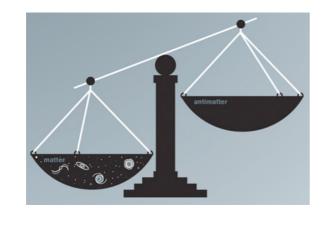
Particle Physics

Dr Victoria Martin, Spring Semester 2012 Lecture 14: *CP* and *CP* Violation

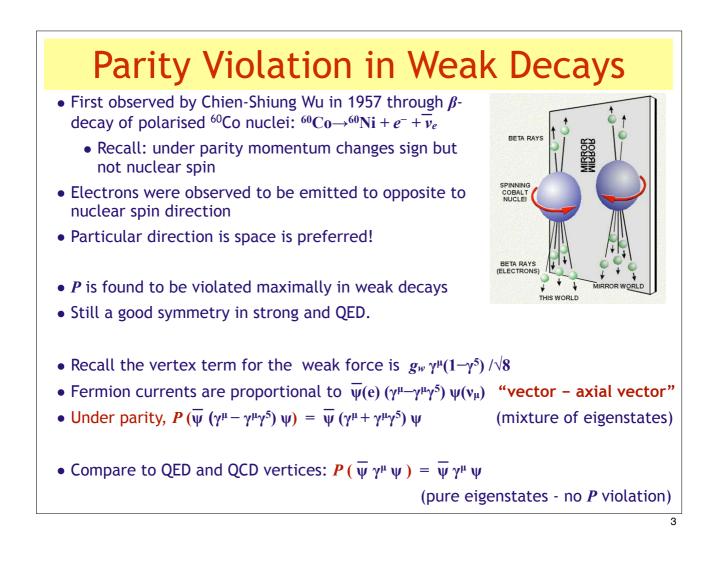


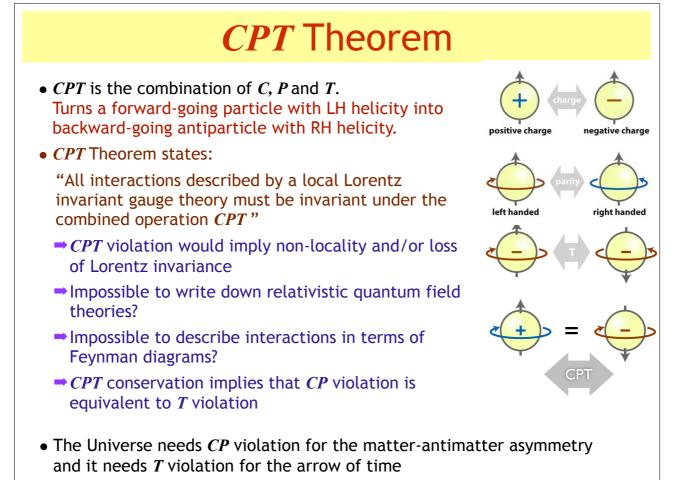
*Parity Violation in Weak Decay
*CP and CPT
*Neutral meson mixing

- ★Mixing and decays of kaons
- **\star***CP* violation in \mathbf{K}^0 and \mathbf{B}^0

Summary from Tuesday

- Symmetries play a key role in describing interactions in particle physics.
- QED and QCD obey Gauge symmetries in the Lagrangian corresponding to symmetry groups. These lead to conservation of electric and colour charge.
- Three important discreet symmetries: Charge Conjugation (*C*), Parity (*P*) and Time reversal (*T*).
 - C: changes the sign of the charge
 - *P*: spatial inversion, reverses helicity. Fermions have *P*=+1, antifermions *P*=-1
 - T: changes the initial and final states
 - Gluons and photons have C = -1, P = -1
- C and P are conserved in QED and QCD, maximally violated in weak
- Only LH neutrinos and RH anti-neutrinos are found in nature.
- *CPT* is thought to be absolutely conserved (otherwise RQFT doesn't work!)

