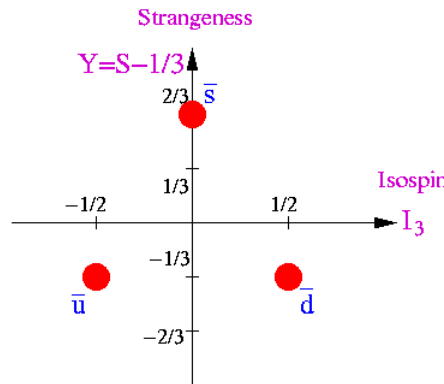
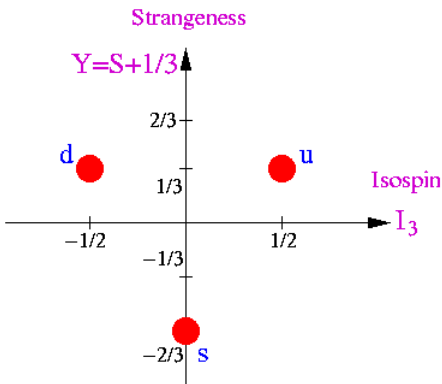
Quark	Charge $Q [e]$	Isospin $ I, I_3 \rangle$	Strange- ness S
up (u)	+2/3	$ \frac{1}{2}, +\frac{1}{2} \rangle$	0
down (d)	-1/3	$ \frac{1}{2}, -\frac{1}{2} \rangle$	0
strange (s)	-1/3	$ 0, 0 \rangle$	-1



Gell-Mann



Zweig



Charge, Isospin and Strangeness

Additive quark quantum numbers are related

$$Q = I_3 + \frac{1}{2}(S + B)$$

not all independent

Gell-Mann Nishijima predates quark model
valid also for hadrons

Baryon number B
quarks $B = +1/3$
anti-quarks $B = -1/3$

Hypercharge $Y = S + B$ is useful quantum number
Quark model gives natural explanation
for Isospin and Strangeness

Mesons



Bound $q\bar{q}$ States

Zero net colour charge

Zero net baryon number

$$|\psi\rangle = \frac{1}{\sqrt{3}} |r\bar{r} + g\bar{g} + b\bar{b}\rangle$$

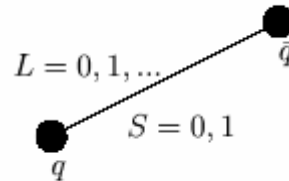
$$B = +1/3 + (-1/3) = 0$$

Angular Momentum \vec{L}

For lightest mesons

Ground state

$L = 0$ between quarks



Parity P

Intrinsic quantum number of quarks and leptons

$P=+1$ for fermions $P=-1$ for anti-fermions

$$P(q\bar{q}) = P_q P_{\bar{q}} (-1)^L$$

$$= (+1)(-1)(-1)^L = -1 \quad \text{for } L=0$$

Total Angular Momentum \vec{J}

$\vec{J} = \vec{L} + \vec{S}$ include quark spins

$S = 0$ $q\bar{q}$ spins anti-aligned $\uparrow\downarrow$ or $\downarrow\uparrow$

