Subatomic Physics:

Particle Physics Handout 8

The Weak Force



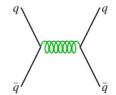
- Weak interactions
- * W and Z interactions at low energy
- * Fermi theory
- * Electroweak theory
- * W and Z bosons at high energy

QCD Summary

QCD: Quantum
Chromodymanics is the quantum description of the strong force.

Only quarks feel the strong force.

Gluons are the propagator of the strong force



Change solf interest

colour charge.

Quarks and gluons carry

- Gluons self-interact:
- Electromagnetic coupling constant α decreases as a charged particles get further apart.
- Strong coupling constant α_S increases as further apart quarks become.

Hadrons can be described as consisting of partons: quarks and gluons, which interact independently

Colour Confinement energy required to separate quarks $\rightarrow \infty$ quarks are confined to hadrons

Quarks and gluons produced in collisions hadronise: hadrons are produced.

000000

The decay products of the hadrons appear in the detector as **jets**.

2

Introduction to the Weak Force

The weak force is responsible for some of the most important phenomena:

- Decays of the muon and tau leptons
- Neutrino interactions
- Decays of the lightest mesons and baryons
- Radioactivity, nuclear fission and fusion

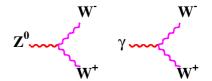
Characteristics of Weak Processes:

- Long lifetimes 10⁻¹³ 10³ s
- Small cross sections 10⁻¹³ mb

Boson	₩±	Z^0
Mass GeV/c ²	80.4	91.2
charge, e	±1	0
spin	1ħ	1ħ

Weak Force is propagated by massive W^+ , W^- and Z^0 bosons

- The interactions of W^{\pm} and Z^{0} are different (related by symmetry of the weak interaction)
 - W^{\pm} and Z^{0} can interact with each other
 - W^{\pm} and γ can interact (as W^{\pm} bosons are charged)



3

Weak Vertices

QED	W-boson	
mediated by the exchange of virtual photons	mediated by the exchange of $\it W$ boson	
acts on all charged particles	acts on all quark and leptons	
coupling strength $\propto e \propto \sqrt{\alpha}$	coupling strength $\propto g_W \propto \sqrt{\alpha_W}$	
propagator term: $1/(q^2-m_{\gamma}^2)=1/q^2$	propagator term: $1/(q^2 - m_W^2)$	
For many processes: $\mathcal{M}_{\sim} e^2/q^2$	For many processes: $\mathcal{M}_{\propto} g_{W^2}/(q^2-m_{W^2})$	
e e r	e - g _w V _e W.	

Recall: matrix element, \mathcal{M} , is the amplitude of a process. Scattering cross section, $\sigma \propto \mathcal{M}^2$. Decay width, $\Gamma \propto \mathcal{M}^2$

Interactions of the W^{\pm} boson

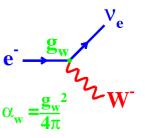
- Known as "charged current interactions"
- W^{\pm} boson interacts with all fermions (all quarks and leptons)
- Charged current changes the flavour of the fermion:
 - *e.g.* electron emitting an *W*-boson can't remain an electron violates conservation of charge!

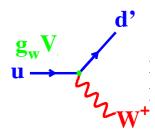


 an up quark turns into a down quark and vice versa!



- ullet Propagator term describing the W-boson $\propto \overline{(\underline{q}^2-m_W^2)}$
 - ullet is the four-momentum transferred by the W-boson





5

Allowed Flavour Changes

At a W-boson vertex:

- Lepton numbers: L_e , L_μ and L_τ , is conserved: Allowed lepton flavour changes: $e^- \leftrightarrow v_e \quad \mu^- \leftrightarrow v_\mu \quad \tau^- \leftrightarrow v_\tau$
- ullet Total Quark Number, N_q , is conserved
- Individual quark flavour numbers: $N_{\rm u}$, $N_{\rm d}$, $N_{\rm s}$, $N_{\rm c}$, $N_{\rm b}$, $N_{\rm t}$ are **not** conserved

