

Discrete Groups and Flavour Physics *

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CP3 Origins
<http://cp3-origins.dk/>

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The rare and the beautiful -- Wylerfescht

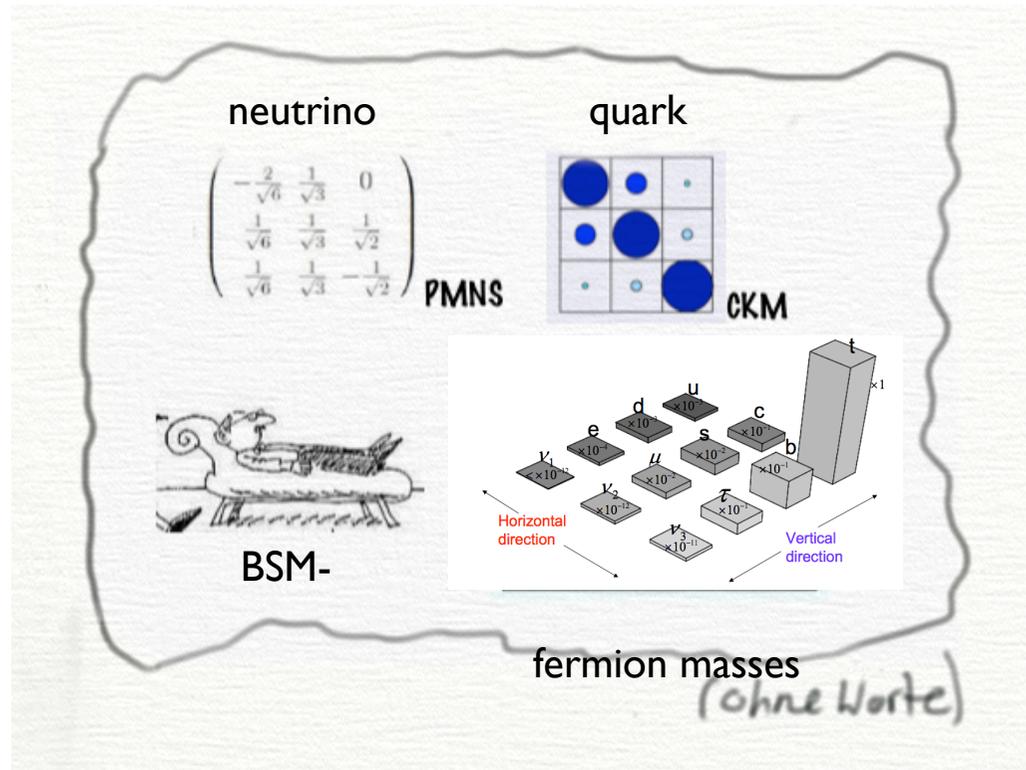
Overview

- Discrete groups -- discrete SU(3) subgroups
- A_4 an example -- 79' model building
02'.. tri-bi-mixing lepton sector

conceptual cut -- discrete groups link

- Discrete Minimal Flavour Violation

Why discrete groups? ... Theory of flavour?



Are these patterns random?

no

yes

- Can (discrete/continuous) **family-symmetries** explain them?
Who gave us the Yukawa matrices?
- Is the mechanism linked with the TeV-scale?
(SM suggests that in a way: why $m_t \sim \Lambda_{EW}$?)



Def: **discrete group** = group with countable many elements

two facts (for general orientation):

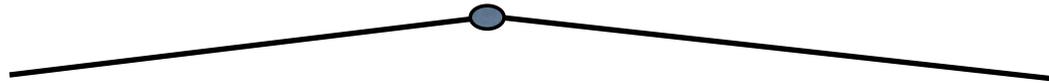
- Any discrete group can be embedded into permutation group S_n (analogue manifold & \mathbb{R}^n) *
- Order & irreducible representations (irreps): $|D| = \sum |\text{irrep}(D)|^2$ \Rightarrow finite many of them

Which discrete groups are subgroups of $SU(3)$? (3 because of 3 families)

* 19th century used to be the definition (as opposed to abstract definition)

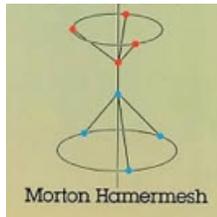
Discrete subgroups of SU(3)