$\rightarrow J^P = 0^-$ **Pseudo-scalar mesons**

$S = 1$ $q\bar{q}$ spins aligned $\uparrow\uparrow$ or $\downarrow\downarrow$

$\rightarrow J^P = 1^-$ **Vector mesons**

Quark flavours

$u\bar{d}, u\bar{s}, d\bar{u}, d\bar{s}, s\bar{u}, s\bar{d}$

non-zero flavour states

$u\bar{u}, d\bar{d}, s\bar{s}$

zero net flavour states

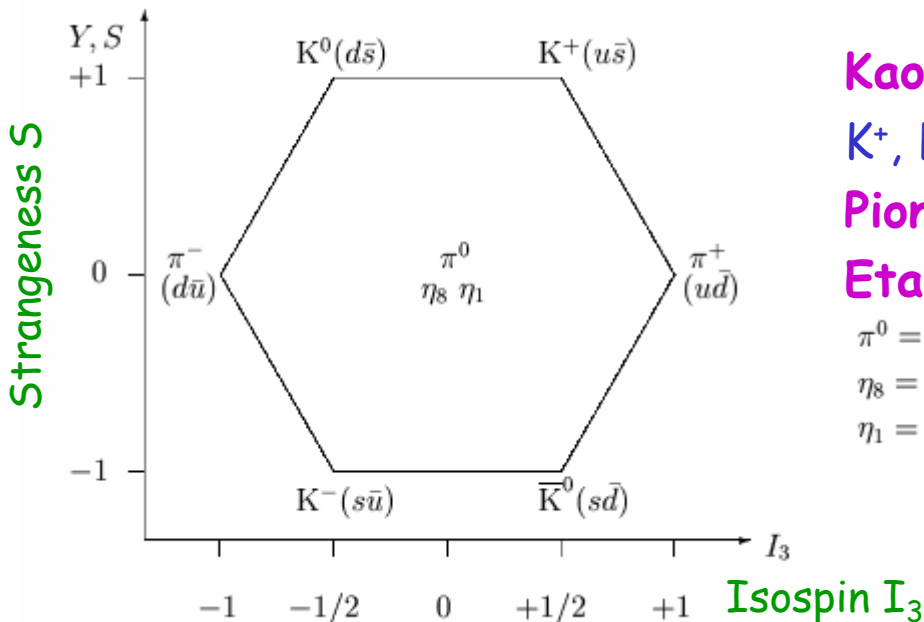
have identical additive quantum numbers

Physical states are mixtures

Mesons



Pseudoscalar Mesons $J^P = 0^-$



Kaons:

$K^+, K^0, \text{anti-}K^0, K^-$

Pions: π^+, π^0, π^-

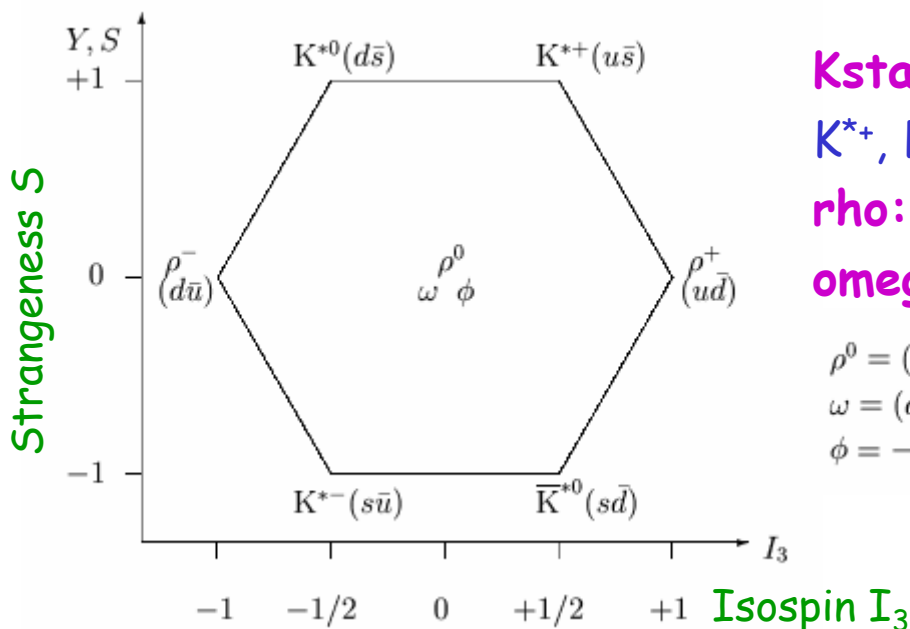
Etas: η, η'

$$\pi^0 = (d\bar{d} - u\bar{u})/\sqrt{2}$$

$$\eta_8 = (d\bar{d} + u\bar{u} - 2s\bar{s})/\sqrt{6}$$

$$\eta_1 = (d\bar{d} + u\bar{u} + s\bar{s})/\sqrt{3}$$

Vector Mesons $J^P = 1^-$



Kstar:

$K^{*+}, K^{*0}, \text{anti-}K^{*0}, K^{*-}$

rho: ρ^+, ρ^0, ρ^-

omega/phi: ω, ϕ

$$\rho^0 = (d\bar{d} - u\bar{u})/\sqrt{2}$$

$$\omega = (d\bar{d} + u\bar{u})/\sqrt{2}$$

$$\phi = -s\bar{s}$$

Baryon Decuplet



Baryon Wavefunction

$$\Psi(\text{total}) = \Psi(\text{space}) \Psi(\text{spin}) \Psi(\text{flavour}) \Psi(\text{colour})$$

Space symmetric - $L = 0$

Flavour symmetric, e.g. uuu , $(udu+duu+uud)/\sqrt{3}$

Spin symmetric

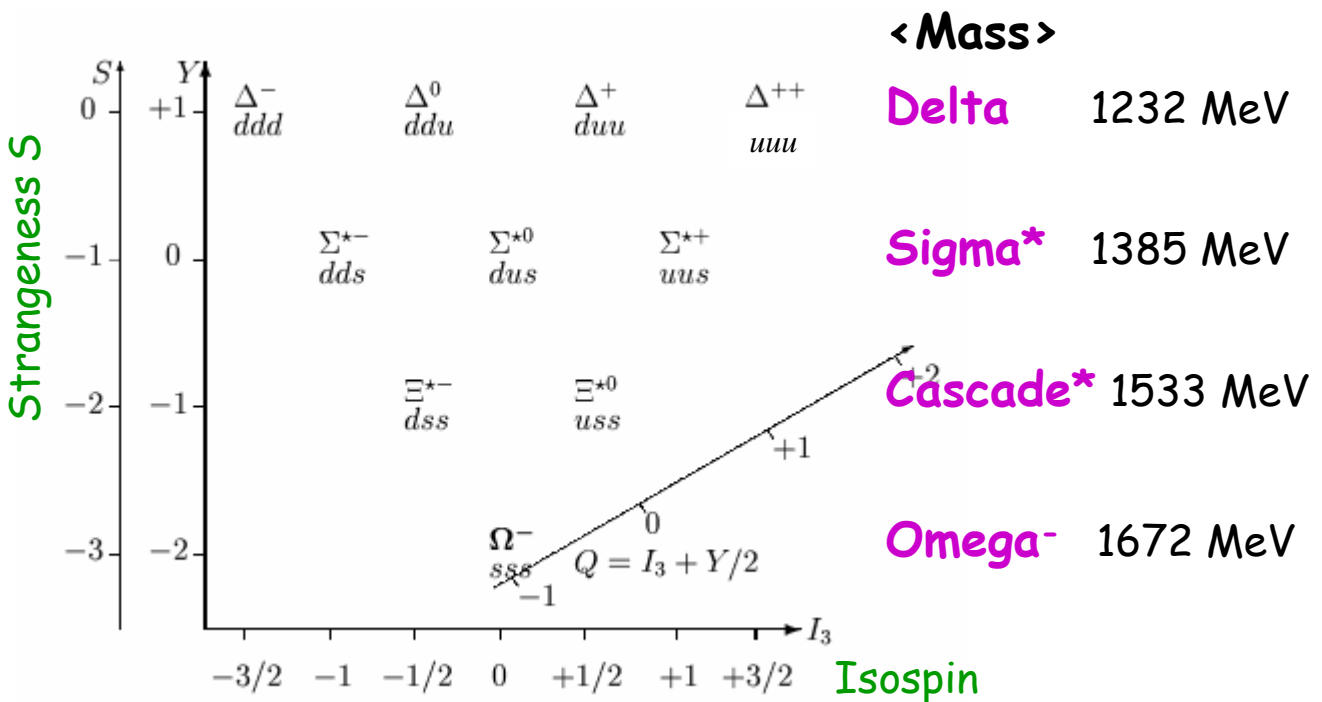
all 3 quarks aligned $\rightarrow S = 3/2$

Colour antisymmetric

$$(rgb - rbg + gbr - grb + brg - bgr)/\sqrt{6}$$

Total antisymmetric - obeys Pauli Exclusion Principle

Baryon Decuplet $J^P = 3/2^+$



Quark model predicted unobserved state Ω^- (sss)