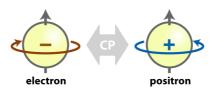




Tests of CPT Invariance**CPT** invariance implies particles and antiparticles must have equal masses: $\frac{M(K^0) - M(\overline{K}^0)}{\frac{1}{2}[M(K^0) + M(\overline{K}^0)]} < 10^{-18}$ **OutputParticle and antiparticles must have equal lifetimes:** $\frac{\Gamma(K^0) - \Gamma(\overline{K}^0)}{\frac{1}{2}[M(K^0) + M(\overline{K}^0)]} < \times 10^{-17}$ **OutputParticle and antiparticles must have equal and opposite charges and magnetic moments** $Q(p) + Q(\bar{p}) < 10^{-21}e$ $\frac{g(e^+) - g(e^-)}{\frac{1}{2}[g(e^+) + g(e^-)]} < 2 \times 10^{-12}$ **Hydrogen and anti-hydrogen atoms have identical spectra**

CP Symmetry

- C is also found to be maximally violated in weak decays.
- Experimental results suggest the combination *CP* is a conserved symmetry.
- *CP* turns a particle into its antiparticle with opposite helicity: it is a symmetry between matter and anti-matter



- CP is a conserved quantity in strong and electromagentic interactions.
- It is *nearly* a conserved quantity in weak interactions, but not quite.
- Violation of *CP* symmetry is required to explain the different between the matter and anti-matter content of the universe.
- We will see that *CP* comes about due to a complex phase in the CKM matrix.

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LHCb and the Matter-Antimatter Asymmetry

• <u>http://www.youtube.com/watch?v=EVQnjuCej6l&feature=related</u>

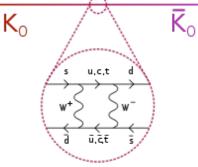
Neutral Meson Mixing

- Second order weak interactions can mix long-lived neutral mesons with their antiparticles:
 - $\stackrel{\bullet}{\rightarrow} K^0 \left(\ \overline{s} \ d \ \right), \ D^0 \left(\ \overline{c} \ u \ \right), \ B^0 \left(\overline{b} \ d \right), \ B_s \left(\overline{b} \ s \right) \qquad K^0 \leftrightarrow \overline{K}^0 \ D^0 \leftrightarrow \overline{D}^0 \ B^0 \leftrightarrow \overline{B}^0 \ B_s \leftrightarrow \overline{B}_s$
- e.g. take the neutral kaons $\mathbf{K}^0 \overline{\mathbf{K}}^0$ as an example:

$P \left \mathrm{K}^{0} \right\rangle = - \left \mathrm{K}^{0} \right\rangle$	$P \ket{\overline{\mathrm{K}}^0} = - \ket{\overline{\mathrm{K}}^0}$
$CP \left \mathbf{K}^{0} \right\rangle = - \left \overline{\mathbf{K}}^{0} \right\rangle$	$CP \overline{\mathbf{K}}^0 \rangle = - \mathbf{K}^0 \rangle$

• The CP eigenstates are:

$$|\mathbf{K}_{1}\rangle = \frac{1}{\sqrt{2}} \left(|\mathbf{K}^{0}\rangle - |\overline{\mathbf{K}}^{0}\rangle \right) \qquad CP = +1$$
$$|\mathbf{K}_{2}\rangle = \frac{1}{\sqrt{2}} \left(|\mathbf{K}^{0}\rangle + |\overline{\mathbf{K}}^{0}\rangle \right) \qquad CP = -1$$



- Decay eigenstates are (approximately) K_1 and K_2 not the same as the flavour eigenstates K^0 and $\bar{K^0}$
- \bullet Two common decay modes are 2π and 3π
 - $\pi^0\pi^0$ and $\pi^+\pi^-$ have CP = +1
 - $\pi^0 \pi^0 \pi^0$ and $\pi^+ \pi^- \pi^0$ have CP = -1

Neutral Kaons continued

- CP is almost conserved in the decay of the neutral kaons.
- Decay $K \rightarrow \pi\pi$ has large phase space \Rightarrow quick decay, travels ~ cm before decay • named "K-short" or K_S with $\tau_S = 0.09$ ns
- Decay $K \rightarrow \pi \pi \pi$ has small phase space \Rightarrow slow decay, travels ~ 10 m before decay • "K-long" or K_L with $\tau_L = 51$ ns
- After a neutral kaon is produced, at some time, *t*, later it will be described by: $\psi(t) = a(t) |\mathbf{K}^0\rangle + b(t) |\overline{\mathbf{K}}^0\rangle$
- The evolution of the kaon state described by mass & decay matrices: M, Γ