- Classified in a classic book Miller, Dickson, Blichfeld '1916
 Analyzed further eightfold way Fairbairn, Fulton, Klink '64
 Further analyzed (lattice ...) Bovier, Luling, Wyler '81 *
 Rescrutinized tri-bi-hype Luhn, Nasri Ramand., '06-08



Trihedral like: $\Delta(3n^2) / \Delta(6n^2)$

- $Z_n \times Z_n \rtimes Z_3 / S_3$
- largest irreps 3/6-dim
- Analogue Dihedral group (chemistry) $Z_n \rtimes Z_2$

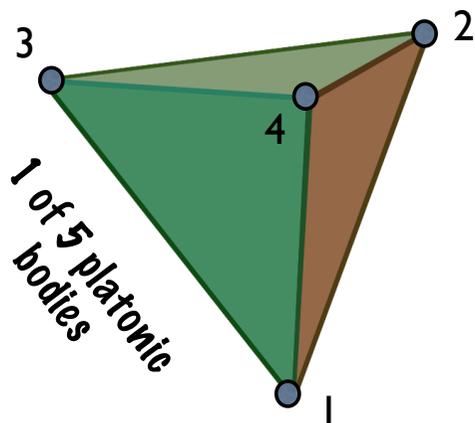


Crystallographic groups, Σ

- finite many of them
- maximal subgroups
 $\Sigma(168) \sim \text{PSL}(2,7)$
 $\Sigma(216\varphi)$ hessian group
 $\Sigma(360\varphi)$
 $(\varphi=1,3 \text{ related center } \text{SU}(3))$
- used in lattice $\text{SU}(3)_{\text{color}}$ discretizations '80

* some confusion (D)-groups, reemphasized Ludl'09, argue (D)-groups embedded in $\Delta(6n^2)$ Fischbacher RZ'09

Tetrahedral group $\cong A_4$



- 4 corners, fix one e.g. 4 120°-rotation: (4)(234) & (4)(243) = 8
 - 3 opposite edges with 180°-rotation: (12)(34) ... = 3
 - identity = 1
-
- elements 12

$\Rightarrow T \cong A_4$ (even four permutation) = $\Delta(12)$

- **Algebraic def:** $S^2=I, T^3=I, (ST)^3=I$
- **Irreps:** $|A_4| = 4!/2 = 12 = |I|^2 + |I'|^2 + |\bar{I}'|^2 + |3|^2$
- Example Kronecker product: $3 \times 3 = I + I' + \bar{I}' + 3_s + 3_a$
denote: $3 \sim (x_1, x_2, x_3)$ & $3 \sim (y_1, y_2, y_3)$ $I' \sim (\omega^2 x_1 y_1, \omega x_2 y_2, x_3 y_3)$ $\omega = \exp(2\pi i/3)$
; $T: I' \rightarrow \omega^2 I'$

A_4 in model-building

Model building 70'

Tri-bi-maximal mixing '00

Connection? ... they were just first in line ...

A₄ quark sector in the '70

Wyler'79

- Not much known about 3rd generation (basically m_b) -- $\Theta_c \sim (m_d/m_s)^{1/2}$ Cabibbo universality

- Assume: $\mathcal{L}_{\text{Yuk}} \sim \sum_n (I^n)_{abc} \bar{Q}_L^a D_R^b H^c$

a,b,c index family-symmetry
Iⁿ invariant (constant tensor)

- Need at least two invariants (o/w $m_d/m_u = m_s/m_c$) -- $\mathbf{3}^* \times \mathbf{3} = \mathbf{2} \oplus \mathbf{3}$... (not simply reducible)
- **A₄** candidate with low order ($\mathbf{3} \times \mathbf{3} = \mathbf{2} \oplus \mathbf{3} + \dots$) \Rightarrow 2 invariants \Rightarrow 2 Yukawas (instead of 3)
- Work out (m_d, m_s, m_b) $\langle H \rangle = (v, v, v_3)$ with $v \ll v_3$ (\Rightarrow Higgs potential add. family singlet)
third mass adjusting the 2 Yukawas

Results:

- Cabibbo universality ✓
- $|V_{ub}|/|V_{cb}| \sim 10$ (reversed hierarchy)
- 2 Yukawa & 3 masses \rightarrow relation: $m_d m_s / m_b^2 = m_u m_c / m_t^2 \Rightarrow m_t \approx 15 \text{ GeV}$