Baryon Octet



Baryon Wavefunction

$\Psi(\text{space})$ symmetric ($L = 0$) $\Psi(\text{colour})$ antisymmetric

Mixed symmetric $\Psi(\text{spin, flavour})$

Flavour mixed symmetric: e.g. $(ud - du) u/\sqrt{2}$

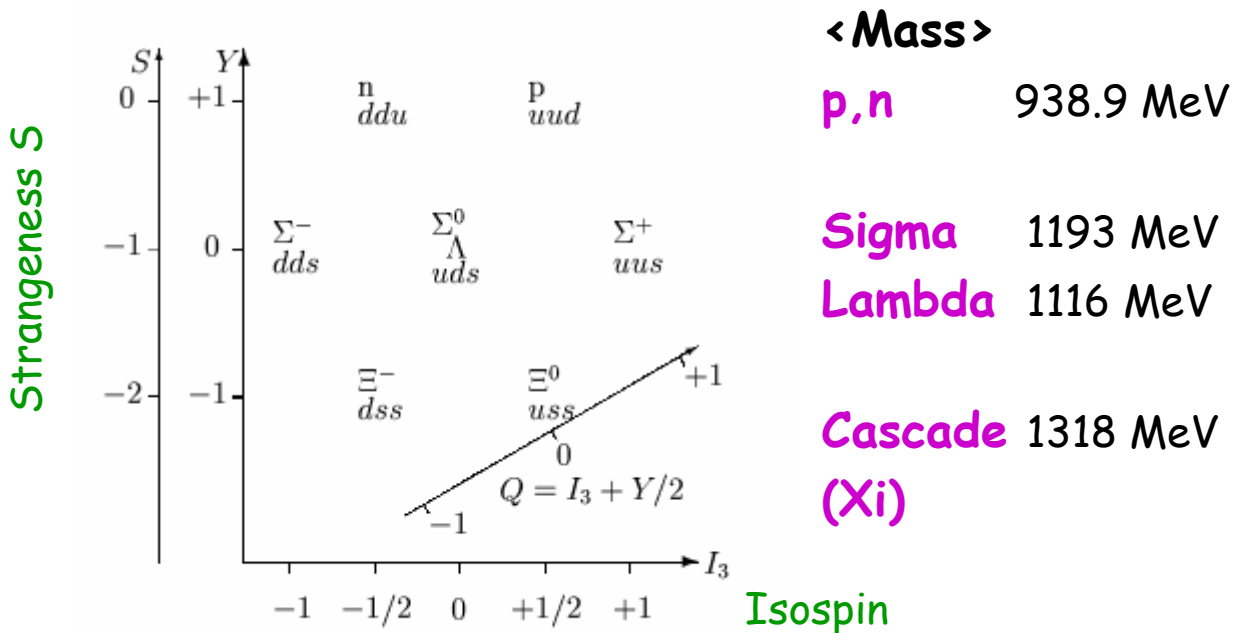
Spin as flavour: e.g. $(\uparrow\downarrow - \downarrow\uparrow) \uparrow/\sqrt{2}$

Spin-flavour e.g. $(u\uparrow d\downarrow - d\uparrow u\downarrow - u\downarrow d\uparrow + d\downarrow u\uparrow) u\uparrow/\sqrt{6}$

Symmetrisation by cyclic permutations

$$\Psi(\text{proton, } s=+\frac{1}{2}) = (2u\uparrow u\uparrow d\downarrow - u\uparrow u\downarrow d\uparrow - u\downarrow u\uparrow d\uparrow + 2d\downarrow u\uparrow u\uparrow - d\uparrow u\uparrow u\downarrow - d\uparrow u\downarrow u\uparrow + 2u\uparrow d\uparrow u\downarrow - u\uparrow d\downarrow u\uparrow - u\downarrow d\uparrow u\uparrow) / \sqrt{18}$$