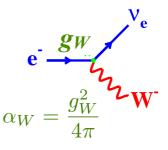
Allowed quark flavour changes:

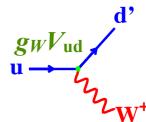
$$(Q=+2/3 e \text{ quark}) \leftrightarrow (Q=-1/3 e \text{ quark})$$

 $(d s b) \leftrightarrow (u c t)$

- Each of the nine possible quark flavour changes has a different coupling strength: e.g. gwV_{ud} for u to d quarks
 (Vs are terms in CKM matrix more later)
- Main quark flavour changes are within a generation:

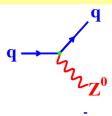
$$d \leftrightarrow u \quad s \leftrightarrow c \quad b \leftrightarrow t$$

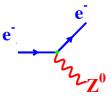


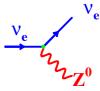


Interactions of the Z^0 boson

- Known as "neutral current interactions"
- Acts on all fermions (all quarks and leptons)
- Neutral current conserves flavour of the fermion
- No allowed fermion flavour changes
- \bullet Propagator term $\propto \frac{1}{\left(\underline{q}^2 m_Z^2\right)}$
- Coupling strength depends on fermion flavour we won't consider this in this course







Anywhere a photon could be exchanged a \mathbb{Z}^0 boson can be exchanged. (Almost vice-versa, except \mathbb{Z}^0 boson also has neutrino interactions too!)

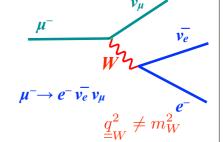
Electromagnetic and weak neutral current interactions are linked!

7

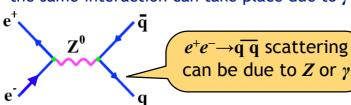
W and Z bosons at low energy

• If the momentum of the $\it W$ or $\it Z$ boson, $\it \underline{q} \ll m_Z, m_W$ the bosons are virtual: $\it \underline{q}^2_Z = \it E^2_Z - \vec{p}_Z \cdot \vec{p}_Z \neq m_Z^2$

$$\begin{array}{ll}
\stackrel{=}{=}Z & Z & IZ & IZ \\
\underline{q}^2 & = & E_W^2 - \vec{p}_W \cdot \vec{p}_W \neq m_W^2
\end{array}$$

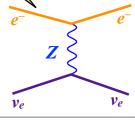


- Virtual *W*-bosons are responsible for the decays of the leptons, and the lightest hadrons e.g. $\mu^- \rightarrow e^- v_e^- v_\mu$
- virtual W-boson interactions can also be described by Fermi Theory
- Many interactions of virtual Z-bosons are not clearly evident, as the same interaction can take place due to γ exchange.



Interactions involving v must be due to Z (or W)

- At low energies we see the effect of Z boson mainly in scattering involving neutrinos e.g. $v_e e^- \rightarrow v_e e^-$
 - as γ cannot couple to the neutral neutrinos



Fermi Theory

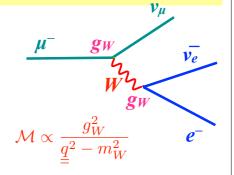
W-boson interactions at Low Momentum Transfer

For muon decay, and many other weak processes:

$$\mathcal{M} \propto rac{g_w^2}{(\underline{\underline{q}^2 - m_W^2})}$$

At low momentum transfer $\underline{q}^2 \ll m_W^2$

$$\mathcal{M}
ightarrow \propto \frac{g_w^2}{m_W^2}$$



Introduce Fermi coupling constant:

$$G_F \propto rac{g_w^2}{m_W^2} ~~G_F = rac{\sqrt{2}\,g_w^2}{8\,m_W^2}$$

- Dimension $[E]^{-\frac{1}{2}}$
- From experimental measurements: $G_F = 1.16637 \times 10^{-5} \text{ GeV}^{-2}$