(not known '79)

Tri-bi-maximal mixing & A_4 or rather S_4

Data suggests (not exclude):

$$U_{\text{TB}} = \begin{pmatrix} \sqrt{2/3} & 1/\sqrt{3} & 0 \\ -1/\sqrt{6} & 1/\sqrt{3} & -1/\sqrt{2} \\ -1/\sqrt{6} & 1/\sqrt{3} & +1/\sqrt{2} \end{pmatrix}$$

Harrison, Perskins, Scott '99'02

- Neutrino sector: 'know' mixing -- masses less known
- Go into basis leptons are diagonal

$$\begin{aligned} M_\nu &= M_{\text{TB}}^\nu \equiv U_{\text{TB}} \text{diag}(m_{\nu_1}, m_{\nu_2}, m_{\nu_3}) U_{\text{TB}}^T \\ M_l &= \text{diag}(m_e, m_\mu, m_\tau) \end{aligned}$$

- $AM_\nu A^\dagger = M_\nu$ for $A \in \mathbf{Z}_2 \times \mathbf{Z}_2$ generated S, U
- $AM_l A^\dagger = M_l$ for $A \in \mathbf{Z}_3$ generated by T
- The three generators S, T, U define \mathbf{S}_4 (and not A_4 ... but was a good start)
Suggests (original) family symmetry S_4 (or any group containing S_4 e.g. $\text{PSL}(2,7)$...)
- Model building **flavon** ϕ_T with T-invariant VEV, Frogatt-Nielsen etc ..

$$\mathcal{L}_{\text{YUK}}^L \sim \Psi(\phi_T + \phi_0)\Psi^c$$

Minimal Flavour Violation

Is there something (very) special about the Yukawa matrices?

Minimal Flavour Violation

- Yukawa** = 0 continuous global symmetry: $G_F = U(3)^5 = G_q \times G_l$, $G_q = U(3)_Q \times U(3)_{UR} \times U(3)_{DR}$
Yukawa $\neq 0$ breaks down to: $G_q = U(3)_q^3 \rightarrow U(1)_{\text{Baryon}}$
- Let **Yukawa** formally transform as $Y_D \sim (3^*, 1, 3)_{G_q}$ & $Y_U \sim (3^*, 3, 1)_{G_q}$

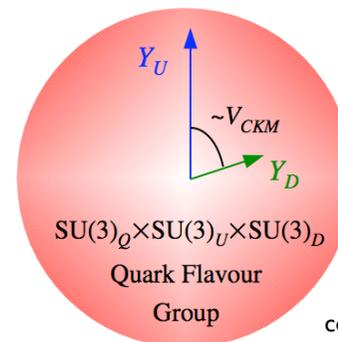
$\Rightarrow G_F$ symmetry restored

MFV: effective field theory invariant under global G_F
(criterion of naturalness applied coefficient $O(1)$)

D'Ambrosio, Giudice, Isidori & Strumia '02

- Yukawa's** promoted to spurions
 $\langle Y_{U/D} \rangle \neq 0$ VEV breaks G_q

N.B. interpretation of symmetry breaking
 other options a) explicit(soft) breaking
 b) anomalous breaking



courtesy G.Isidori

CKM due to mis-alignment of Yukawa's

* Add. assumption: CP-invariance
 No new Lorentz structure

AGIS'02
 Buras, Gambino, Gorgahn Jager, Silvestrini ... '00

A few remarks on MFV

(partly skip)

- Restricts the number of operators: (denote: $D = (d,s,b)$)

Dynamics without dynamics

$$O^{\Delta F=1'} = (\bar{D}_L Y_U Y_U^\dagger Y_D \sigma \cdot F D_R)$$

$$b \rightarrow s\gamma \text{-type}$$

$$O^{\Delta F=1} = (\bar{D}_L Y_U Y_U^\dagger D_L) \cdot \bar{D}_L D_L$$

$$B \rightarrow K\pi\pi$$

$$O^{\Delta F=2} = (\bar{D}_L Y_U Y_U^\dagger D_L)^2$$

$$B_d \rightarrow \bar{B}_d \text{-type}$$

- \Rightarrow **correlations**: $b \rightarrow s$, $b \rightarrow d$, $s \rightarrow d$ transitions e.g. $\Delta M_d / \Delta M_s$ as in SM

testable hypothesis!