Baryon Octet $J^P = \frac{1}{2}^+$



Lightest baryons

stable or long-lived

Antibaryons (\bar{p}, \bar{n}, \dots)

also form Octet and Decuplet

Discovery of Ω^-



Ω^- (sss) Hyperon

Hyperon - baryon with at least one s quark

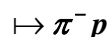
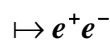
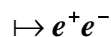
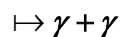
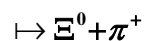
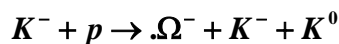
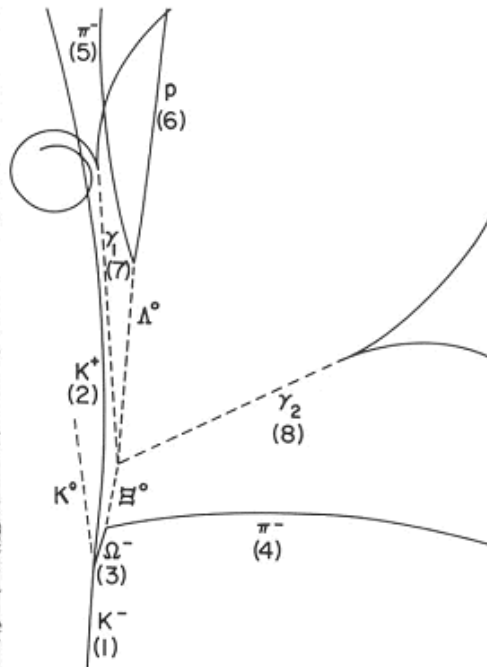
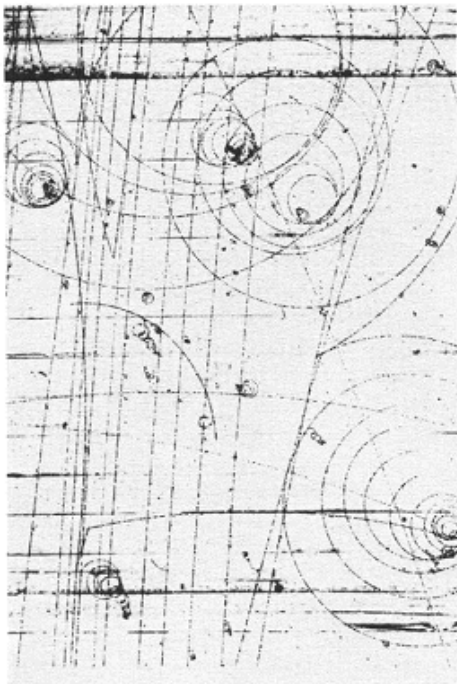
Quark model predicted existence and mass

Missing member of baryon decuplet $J^P = 3/2^+$

discovered 1964 at Brookhaven

K⁻ beam onto hydrogen target

Bubble Chamber detector



Hadron Masses



Quark Masses

u, d & s quark masses light at short distance

$$q^2 > 1 \text{ GeV}^2 \quad m_u < m_d \sim 5 \text{ MeV} \quad m_s \sim 100 \text{ MeV}$$

Constituent mass is relevant for quark model

$$q^2 < 1 \text{ GeV}^2 \quad m_u = m_d \sim 300 \text{ MeV} \quad m_s \sim 500 \text{ MeV}$$

Meson Masses

$m(K) > m(\pi)$ due to $m_s > m_u, m_d$

$m(\rho) > m(\pi)$ same quark content e.g. ρ^+, π^+ : (u-dbar)

Mass difference is due to quark spins

Chromomagnetic Mass Splitting

Spin-spin coupling of quarks $S_1 = S_2 = 1/2$

analogous to hyperfine splitting in el. mag. interaction

$$\Delta E \propto \alpha_s \frac{\vec{S}_1 \cdot \vec{S}_2}{m_1 m_2}$$

$$m(q\bar{q}) = m_1 + m_2 + A \frac{\vec{S}_1 \cdot \vec{S}_2}{m_1 m_2}$$

$$\vec{S}_1 \cdot \vec{S}_2 = \frac{1}{2} (\vec{S}^2 - \vec{S}_1^2 - \vec{S}_2^2) = \frac{1}{2} (S(S+1) - S_1(S_1+1) - S_2(S_2+1))$$

$$= \begin{cases} 1 - \frac{3}{4} = \frac{1}{4} & S=1 \\ 0 - \frac{3}{4} = -\frac{3}{4} & S=0 \end{cases}$$

Meson Masses

$$m_u = m_d = 310 \text{ MeV}$$

$$m_s = 483 \text{ MeV}$$

$$A = (2m_u)^2 \cdot 160 \text{ MeV}$$

Excellent agreement

What about eta(')?