Measurements of
$$G_F \& M_W \Rightarrow g_w = 0.66 \Rightarrow \alpha_w = \frac{g_w^2}{4\pi} = \frac{1}{29} > \alpha_{\rm EM} = \frac{1}{137}$$

• Recall, from problem sheet 2, Q7, range of W boson:
$$\Delta x \approx \frac{\hbar}{\Delta p} = \frac{\hbar}{m_W c} = 0.002 \ {\rm fm}$$

Weak interaction not intrinsically weak - appears weak due to large boson masses.

9

W interactions with quarks and leptons

Lepton universality:

- Coupling of to W-boson to all leptons is equal = g_W
- electrons, muons and taus all interact identically
- interact with the same bosons with same coupling strength

Quark interactions:

- In general, any vertex W-(Q=+2/3 e quark)-(Q=-1/3 e quark) is valid.
- ullet W-boson coupling to quarks suppressed by a flavour-dependent factor V

u d What gwVud	$\frac{\mathbf{u}}{W^{2}}\frac{\mathbf{s}}{g_{W}V_{\mathbf{u}\mathbf{s}}}$	u b Wh gwVub
c d Wh gwVcd	$\frac{c}{W^2} \frac{s}{g_W V_{cs}}$	$\frac{c}{W^2}$ g_WV_{cb}
t d Wh gwVtd	$\frac{t}{W^2} \frac{s}{g_W V_{ts}}$	$\frac{\mathbf{t}}{W^{2}}$ $g_{W}V_{tb}$

Known as the CKM matrix. (values from experimental measurements)

V _{ud} =0.974	V _{us} =0.227	V _{ub} =0.004
$V_{\rm cd} = 0.230$	$V_{\rm cs} = 0.972$	$V_{\rm cb} = 0.042$
$V_{\rm td} = 0.008$	$V_{\rm ts} = 0.041$	V _{tb} =0.999

Largest couplings are within a generation:

$$d \leftrightarrow u \quad s \leftrightarrow c \quad b \leftrightarrow t$$

Electroweak Theory

We've seen already that wherever a γ boson can be exchanged a Z can also be exchanged:

- The weak and electromagnetic force are linked.
- At short distances (or high energies) the strength of the electromagnetic force and the weak force are comparable. Can be related by a parameter, $\sin \theta_W$

$$e = g_W \sin \theta_W$$

The weak and electromagnetic interactions are manifestations of a underlying force: the electroweak force.

- ullet Couplings of the γ, W (and Z) bosons are related: $e=g_W\sin heta_W$
- ullet Mass of the $extit{\it W}$ and $extit{\it Z}$ bosons are related: $m_Z^2 = m_W^2/\cos^2 heta_W$

Just three fundamental parameters are required to describe:

- couplings of W, Z and γ to quarks and leptons
- masses of the W, Z, γ bosons
- interactions of the *W*, *Z*, γ bosons with each other

Normally formulated in terms of most accurately measured parameters: e, G_F , m_Z

11

W and Z boson at high energies

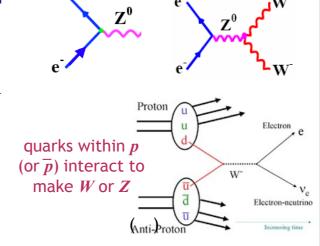
Using a collider, we can create high enough energies to make real W and Z bosons:

$$\underline{\underline{q}} \sim m_Z, m_W$$

- e.g. LEP collider $e^+e^- \rightarrow Z$, $e^+e^- \rightarrow W^+W^-$
- e.g. at the LHC: $pp \rightarrow Z+X$, $pp \rightarrow W+X$

Can study the properties of the W and Z bosons in detail.