- Is there still room for large effects? Yes in certain channels e.g. $B \rightarrow ll$ -- enhancement due to large $\tan\beta = v_u / v_d$

- SUSY & MFV make proton long lived! No need for R-parity!

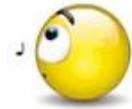
Nikoladitis & Smith '07

- No model of MFV (... seems as hard as creating a **Theory of Flavour**)

- MFV is also a language -- you can compare your BSM-flavour physics to MFV

If we take it beyond that

Consequences of spurious Goldstone bosons



- If G_q -symmetry SSB by Yukawa's $\langle Y_{U/D} \rangle \neq 0 \Rightarrow 3 \times 8 + 2 = 26$ (massless) Goldstone bosons
- Mid 70's 80's study breaking of continuous family symmetries -- dubbed Goldstone modes **familons**
- Physics should be the same:

$$\mathcal{L}^{\text{eff}} \simeq \frac{1}{\Lambda_F} (\partial_\mu \phi_F) (\bar{s} \gamma_\mu d_L) + \dots$$

Φ_F weakly coupled (not detected) ---- $K^+ \rightarrow \pi^+ \nu \nu$ vs $K^+ \rightarrow \pi^+ \Phi_F$
 $\Rightarrow \Lambda_F > 10^8 \text{ TeV}$ (infer from e.g. Feng et al '97)

- If $\Lambda_F \sim \Lambda_{\text{MFV}}$ flavour difficult to detect

caveat:

If flavour violation in soft-breaking terms
 e.g. SUSY-GUT then $\Lambda_F \sim \Lambda_{\text{GUT}} \dots \Lambda_{\text{MFV}} \sim \Lambda_{\text{SUSY}}$

ways out

- **Discrete symmetry**
 (no Goldstone modes)

this work ..

- **Gauge the symmetry**
 (new massive gauge bosons)

Albrecht, Feldmann, Mannel 'announced'

**There are (plenty) discrete subgroups
very good -- end of the day?**

By going to a discrete symmetry

1. Get rid of familons (goldstone bosons)
2. Reduce symmetry \Rightarrow new flavour structure (dangerous?)



Is there room for a TeV-scale dMFV-scenario ?

Formulation: discrete Minimal Flavour Violation (dMFV)

Fischbacher RZ 2008 PRD79

1. $G_q \rightarrow D_q = \mathbf{D3}_Q \times \mathbf{D3}_{UR} \times \mathbf{D3}_{DR}$ $D3 \subset SU(3)$, not discuss $U(1)$'s could be Z_n
2. Specify the 3D **irrep** of $D3$
3. (possibly) Yukawa expansion $Y \rightarrow \kappa Y \quad \kappa \leq 1$ *

- Model independent approach: \Rightarrow study of **invariants** (\sim effective operators)

$$\mathcal{L}^{\text{eff}} \sim \sum_n \frac{c_n}{\Lambda^{\dim(\mathcal{I}_n)-4}} \mathcal{I}_n(u, d, Y_U, Y_D) + h.c.$$

- Cutting a long story short: new invariants = non-MFV transitions
 1. New invariants (typically) \Rightarrow anarchic flavour transitions
 2. Moreover: **Yukawas** (modulo CKM) diagonalized via G_q
If we break it down to D_q (\Rightarrow more observable mixing angels!)

* $\kappa \equiv 1$ non-linear MFV (σ -model ... [Feldmann, Mannel '08](#), [Kagan et al '09](#)) $\kappa \ll 1$ linear MFV

Invariants

Let $\mathbf{A}, \mathbf{B}, \mathbf{C}, \dots$ be irreps of some group then $\mathbf{A} \times \mathbf{B} \times \mathbf{C} \times \dots = n\mathbf{1} + \dots$

n: Number of invariants = Number times $\mathbf{1}$ appears in (follows from $\mathbf{V} \times \mathbf{V}^* = \mathbf{1} + \dots \Leftrightarrow \mathbf{V}$ irrep)