	Mass [MeV]	
Meson	Prediction	Experiment
π	140	138
K	484	496
ρ	780	770
ω	780	782
K^*	896	894
ϕ	1032	1019

Heavy Quarks



Charm and bottom quarks

Charmonium (c-cbar) --- see QCD lecture

1977 Discovery of Upsilon States

Interpretation is

Bottomonium (b-bbar)

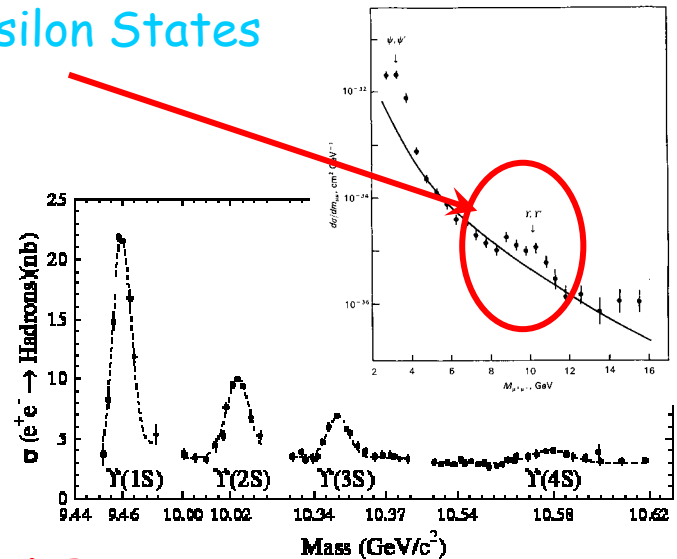
Spectroscopy

Charmonium

and Upsilon

$m_c \sim 1.1 \dots 1.4 \text{ GeV}$

$m_b \sim 4.1 \dots 4.5 \text{ GeV}$



Heavy-light Mesons and Baryons

Charmed (c-quark) hadrons

$$J^P = 0^- \quad D^0 = c\bar{u}, \quad D^+ = c\bar{d}, \quad D_s^+ = c\bar{s},$$

$$J^P = 1^- \quad D^{*0} = c\bar{u}, \quad D^{*+} = c\bar{d}, \quad D_s^{*+} = c\bar{s},$$

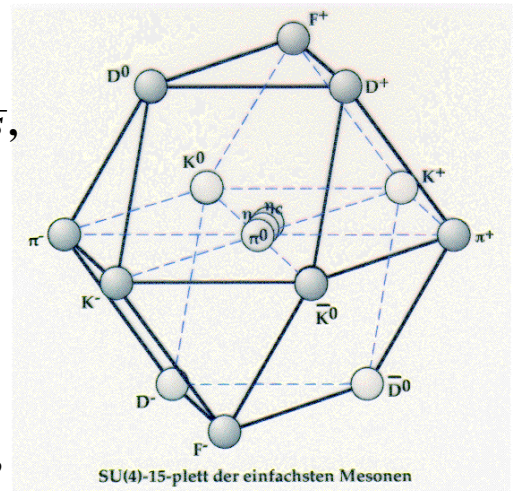
$$J^P = \frac{1}{2}^- \quad \Lambda_c^+ = cud$$

Bottom-quark hadrons

$$J^P = 0^- \quad B^+ = u\bar{b}, \quad B^0 = d\bar{b}, \quad B_s^0 = s\bar{b},$$

$$J^P = 1^- \quad B^{*+} = u\bar{b}, \quad B^{*0} = d\bar{b}, \quad B_s^{*0} = s\bar{b},$$

$$J^P = \frac{1}{2}^- \quad \Lambda_b^0 = bud$$



Top quark

Decays before forming bound states

$m_t \sim 174 \text{ GeV}$

discovered in 1995 at Fermilab