$$\begin{split} i\frac{\partial\psi(t)}{\partial t} &= \hat{H}\psi(t) = (\hat{M} - \frac{i}{2}\hat{\Gamma})\psi(t) \\ \hat{M} - \frac{i}{2}\hat{\Gamma} = \begin{pmatrix} M_{\rm K} - \frac{i}{2}\Gamma_{\rm K} & \Delta m_{\rm K} - \frac{i}{2}\Delta\Gamma_{\rm K} \\ (\Delta m_{\rm K})^* - \frac{i}{2}(\Delta\Gamma_{\rm K})^* & M_{\rm K} - \frac{i}{2}\Gamma_{\rm K} \end{pmatrix} \end{split}$$
• Mass difference $\Delta m_{\rm K} = m_{\rm S} - m_{\rm L} = 3.52(1) \ \mathrm{x10^{-12}} \ \mathrm{MeV} = 0.53 \ \mathrm{x} \ 10^{-10} \ \mathrm{s}^{-1} \ \mathrm{is}$ oscillation frequency

- **EXAMPLE 1** The time evolution of the Ks and KL states are: $K_{S}(t) = K_{S}(0) \exp\{-(\Gamma_{S}/2 + im_{S})t\}$ $K_{L}(t) = K_{L}(0) \exp\{-(\Gamma_{L}/2 + im_{L})t\}$ • In terms of the weak (or flavour) eigenstates: $K^{0}(t) = \frac{1}{4} \left[e^{-\Gamma_{L}t} + e^{-\Gamma_{S}t} + 2e^{-\Gamma t} \cos \Delta m_{K}t \right]$ • where $\Gamma = \frac{1}{2} \left[e^{-\Gamma_{L}t} + e^{-\Gamma_{S}t} - 2e^{-\Gamma t} \cos \Delta m_{K}t \right]$ • where $\Gamma = \frac{1}{2} (\Gamma_{S} + \Gamma_{L})$
 - \bullet If you start with beam of neutral kaon (either $K^0,\ \overline{K}{}^0$ or a mixture) it will end up a beam of almost pure K_L

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Neutral Kaons with CP Violation

• The CP eigenstates are K_1 and K_2

$$|\mathbf{K}_{1}\rangle = \frac{1}{\sqrt{2}} \left(|\mathbf{K}^{0}\rangle - |\overline{\mathbf{K}}^{0}\rangle \right) \qquad CP = +1$$
$$|\mathbf{K}_{2}\rangle = \frac{1}{\sqrt{2}} \left(|\mathbf{K}^{0}\rangle + |\overline{\mathbf{K}}^{0}\rangle \right) \qquad CP = -1$$

• The decay K_S and K_L are not quite the same as the decay eigenstates, instead in terms of a small parameter ϵ :

$$\begin{split} |\mathbf{K}_{\mathrm{S}}\rangle &= \frac{1}{N} \left((1-\epsilon) |\mathbf{K}^{0}\rangle - (1+\epsilon) |\overline{\mathbf{K}}^{0}\rangle \right) \\ |\mathbf{K}_{\mathrm{L}}\rangle &= \frac{1}{N} \left((1+\epsilon) |\mathbf{K}^{0}\rangle + (1-\epsilon) |\overline{\mathbf{K}}^{0}\rangle \right) \end{split}$$

• The decay states contain both CP = +1 and CP = -1, CP is violated.

$$|\mathbf{K}_{\mathrm{S}}\rangle = \frac{1}{N} (|\mathbf{K}_{1}\rangle - \epsilon |\mathbf{K}_{2}\rangle)$$
$$|\mathbf{K}_{\mathrm{L}}\rangle = \frac{1}{N} (|\mathbf{K}_{2}\rangle + \epsilon |\mathbf{K}_{1}\rangle)$$

• ε is measured to be $|\varepsilon| \sim 2 \times 10^{-3}$, the amount of indirect *CP* violation due to mixing of different *CP* eigenstates K_1 and K_2

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Summary

- Parity *P* and Charge Conjugation *C* are maximally violated in weak interactions due to vector axial vector structure of interaction vertex.
 - Conserved in strong and electromagnetic interactions.
- The combined symmetry *CP* describes the difference between matter and anti-matter
 - almost a good symmetry in the weak interactions.
- *CPT* symmetry must be conserved... it's one of the foundations of QM and field theory!
- Small amounts of CP violation observed in K^0 and B^0 decays and mixing.
 - (New this year also in charm-mesons: \mathbf{D}^0)
- Three types of *CP* violation:
 - 1. Direct CP violation in decay amplitudes
 - 2.CP violation in neutral meson mixing
 - 3. Indirect *CP* violation due to interference of mixing and decay.
- The amount of *CP* observed is not enough to explain the matter antimatter asymmetry of the universe.