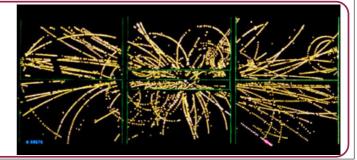
• masses, lifetime/width, coupling strengths, decay modes, spin...



Discovery of \overline{W} **and** \overline{Z} **bosons at CERN in 1983 at the SppS collider**

$$p \, \overline{p}$$
 collider with $E_p = E_{p^-} = 270 \, \text{GeV}$

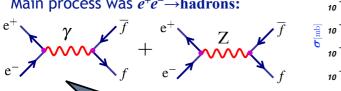
 $p \, \overline{p} \rightarrow W^- \rightarrow e^- \, \overline{v_e}$ event at UA1 experiment

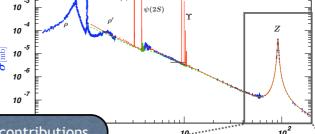


Measuring Z^{θ} properties

Most studies of the Z boson were made at LEP.

• Main process was $e^+e^- \rightarrow hadrons$:





Mixture of EM (γ) & Weak (Z) contributions

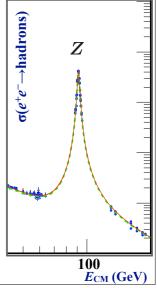
- At low $E_{\rm CM}$ mainly γ interactions.
- At $E_{\rm CM} = \sqrt{s} \sim m_Z$, mainly Z-boson interactions:

$$\sigma(e^+e^- o Z o f\overline{f})\propto rac{1}{({ar q}^2-m_Z^2)^2}$$

 Aside: for an massive boson, shape of cross section also depends on total decay width $\Gamma_Z = \hbar/\tau_Z$:

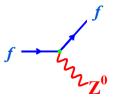
$$\sigma(E = \underline{\underline{q}}^2) = \sigma_{\text{max}} \frac{\Gamma_Z^2/4}{(\underline{\underline{q}}^2 - m_Z)^2 + \Gamma_Z^2/4}$$

- Measurements of $\sigma(e^+e^-\rightarrow hadrons)$ at LEP determined:
 - $m_Z = 91.188 \pm 0.002 \text{ GeV}/c^2$
 - $\Gamma_Z = 2.4953 \pm 0.002 \text{ GeV}$



Interactions of the Z boson

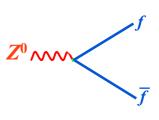
Z boson interacts with all quarks and leptons. No change of quark or lepton flavour at **Z** boson vertex.



Z boson decay

- Z can decay into any fermion-anti-fermion pair, $f\overline{f}$.
- except top quark pair, as $m_Z < 2m_t$

$Z \rightarrow e^+e^-$	$Z \rightarrow v_e v_e^-$	<i>Z</i> →uū	$Z \rightarrow d\overline{d}$
$Z\!\! o\!\!\mu^+\!\mu^-$	$Z\!\! o\!\!v_{\mu}v_{\overline{\mu}}^{\overline{}}$	$Z\!\! ightarrow\!c\overline{c}$	$Z \rightarrow s\overline{s}$
$Z{ ightarrow} au^+ au^-$	$Z{ ightarrow}v_{ au}v_{ au}^{-}$		$Z\!\! o\!b\overline{b}$



Lepton Universality ⇒

- \mathcal{M} for e, μ, τ decay is same $\Rightarrow \Gamma(Z \rightarrow e^+e^-) \approx \Gamma(Z \rightarrow \mu^+\mu^-) \approx \Gamma(Z \rightarrow \tau^+\tau^-)$
- \mathcal{M} for v_e , v_μ , v_τ decay is same $\Rightarrow \Gamma(Z \rightarrow v_e v_e) = \Gamma(Z \rightarrow v_\mu v_\mu) = \Gamma(Z \rightarrow v_\tau v_\tau)$

Number of Neutrinos

Total width of the Z-boson (Γ_Z) is sum of all partial widths:

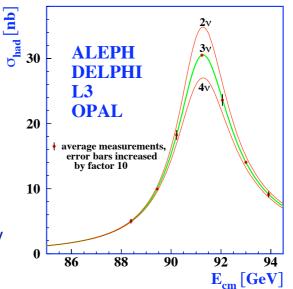
$$\Gamma_Z = \Gamma(Z \to q\bar{q}) + \Gamma(Z \to e^+e^-) + \Gamma(Z \to \mu^+\mu^-) + \Gamma(Z \to \tau^+\tau^-) + N_\nu \Gamma(Z \to \nu\bar{\nu})$$

LEP directly measured:

- partial widths: $\Gamma(Z \rightarrow e^+e^-)$, $\Gamma(Z \rightarrow \mu^+\mu^-)$, $\Gamma(Z \rightarrow \tau^+\tau^-)$, $\Gamma(Z \rightarrow hadrons) = \Gamma(Z \rightarrow q\bar{q})$
- total width, Γ_Z

Cannot measure $\Gamma(Z \rightarrow v \overline{v})$ directly as neutrinos leave no signal in the detector.

- Can predict $\Gamma(Z \to v \overline{v})$ using $\mathcal{M}(Z \to v \overline{v})$
- Use prediction & measurements to find number of neutrinos contributing to Γ_Z
- $N_v = 2.999 \pm 0.011$



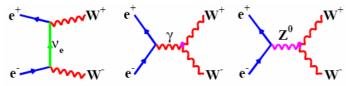
Consistent with exactly three neutrinos!

⇒ Gives us confidence that there are exactly three generations of quarks and leptons

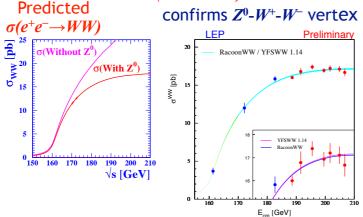
Measuring W-boson properties

At LEP: W-bosons produced in pairs

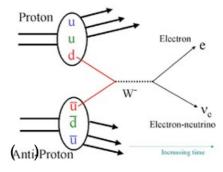
Three possible modes:



 $\sigma(e^+e^- \rightarrow WW)$ measured at LEP confirms Z^0 - W^+ - W^- vertex



At hadron colliders: Wbosons produced from quarks and anti-quarks



From measurements at LEP & **Tevatron:**

- $m_W = 80.413 \pm 0.048 \text{ GeV}/c^2$
- $\Gamma_W = 2.141 \pm 0.041 \text{ GeV}$

Measurements of W and Z bosons validate the Electroweak Model beautifully

Weak Force Summary

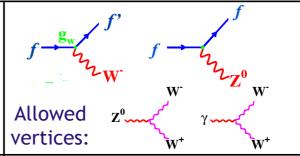
- The weak force acts on all quarks and leptons.
- Two massive bosons propagate the weak interaction: W^{\pm} and Z^{0} .
- Interactions characterised by: long lifetimes 10⁻¹³ - 10³ s and small cross sections 10⁻¹³ mb.

At low energies, virtual W and Z bosons responsible for lepton and lightest hadron decays and neutrino scatterings.

Fermi theory describes *W*-boson interactions for low boson momentum.

Described by Fermi constant: $G_F \propto g_w^2/m_W^2$

High energies colliders can produce real W and Z bosons for study. These validate electroweak model e.g. Z decays suggest exactly three generations of quarks and leptons.



- Lepton interactions are universal: same coupling to *W*, *Z* bosons
- Quarks interactions not universal:
 - → W-(Q=+2/3 quark)-(Q=-1/3 quark) coupling is $g_W V$, where V depends on flavours of quark involved.
 - → *V* is largest within a generation: *V*ud, *V*cs, *V*tb

Electromagnetic & weak are manifestations of a single unified electroweak interaction, described by just 3 parameters.