- Denote tensor n $\mathbf{3}$ and m $\mathbf{3}^*$ indices by $T^{(m,n)}$ $D \equiv (d, s, b)_L \in T^{(1,0)}$, $\Delta \equiv Y_U^\dagger Y_U \in T^{(1,1)}$ *

dMFV

MFV

$$\begin{array}{l} \Delta F = 1' \\ \Delta F = 2 \end{array} \quad \begin{array}{l} \mathcal{I}_n^{(2,2)} = (\mathcal{I}_n)^{ab}_{rs} (\bar{D}^r \Delta_a^s D_b) \\ \mathcal{I}_n^{(4,4)} = (\mathcal{I}_n)^{abcd}_{rstu} (\bar{D}^r \Delta_a^s D_b) (\bar{D}^t \Delta_c^u D_d) \end{array} \quad \begin{array}{l} (\bar{D}_L \Delta_U Y_D \sigma \cdot F D_R) \\ (\bar{D}_L \Delta_U D_L)^2 \end{array}$$

Q: Are there $SU(3)$ subgroups with no new a: $\mathbf{1}^{(4,4)}$ b: $\mathbf{1}^{(2,2)}$ -invariants?

*assuming D_{UR} indices can be contracted

.... results on invariants:

- a: **no** (no discrete SU(3) subgroups no new $\mathbf{I}^{(4,4)}$ -invariants!)
- b: **yes** (are discrete SU(3) subgroups no new $\mathbf{I}^{(2,2)}$ -invariants!)

group	order	pairs ($\mathbf{3}, \bar{\mathbf{3}}$)	$\mathcal{I}^{(2,2)}$	$\mathcal{I}^{(3,3)}$	$\mathcal{I}^{(4,4)}$
SU(3)	∞	1	2	6	23
$\Sigma(360\varphi)$	1080	2	2	6	28
$\Sigma(216\varphi)$	648	3	2	7	40
$\Sigma(168)$	168	1	2	7	44
$\Sigma(72\varphi)$	216	4	2	11	92

subgroup

- Show: absence 27-dim irrep \Rightarrow new $\mathbf{I}^{(4,4)}$ -invariants (c.f. backup-slide)
 absence 8-dim irrep \Rightarrow new $\mathbf{I}^{(2,2)}$ -invariants

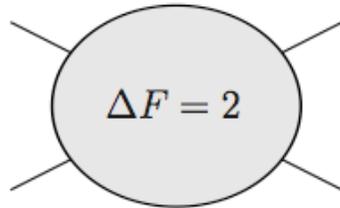
Q: Does this mean there's no TeV-scale dMFV?

Have to refine notion of model independence.

The $\Delta F=2$ generation mechanism has to be reflected upon.

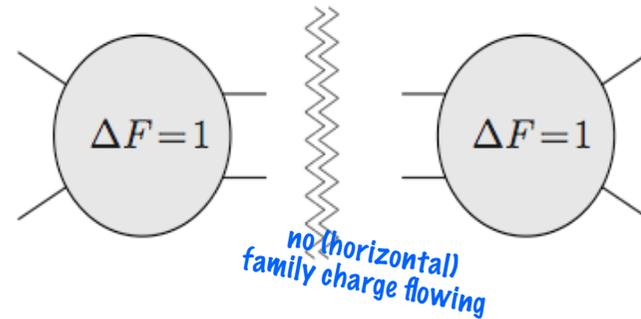
“Family (ir)reducible”

“family **ir**reducible” $I^{(4,4)}$

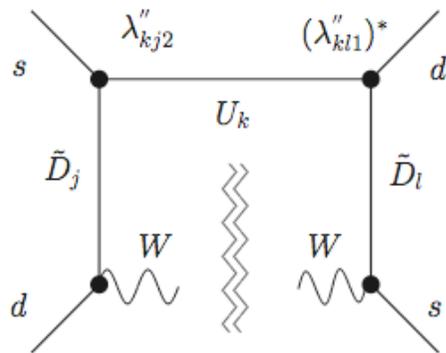


“family reducible” $I^{(4,4)} \rightarrow I^{(2,2)} \times I^{(2,2)}$

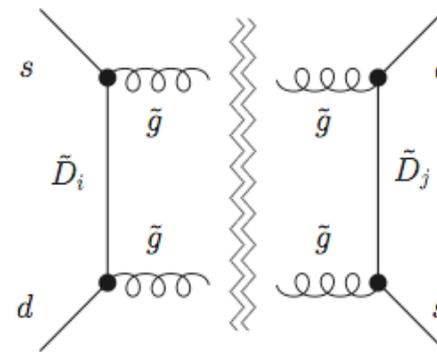
factorization



- Composite TC model, R-parity violating MSSM



- SM, R-parity conserving MSSM:



TeV-scale dMFV scenario

“family reducibility” is sufficient property for “TeV-scale dMFV scenario” for D_q :

$$\Sigma(168), \Sigma(72\varphi), \Sigma(216\varphi) \text{ and } \Sigma(360\varphi)$$

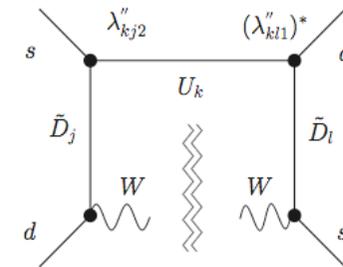
Dangerous invariants factorize: $\mathbf{I}^{(4,4)} \rightarrow \mathbf{I}^{(2,2)} \mathbf{I}^{(2,2)}$

- TeV-scale? Recall: $C_{SM}/C_{MFV} \geq (0.5 \text{ TeV}/m_W)^2$ -- Yukawa expansion: what κ bound $C_{MFV} \sim C_{dMFV}$?
 - most suitable candidate $\Sigma(360\varphi)$ only $\mathbf{I}^{(4,4)}$ new invariants
 - MFV**: $s \rightarrow d$ $O(\lambda^5)$ strong suppression, $\Delta S = 2$ real part $O(\lambda^{10})$
 - dMFV**: $s \rightarrow d$ $O(\lambda)$ (from examples ..“worst case”)

$$\Rightarrow \mathbf{MFV} : \mathbf{dMFV} = \lambda^{10} : \lambda^6 \kappa^4 \text{ equal} \quad \kappa_{\Sigma(360\varphi)} \approx \lambda \approx 0.2$$

- sufficient but not necessary! Consider R-parity violating MSSM
Can convince yourself that not lead to “dangerous” non-fac. $\mathbf{I}^{(4,4)}$

(essentially) each vertex is D_q -invariant!



Epilogue

- Discrete groups are fun (and have potential)
- Would be good to work out in more generality how breaking patterns
 $G_{\text{cont}} \rightarrow G_{\text{discrete}} \rightarrow G'_{\text{discrete}}$ works out (systematically)
- TeV-scale dMFV scenario possible for crystal-like groups $\Sigma(360\varphi)$, $\Sigma(216\varphi)$, $\Sigma(72\varphi)$, $\Sigma(168)$
(with moderate ($\kappa \sim O(0.2)$) possible (model-independent))
- *Happy Birthday* greetings from Bob Shrock, Pasquale Di Bari ,



Backup Slides

Necessarily new $I^{(4,4)}$ -invariants !!

- **What level?:** Kronecker product decompose any different than SU(3) !

Look at:
$$\mathcal{I}_n^{(4,4)} = (\mathcal{I}_n)_{rstu}^{abcd} (\bar{D}^r \Delta_a^s D_b) (\bar{D}^t \Delta_c^u D_d)$$



- SU(3): $(\mathbf{3} \times \mathbf{3}^* \times \mathbf{3} \times \mathbf{3}^*)_S \times (\mathbf{3} \times \mathbf{3}^* \times \mathbf{3} \times \mathbf{3}^*)_S =$
 $(\mathbf{8} \times \mathbf{8})_S \times (\mathbf{8} \times \mathbf{8})_S + \dots =$
 $(\mathbf{1} + \mathbf{8} + \mathbf{27}) \times (\mathbf{1} + \mathbf{8} + \mathbf{27}) + \dots$

\Rightarrow if $D_Q \subset \text{SU}(3)$ has no **27** \Rightarrow new $I^{(4,4)}$ invariants (new $\Delta F = 2$ structure)

- 1. **Dihedral groups** $\Delta(3n^2)$, $\Delta(6n^2)$ max 3,6D irrep \Rightarrow out
- 2. **Crystallographic groups** .. look at character tables reveals there is none (N.B. $27^2 = 729$ almost saturates the largest group ($3 \times 360 = 1080$) ...)